# Tabletop Analog Black Holes to Investigate Information Loss Paradox

### Pisin Chen

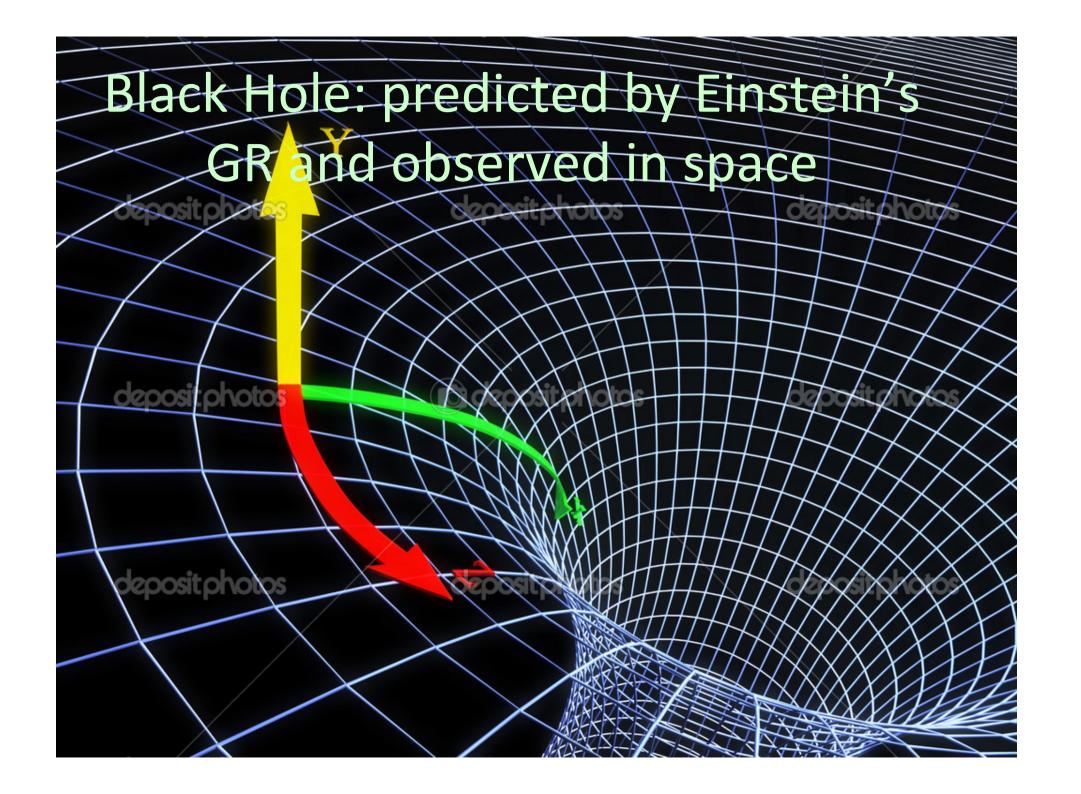
Department of Physics and Graduate Institute of Astrophysics & Leung Center for Cosmology and Particle Astrophysics (LeCosPA)

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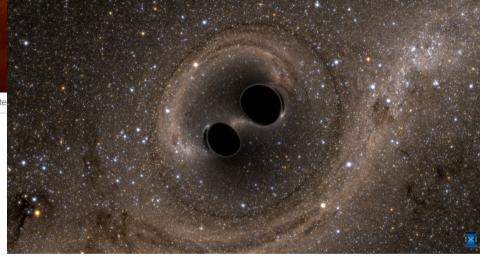
Stanford University

LLNL, July 20, 2017

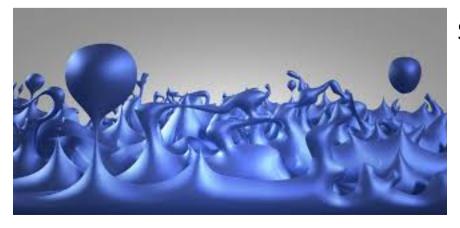


## Three types of black holes

Supermassive BH (100-million to billions solar mass)



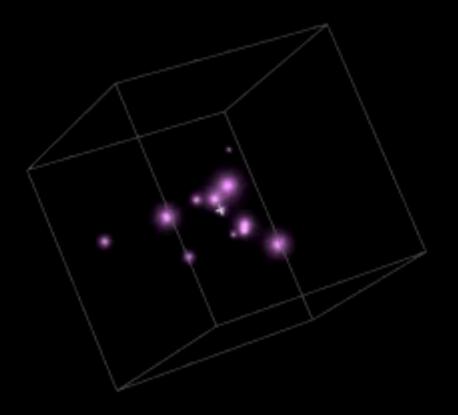
Stellar size BH (several to tens of solar mass)



Primordial BH (born around the big bang, Planck mass and size)

## Supermassive BH at Milky Way galactic center





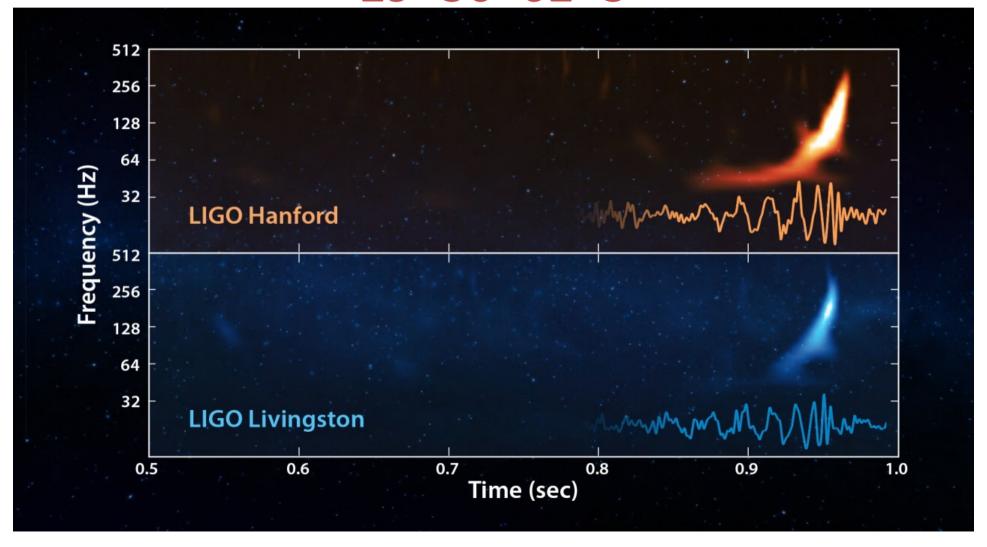
Data: Andrea Ghez, Jessica Lu (UCLA)

Visualization: Dinoj Surendran, Randy Landsberg,

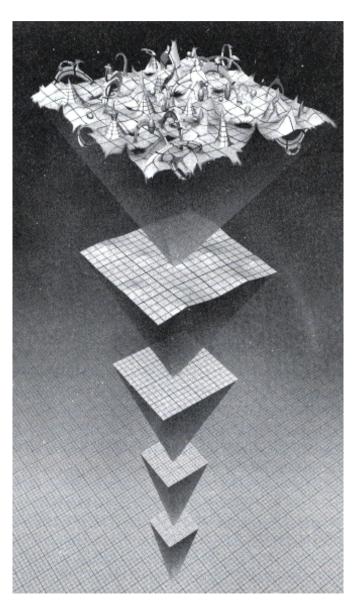


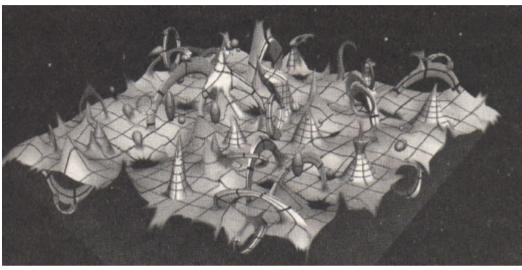
Mark SubbaRoo (UChicago / Adler / KICP)

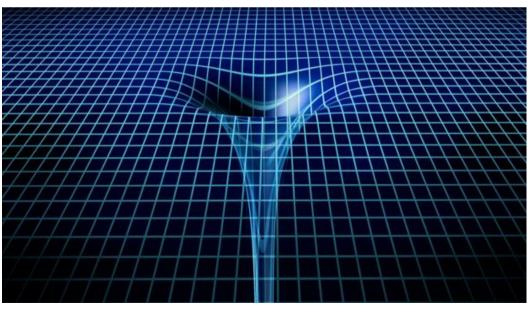
## 2016 discovery of GWs by LIGO: Result of coalescence of two BHs 29+36=62+3



## Primordial BHs born around Big Bang







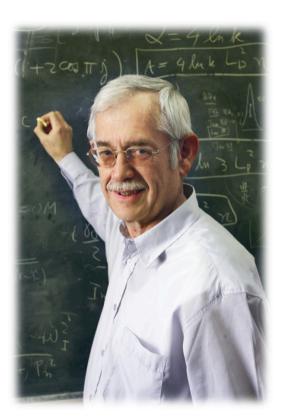
## Black holes have entropy!

• Bekenstein-Hawking BH entropy:

$$S_{BH} = \frac{k_B A}{4l_P^2}$$

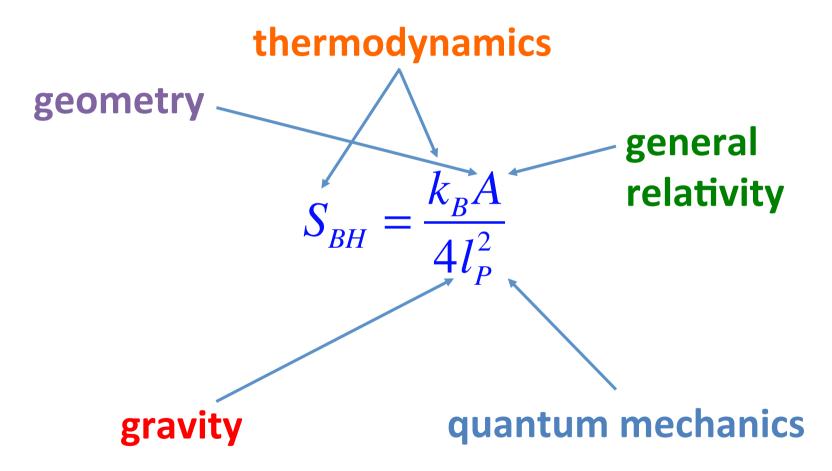


$$E = k_B T$$



Jacob Bekenstein (1947-2015)

## BH entropy: A very unique formula!

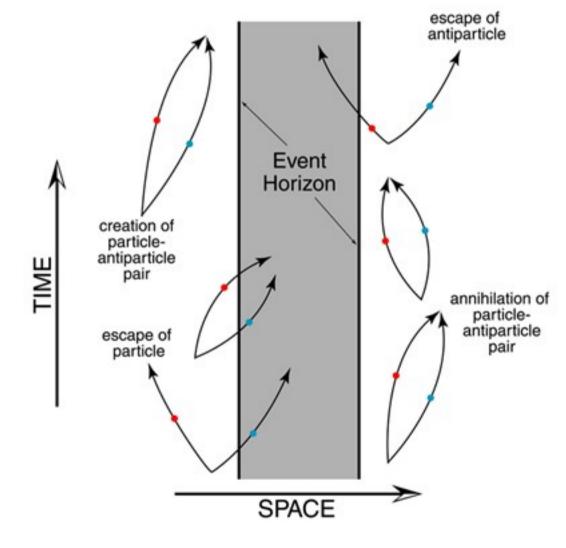


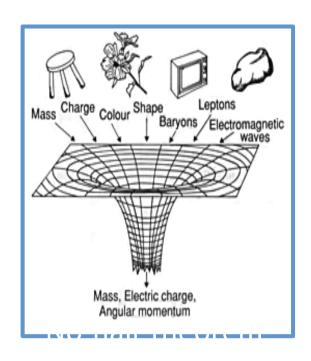
Strangely, it's proportional to BH surface area, not the volume!

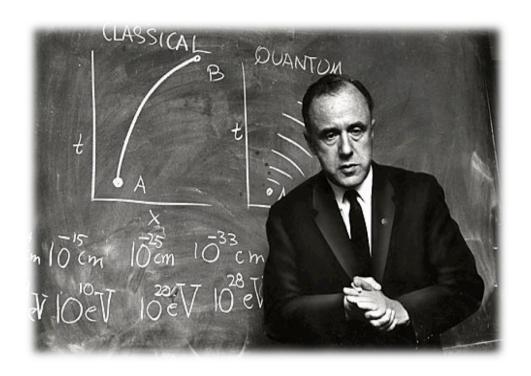
# Black hole Hawking evaporation – Connecting GR, QM, SM in one stroke

 $r_{s} = \frac{2GM}{c^{2}} \qquad g = \frac{GM}{r_{s}^{2}}$  $k_{B}T_{H} = \frac{\hbar c^{3}}{8\pi GM} = \frac{\hbar g}{2\pi c}$ 





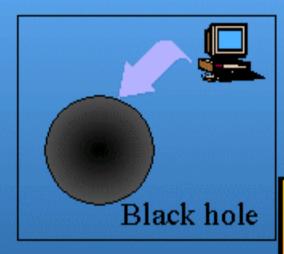




### John Wheeler:

What would happen if I dropped my coffee mug into a black hole?

### **Information Loss**



Heat

- Entanglement of Hawking radiation?
- Firewall?
- Fussball?
- Etc., etc., etc.







End point

40 years after Hawking's discovery, the nature of BH evaporation is still under debate!

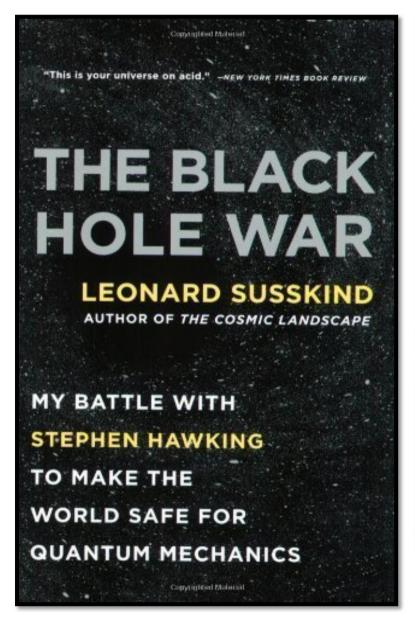
## Hawking evaporation may result in the loss of information!

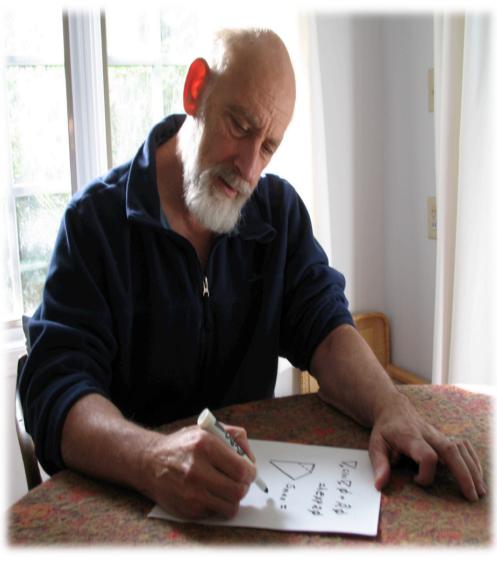
- First pointed out by Hawking himself in 1976
- Endless debates ever since
- Solutions include "black hole complementarity" (Susskind et al.), Firewall (AMPS, AMPSS), etc.
- Entanglement between Hawking radiation and partner particles Wilczek 1987, Schutzhold-Unruh 2010, Hotta-Schutzhold-Unruh (2015)
- Planck size black hole remnants (Chen-Ong-Yeom, Phys. Rep.2015)
- Naked black hole firewalls (Chen-Ong-Page-Sasaki-Yeom, PRL 2016)
- BMS Soft Hairs (Hawking-Perry-Strominger, 2016)
- No firewalls & nothing wrong w. information loss (Unruh-Wald 2017)
- An alternative hairdo based on Kac-Moody symmetry (Addazi-Chen-Marciano-Wu, 2017)

## Can Hawking radiation carry out information after all?



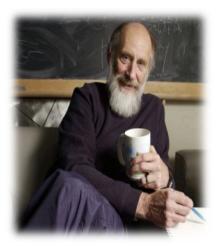
## Recommended reading





## **BH** Completentarity

(Nothing wrong with a process as long as it is consistent with causality).



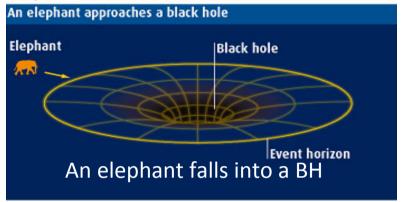


New Scientist, Oct. 2006,

Leonard Susskind, Lárus Thorlacius, John Uglum, "The Stretched Horizon and Black Hole Complementarity"

#### HOW TO BE IN TWO PLACES AT ONCE

Black holes challenge our notion of locality. If no observer can see information lost in the universe, then a thought experiment shows that an object's location in space-time depends on whether its observer is accelerating outside the event horizon or free-falling inside. That means information can be in more than one place at the same time



Observer A sees the elephant get closer and closer to the horizon, while observer B sees the elephant pass through



A sees it continues to approach horizon B sees it crossing the horizon unharmed

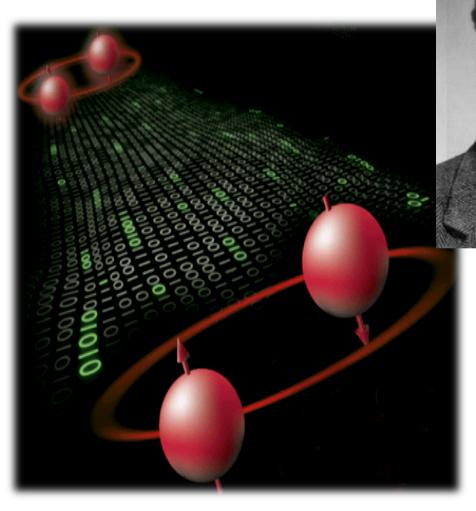
Observer A sees the elephant get thermalised and radiated back out, while B sees it continue to its doom



## Quantum entanglement

Schrödinger: "Verschrankung" (1935) as a result of

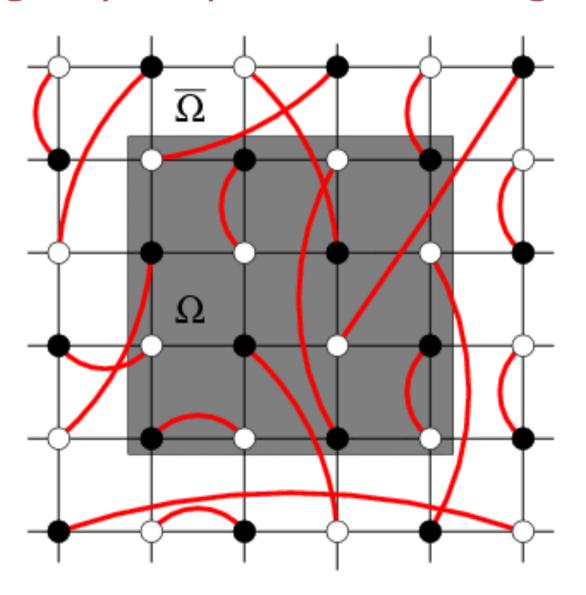
discussing with Einstein



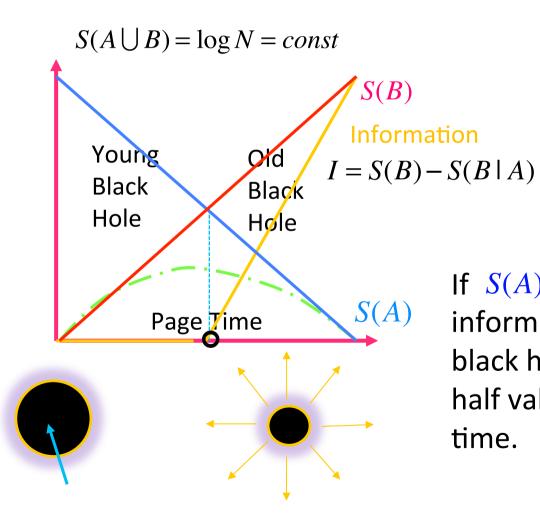
A B

"Quantum entanglement is not just a property of QM, it is THE character of QM. It fundamentally breaks QM from classical physics." (E. Schrodedinger)

## Monogamy of quantum entanglement



### When would BH information come out?



$$S(B) = \log m$$

$$S(B \mid A) = \sum_{k=n+1}^{mn} \frac{1}{k} - \frac{m-1}{2n}$$

(For a pure and random system, Conjectured by Page, 1993; proved by Sen, 1996.)

If  $S(A) \propto Area$ , then the information will come out when the black hole initial area decreases to half value. This is called the Page time.

In 2012, four physicists (AMPS) argued that the 3 basic assumptions that led to the BH complementarity principle, namely,

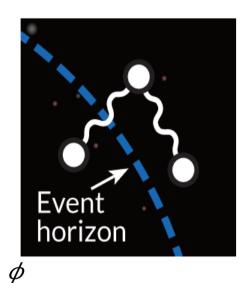
- 1. Unitarity
- 2. Local quantum field theory
- 3. No drama

cannot be all consistent. They suggested that the "most conservative" solution would be that there exists a firewall on the BH surface, anything falls into BH would

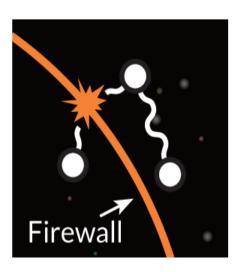
be burned into ashes.

## AMPS black hole firewall

#### **Problem**



#### **Solution: Firewall**

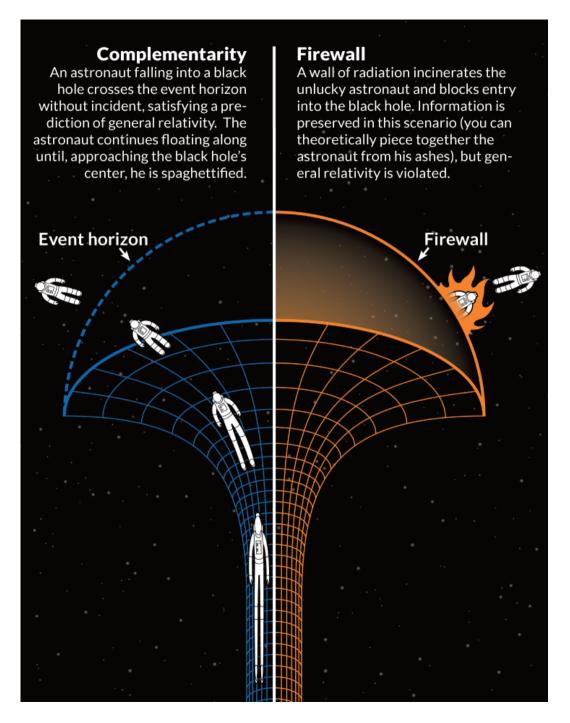


- Ahmed Almheiri,
   Donald Marolf, Joseph
   Polchinski, James Sully,
   "Black Holes:
   Complementarity or
   Firewalls?",
- Donald Marolf, Joseph Polchinski, Douglas Stanford, James Sully, "An Apologia for Firewalls",

Ahmed Almheiri,

The intensity of a quantum field is Determined by the rate of change of the field

For disconnected spacetimes, the magnitudes of the quantum field need not be continuous.



### General relativity:

For a sufficiently large BH, whose curvature is small, objects should pass its horizon uneventfully— "No Drama"

### **AMPS** firewall:

The requirement that Hawking radiation can bring information out from BH would result in the notion of firewall.

## Yukawa Institute of Theoretical Physics, Kyoto University



The five authors of the paper with another colleague during the discussion at the Yukawa Institute for Theoretical Physics: (L to R) Dong-han Yeom, Yen Chin Ong, Pisin Chen, Don Page, Yasusada Nambu, and Misao Sasaki.

## Chen-Ong-Page-Sasaki-Yeom: Why should firewalls be naked?

Phys. Rev. Lett. 116, 161304 (2016)

- Quantum fluctuations in Hawking radiation inevitable
- BH's backreaction to the quantum fluctuations leads to the exposure of the event horizon.
- Curvature of stellar-size BHs small, so GR should be satisfied
- Firewalls conjecture is not a conservative solution to the information paradox

## Investigations of ILP mostly theoretical Astro black holes too cold and too young

Lifetime of solar mass BH: 10<sup>67</sup> years

Age of the universe: 1.38 x 10<sup>10</sup> years



## **Analog Black Holes**

- Sound waves in moving fluids "dumb holes"
   Unruh (1981, 1995)
- Traveling index of refraction in media Yablonovitch (1989)
- Violent acceleration of electron by lasers Chen-Tajima (1999)
- Electromagnetic waveguides
   Schutzhold-Unruh (2005)
- Bose-Einstein condensate
   Steinhauer (2014)
- Accelerating mirror

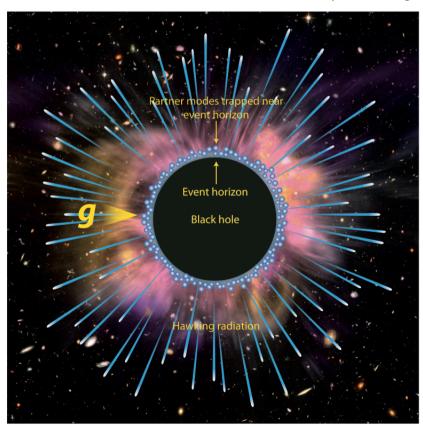
Fulling-Davies (1976), Davies-Fulling-Unruh (1977), Birrell-Davies (1982), Carlitz-Willey (1987), Hotta-Schutzhold-Unruh (2015), Chen-Mourou (2016), Chen-Yeom (2017)

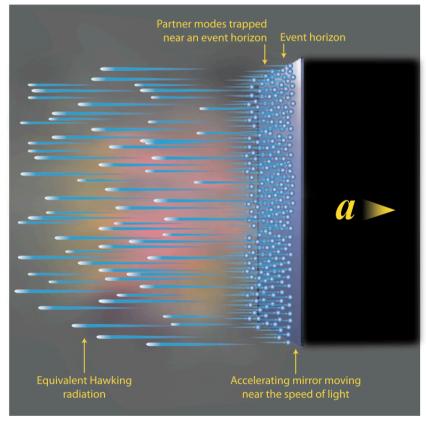
Testing thermal nature of Hawking radiation

## Accelerating mirror as an analog black hole

#### SIMULATING A BLACK HOLE ON A TABLE

New black hole simulator may shed more light on a contradiction in fundamental physics

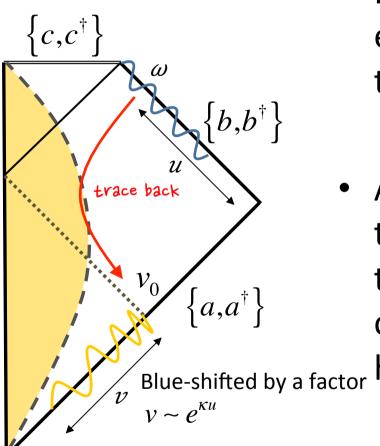




Black hole Hawking evaporation

Accelerating mirror as an analog black hole

## Why is there a radiation?



 Hawking radiation is theoretically estimated by using the Bogoliubov transformation.

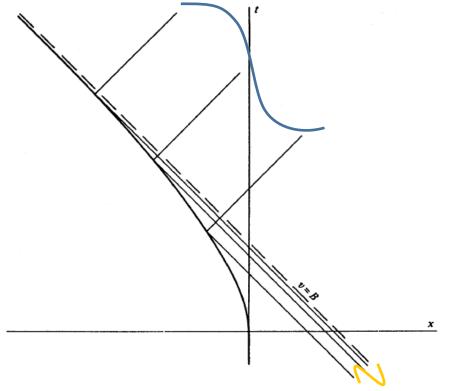
$$\langle n_{\omega} \rangle = \langle b_{\omega}^{\dagger} b_{\omega} \rangle = \sum_{\omega'} \left| \beta_{\omega\omega'} \right|^2$$

 A non-trivial Bogoliubov transformation is possible due to the red-shift of incoming modes caused by the horizon of a black hole

$$p_{\omega} \sim e^{i\omega\kappa^{-1}\ln\left[(v_0 - v)/c\right]} \qquad (v > v_0)$$

$$= \sum_{\omega'} \left[\alpha_{\omega\omega'} f_{\omega'} + \beta_{\omega\omega'} f_{\omega'}^*\right]$$

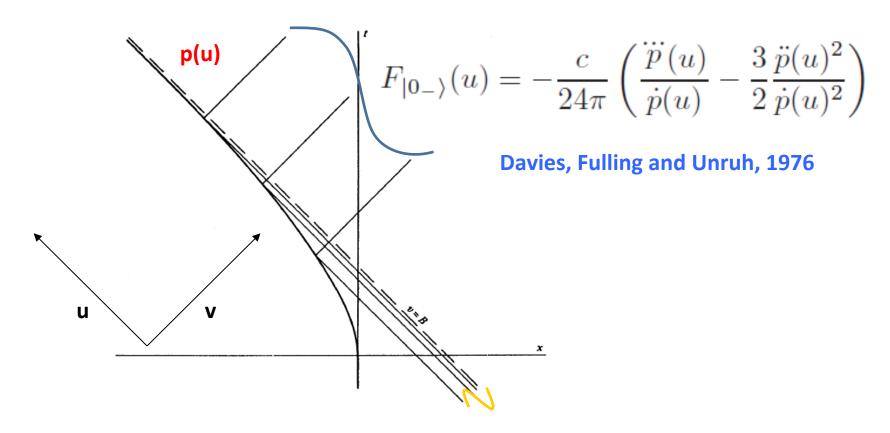
## Red-shift by a mirror



Fulling and Davies, 1976
Birrell and Davies, 1982

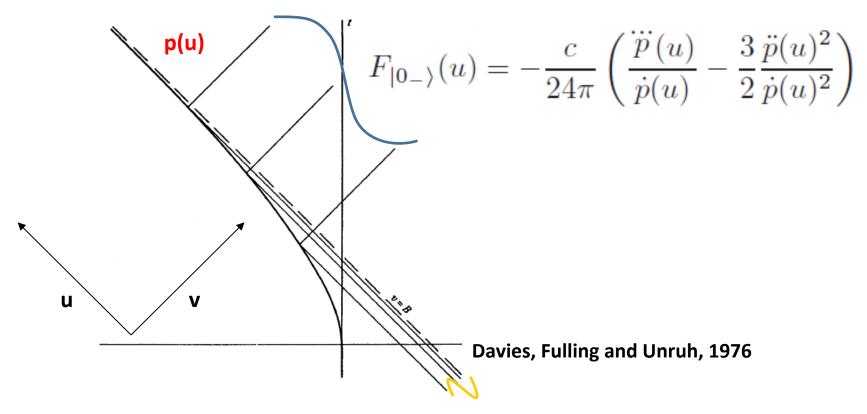
- A moving mirror is a surface that satisfies a reflecting boundary condition.
- If the mirror is moving with a constant acceleration, then it generates a thermal radiation.

## Red-shift by a mirror



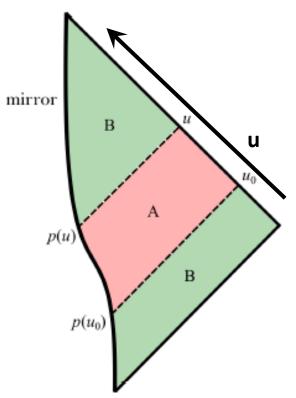
 One can calculate the out-going energy flux as a function of the mirror trajectory (for 2D spacetime).

## Is there information loss?



- Definitely, there should be no information loss in the mirror dynamics.
- Then what can we learn about the entanglement entropy?

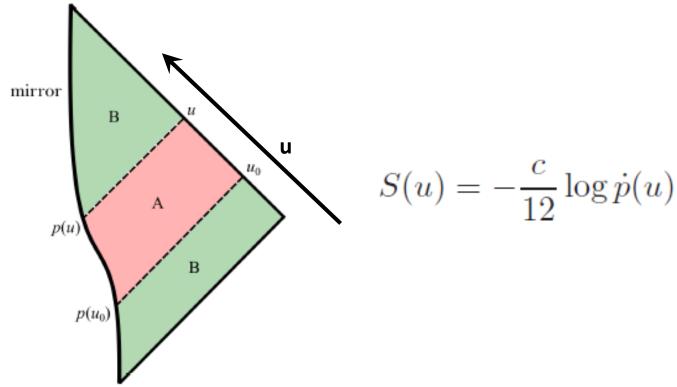
## Entanglement entropy



$$S_{|0+\rangle}^{\epsilon}(A) = \frac{c}{12} \log \frac{(u-u_0)^2}{\epsilon^2}$$

- In order to apply Page's argument, one can calculate the entanglement entropy as a function of u.
- In order to obtain a finite result, renormalization of the cutoff is needed.

## Entanglement entropy formula

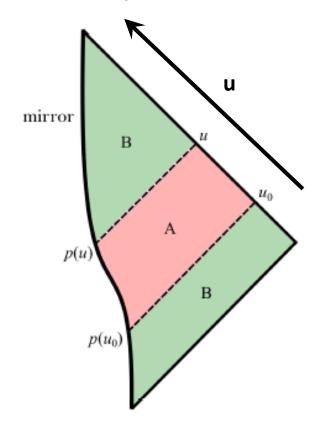


- After a proper renormalization, we obtain the formula (Holzhey, Larsen and Wilczek, 1994; Bianchi and Smerlak, 2014).
- Several authors have tested the consistency of this paper, e.g., Abdolrahimi and Page, 2015.

## Mirror trajectories

P Chen, D-h Yeom, "Entropy evolution of moving mirrors and the information loss problem", arXiv:1704.08613 (Accepted for publication PRD)

 Using this formula, we can test several candidate trajectories.

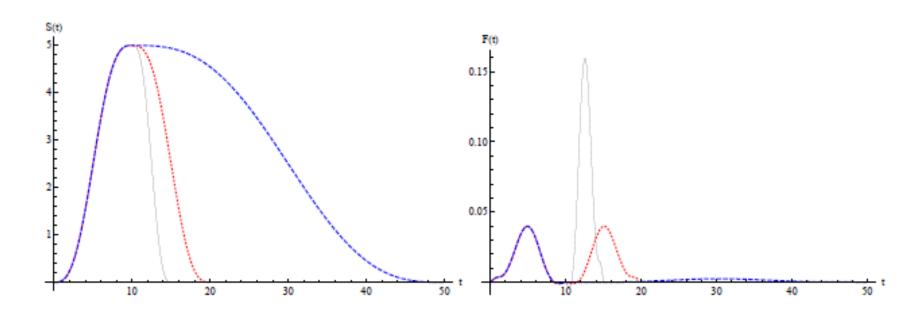


$$\frac{dS(t)}{dt} = A \sin^2 \pi \frac{t}{t_P} \qquad 0 \le t < t_P,$$

$$= -A \frac{t_P}{t_f - t_P} \sin^2 \pi \frac{t - t_P}{t_f - t_P} \qquad t_P \le t < t_f,$$

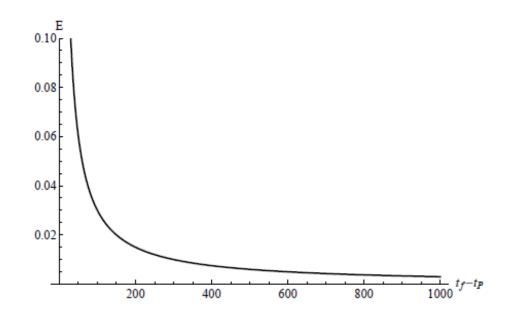
- Suddenly stopping mirror:  $t_f = 15$ ,
- Slowly stopping mirror:  $t_f = 20$ ,
- Long propagating mirror:  $t_f = 50$ .

## Test of scenarios



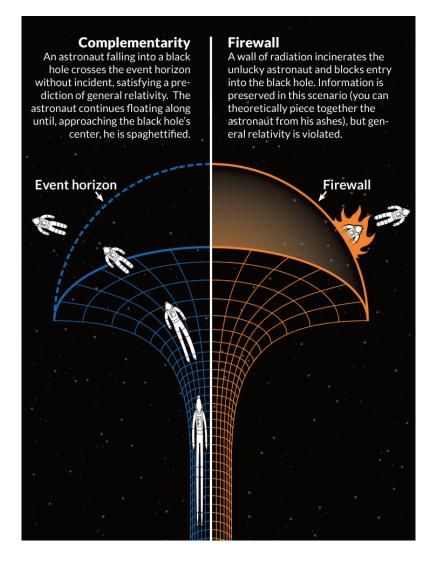
 For a suddenly stopping mirror, there is a large amount of energy emission. In general it is too large and hence it cannot mimic the last burst of a black hole.

### Test of scenarios

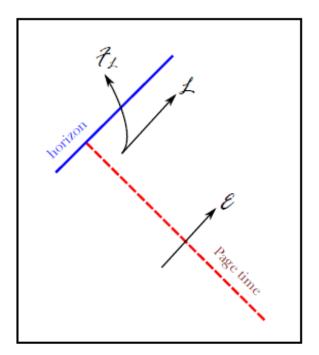


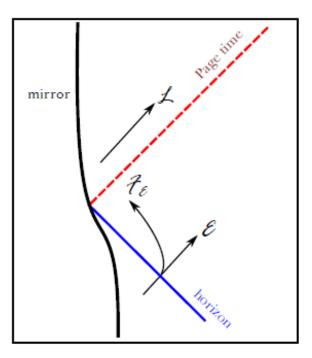
- For a mirror with very long lifetime, the emitted energy can be arbitrarily small.
- This mimics the possibility of correlation between vacuum and radiation or the remnant scenario.

#### Complementarity vs. Firewall



# Consistency test: AMPS thought experiments

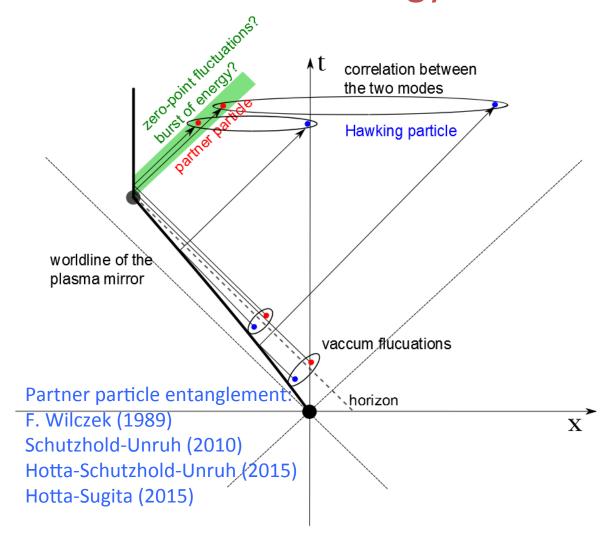




- Due to the monogamy of entanglements, there should be an effect that breaks one link.
- For a black hole case, this is called by (hypothetical) the firewall.
- There should be a violent effect from a mirror: a firewall-like emission?

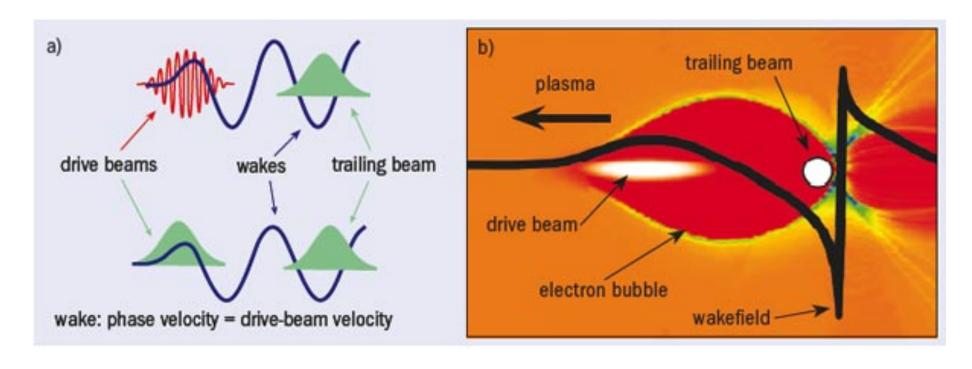
#### Flying Mirror:

Entanglement between Hawking & partner particles Final outburst of energy or not?



#### Plasma Wakefield Acceleration

Tajima-Dawson (1979)- Laser driven (LWFA)
Chen-Dawson-Huff-Katsouleas (1985)- Particle beam driven (PWFA)



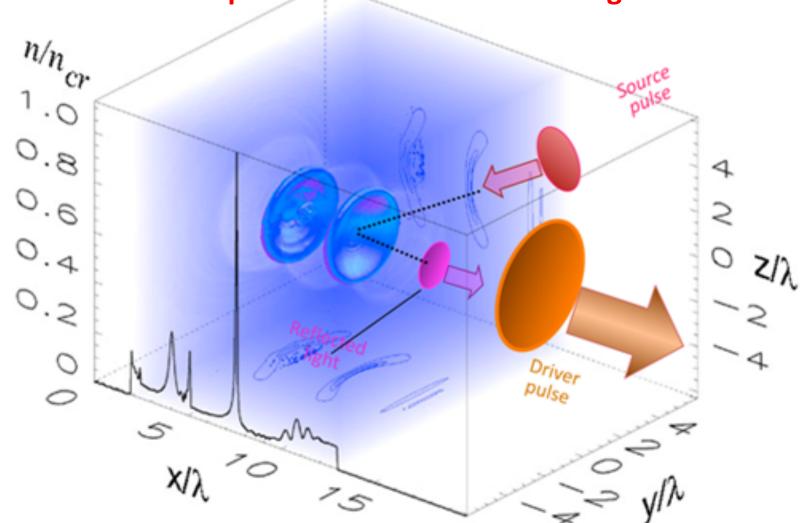
SLAC & LBL- Acceleration of O(100) GeV/m observed! AWAKE- A new experiment at CERN

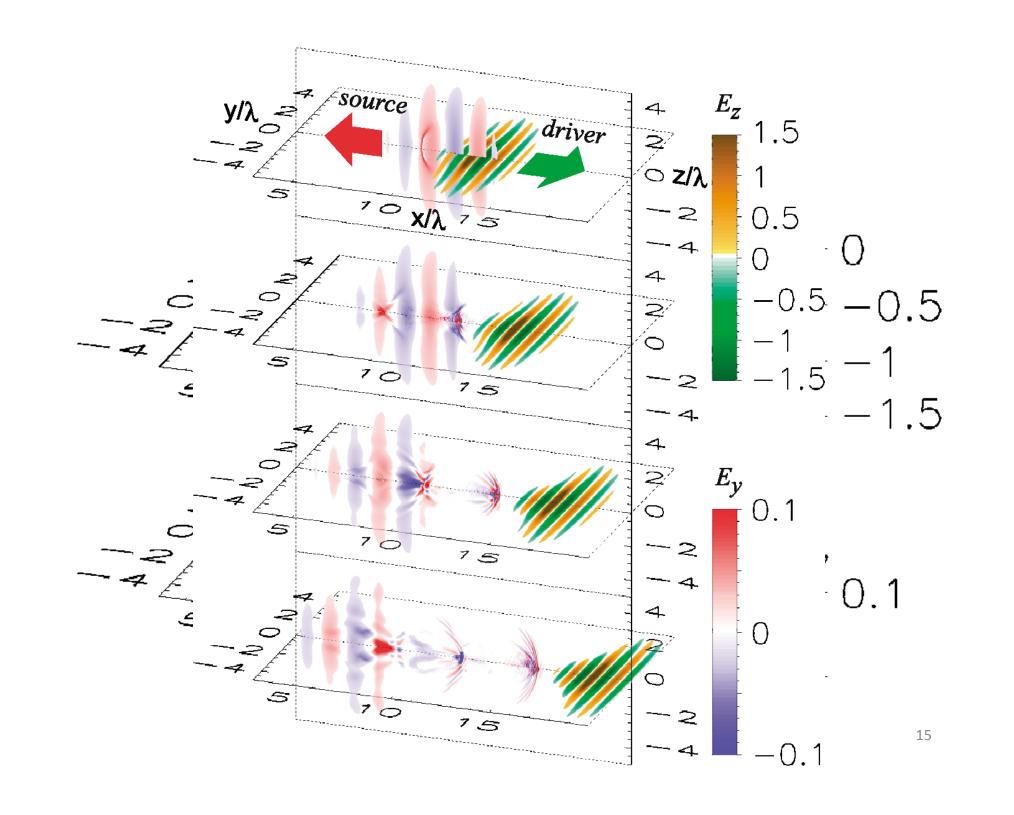
#### Plasma Wake is like a tsunami



#### Relativistic Plasma Mirror

Reflected laser pulse Lorentz-boosted and tighter-focused.





### An accelerating plasma mirror

P Chen, G Mourou, "Accelerating plasma mirrors to investigate black hole information loss Paradox", Phys. Rev. Lett. 118, 045001 (2017).

 For uniform plasmas, the plasma wakefield, i.e., the relativistic mirror, is induced instantly by the impinging laser, under the *Principle of Wakefield*:

#### "Plasma wakefield follows the driver by one wavelength."

 Nonlinear plasma wakefield is described by the (normalized) scalar and vector potentials \$\phi\$ and \$a\$ by the coupled equations

$$\left[\frac{2}{c}\frac{\partial}{\partial \chi} - \frac{1}{c^2}\frac{\partial}{\partial \tau}\right]\frac{\partial a}{\partial \tau} = k_{p0}^2 \frac{a}{1+\phi},$$

$$\frac{\partial^2 \phi}{\partial \chi^2} = -\frac{k_{p0}^2}{2} \left[1 - \frac{(1+a)^2}{(1+\phi)^2}\right].$$

# Natural tendency of laser deceleration due to wakefield excitation

 The deceleration (or redshift) of the laser (and therefore the mirror) is governed by

$$\frac{\partial \omega}{\partial \chi} = -\frac{1}{2} \frac{\omega_p}{\omega} \frac{\partial}{\partial \chi} \frac{1}{1+\phi}.$$

• Let us model the laser envelope as

$$a_L(\chi) = a_{L0} \sin\left(\frac{\pi\chi}{L}\right), \qquad -L \le \chi \le 0.$$

Then the solution is

and 
$$\phi \simeq \frac{a_{L0}^2 k_p^2}{8} \left\{ \chi^2 - 2 \left( \frac{L}{2\pi} \right)^2 \left[ 1 - \cos(2\pi \chi / L) \right] \right\}.$$

$$\frac{\partial \phi}{\partial \chi} \simeq \frac{a_{L0}^2 k_p^2}{4} \left\{ \chi - \frac{L}{\pi} \sin\left(\frac{2\pi \chi}{L}\right) \right\} < 0.$$

### Acceleration of the plasma mirror

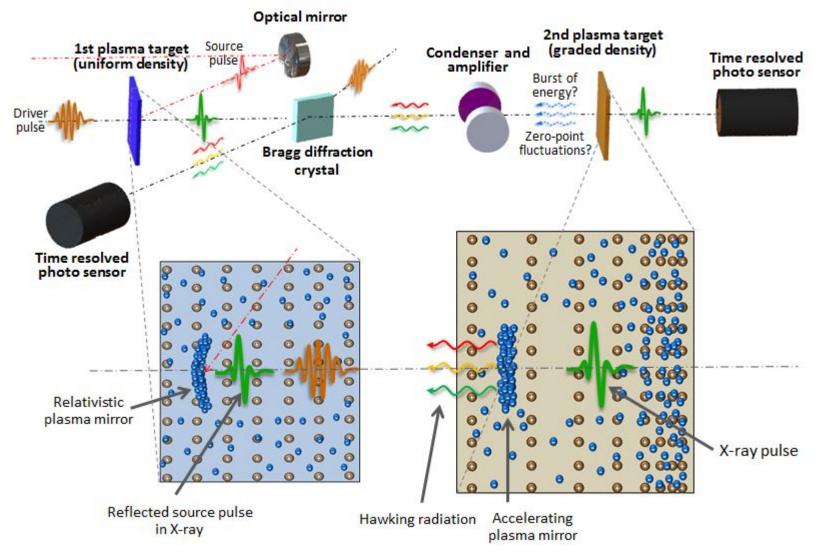
Invoking the "wakefield principle" in nonuniform plasmas,

$$\ddot{x}_{M} = \ddot{x}_{g} - \ddot{\lambda}_{p} = v_{g} \frac{\partial v_{g}}{\partial x} - \ddot{\lambda}_{p} = \eta c^{2} \frac{\partial \eta}{\partial x} - \ddot{\lambda}_{p}.$$

where the refractive index  $\eta = \sqrt{1 - (\omega_p^2 / \omega^2)/(1 + \phi)}$ , Carrying out the calculations, we find

$$\ddot{x}_{M} = \frac{c}{2\eta_{0}} \left[ v_{g} \left( 1 + \frac{\omega_{p0}^{2}}{\omega^{2}} \right) \frac{\omega_{p0}^{2}}{\omega^{2}} \frac{\partial}{\partial x} \frac{1}{1 + \phi} \right] \left( 1 + \frac{\partial \omega_{p}}{\partial x} \frac{t}{k_{p0}} \right) + c\eta_{0} \left( \frac{\partial \omega_{p}}{\partial x} \frac{1}{k_{p0}} + \frac{\partial^{2} \omega_{p}}{\partial x^{2}} \frac{v_{g}t}{k_{p0}} \right).$$
Due to frequency redshift

# A conceptual design of the accelerating plasma mirror experiment



### Plasma density variation

 Invoking nano-fabrication technology for solid plasma targets with, for example, a power-law increase of density:

Si<sub>0.8</sub>Ge<sub>0.3</sub> layer

Si<sub>0.8</sub>Ge<sub>0.2</sub> layer

Si<sub>0.8</sub>Ge<sub>0.1</sub> layer

50 nm

$$n_{p0}e^{\pm x/D}$$
,  $n_{p0}(1+x/D)^{2(1-\eta_0)}$ ,  $0 \le x \le X$ , otherwise

Then the acceleration is

$$\ddot{x}_{M} = \frac{(1 - \eta_{0})c^{2}}{D(1 + x/D)^{2}} \exp\left(\frac{(1 - \eta_{0})x/D}{1 + x/D}\right), \quad 0 \le x \le X.$$

### Example

The 4 length scales should satisfy the inequality:

$$\lambda_x \ll \lambda_p \ll D \ll X.$$
  $(\lambda_x \simeq 7.79nm)$ 

Plasma target based on nanotechnology with

$$\lambda_{p0} = 7.8nm$$
  $D = 10nm$ , thickness  $X = 2D$ , and density  $n_{p0} \sim 5 \times 10^{23} cm^{-3}$ 

- Mirror velocity:  $v_M(0) \sim 0.01c$   $\longrightarrow$   $v_M(2D) \sim 0.997c$
- Reflectivity of plasma mirror at this frequency:  $Y \approx 1$ 
  - Corresponding Hawking temperature:

$$k_B T_H(x) \simeq \frac{\hbar c}{4\pi D} \frac{\omega_{p0}^2}{\omega_0^2} \frac{1}{(1+x/D)^2} \exp\left\{\frac{(1-\eta_0)x/D}{1+x/D}\right\} \sim 1.6 - 0.1 eV.$$

#### Background noise not severe

- One salient feature of this experiment:
  - The Hawking signals propagate backward,
  - whereas most x-ray or optical laser induced background particles would move forward.
- Since the x-ray energy 25 eV  $<< m_e = 0.5$  MeV, Compton backscattering induced by x-ray would have similar frequency at 25 eV
- Bragg diffraction crystal is designed to let pass the 25 eV but divert the 1-10 eV photons, these background signals would therefore be directed to a different path.
- In conclusion, the background in this experiment should be minute.

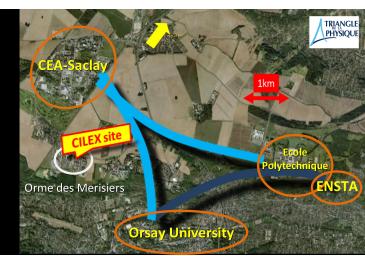
# AnaBHEL Collaboration formed (Analog Black Hole Evaporation via Lasers)

National Taiwan University + Ecole Polytechnique + Kansai Photon Research Inst. + Shanghai Jiao Tong U.

#### Two stages:

- 1. Proof of principle at KPRI Laser facility, presently one of the most powerful lasers in the world @ PW
- 2. Full scale expt. with 10PW APOLLON laser, Saclay when completed in 2018

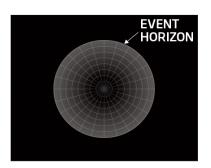




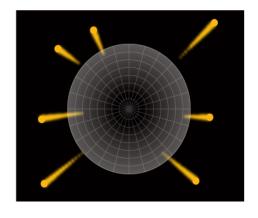
1916 General Relativity predicted A Brief History of BHS the existence of BH

Brief history of black holes

1916 **Black holes** emerge from general relativity: nothing, not even light, escapes the event horizon



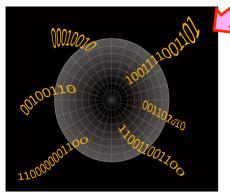
1974 Black holes emit **Hawking** radiation thanks to quantum mechanics



served

2004

Hawking accepts that information escapes from black holes



1974

Hawking theoretically discovered BH evaporation

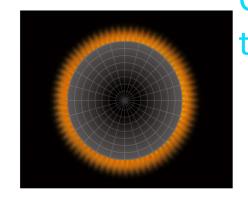
2004

Hawking agreed the

2012 **Escaping** information ignites firewall. which can't be

general relativity

2012 reconciled with AMPS conjectured firewalls



Can the BH war be settled through experiments

2016

COPSY argued that firewalls are naked

#### What can we learn from AnaBHEL?

- First, if we can detect the **thermal radiation**, then in itself it confirms QFT in curved spacetime.
- Second, if we can experimentally measure the entanglement entropy of radiation (this is a challenge) before, during, and after the acceleration, then it can test the renormalization method for the entanglement entropy.
- Third, one can expect that as the mirror stops, there maybe violent effects on the mirror: firewall-like burning mirror?

## Summary

- Hawking evaporation and information loss paradox is one of the fundamental problems in physics.
- So far investigations are essentially theoretical;
   Direct observation of black hole end-stage unlikely.
- Quantum entanglement between Hawking radiation and partner particle may reveal the secrete.
- Accelerating plasma mirrors may serve to address some aspects of this paradox experimentally.