#### Laser-driven coils, how well do they work?

1000 35866 35870 18 Maximum B flux density @ 850 µm (T) 35869  $\lambda_{\rm L} = 1.053 \ \mu {\rm m}$ 35868 100  $\lambda_{\rm L} = 0.53 \,\mu{\rm m}$ **Faraday rotation** measurement Pick-up probe measurement 10 5 6 7 8 9 4 5 6 7 8 9 3 10<sup>15</sup> 10<sup>16</sup> Intensity (W/cm<sup>2</sup>)

> Left: General laser driven coil configuration Center: Result from Fujioka *et al.* claiming impressive, kT magnetic fields Right: Proton radiograph of the center of the laser driven coil

TER

J. Peebles April 1<sup>st</sup>, 2021 HEDS Seminar



J. Davies, D. Barnak, M. Bonino, G. Brent, T. Cracium, R. Betti

Laboratory for Laser Energetics University of Rochester

Thanks to the Office of Fusion Energy Sciences for supporting this work and the Center of Excellence for Advanced Nuclear Diagnostics and Platforms for ICF and HED Physics at OMEGA, NIF and Z

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#### **Outline**



- Laser driven coils (LDC) overview
- How well do LDCs work?
- Electric and magnetic field diagnostics
- Calibrating and validating diagnostics
- Overview of initial and follow up experimental results
- Discussion and conclusions



#### **Outline**



- Laser driven coils (LDC) overview
  - Motivation
  - Principles of operation
- How well do LDCs work?
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#### Strong magnetic fields open new avenues of research in high energy density physics and are highly desirable for inertial confinement fusion research

- LLE's magnetic field generator MIFEDS is heavily subscribed for shots on OMEGA and EP, despite complications in brings to campaigns (~8 shot days in Q4 FY20 alone)
- LLNL is pursuing a project to create its own pulsed power device for NIF to be compatible with their cryogenic system and debris requirements
- Pulsed power systems produce debris that are incompatible with many high performing laser systems and experiments



Single shot MIFEDS debris



#### The Catch-22 with magnetized experiments

- Magnetized experiments use high intensity lasers to generate energetic particles for physics goals or to probe magnetic fields
  - I<sub>laser</sub> > 10<sup>19</sup> W/cm<sup>2</sup>, pulse length = best compression
- Pulsed power devices such as MIFEDS can potentially damage optics, debris shields are required which drop laser intensities

- I<sub>laser</sub> < 10<sup>18</sup> W/cm<sup>2</sup> at BC, ~10<sup>19</sup> W/cm<sup>2</sup> at 10 ps

• Coils are bulky and can block beams preventing physics goals. Magnetic fields are generated at the expense of the rest of the experiment



Not sure there's enough here for a paper...





## Laser driven coils are target based platforms used to generate strong magnetic fields without a pulser



- Early laser experiments measured current traveling through target support structures and the chamber wall\*
- While current was low, it was hypothesized that electrons ejected from targets lead to a neutralizing return current through the target



Proton radiographs taken of a target stalk at different stages of a direct drive laser implosions by M. Manuel inferred a stalk current of 7 kA with large electric field

\*R. F. Benjamin et al., Phys Rev Lett **42** (1978) M. Manuel et al., Appl. Phys. Lett. **100**, 203505 (2012)



### Laser driven coils are target based platforms used to generate strong magnetic fields without a pulsed power device



- Conventional laser driven coils (LDCs) use a laser to drive charge separation, which draws a current to create a field in a loop of wire\*
- This solves many problems by limiting debris is generated and providing magnetic fields to facilities with no pulser





\*H. Daido, et al., Phys. Rev. Lett. 56 846 (1986)

### LDC experiments are often modeled after electronic circuits and are frequently referred to as "capacitor coils"

V<sub>c</sub> I<sub>l</sub> Laser R



#### C. Goyon, et al., Phys. Rev. E 95, 033208 (2017)

thermal electrons that are escaping the potential barrier [30, 31]. A fraction of them are captured by the opposite disk. The target reacts like an RL-circuit to the potential difference between the disks and the subsequent discharge current through the coil-shaped wire. The laser-driven target charging and the discharge current process simultaneously (steady-state regime), during the laser-pulse irradiation, establishing a quasi-static

#### J. J. Santos et al., New J. Phys., 17 083051 (2015)

#### B. J. Zhu *et al.,* Appl. Phys. Lett. 107, 261903 (2015)

With the establishment of an initial potential  $V_0$  between the two disks, the target behaves toward the potential like a resistor–inductor (*RL*) electrical circuit. The time evolution of the current I(t) can be treated using an *RL* model:

Coil 
$$V_0(t) = L \frac{dI(t)}{dt} + RI(t),$$
 (5)  
Laser

Figure 1. Basic geometry of the capacitor-coil target

## LDC experiments are often modeled after electronic circuits and are frequently referred to as "capacitor coils"



#### "Capacitor coils" are a misnomer

The physical dimensions of this plate have no basis on the "capacitance" of this target

$$C = \frac{\varepsilon A}{d} \approx 0.1 \text{pF}$$

Capacitor Energy =  $\frac{1}{2}CV^2 = 0.5 \text{ mJ}$  for 100 kV charge

The plates are considered <u>"current sources"</u>

Circuit theory also makes certain assumptions, such as <u>current</u> <u>uniformity</u>, truncated circuit size and Ohm's law



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## Ongoing research shows interest in laser driven coils (LDCs), in particular with coils producing kiloTesla fields

hit the front target. Strong currents are driven in the two parallel coils connected to the target due to the voltage between the two targets. A magnetic field is generated around the coils through Ampere's law. Experiments show that the magnetic field generated by this method can reach a kilotesla [14–17]. There is a spacer between the coils and the target to prevent the laser-produced high-energy plasma from affecting reconnection, the reconnection plasma is generated by irradiating another laser beam to a foam target. In this setup, the plasma beta can reach a low value ( $\sim$ 0.1).

Kai Huang et al., Plasma Phys. Control. Fusion 63 015010 (2021)

#### All samples from articles released in 2020

As already stated, a strong magnetic field that can confine the relativistic electrons is required to realize this scheme experimentally. A laser-driven capacitor-coil [25], [26], [27], [28], [29], [30] can produce a 1-kT-level magnetic field. The pulse duration of this magnetic field is about 1 ns, sufficiently long relative to the time scale of proton acceleration. It was confirmed in the previous experiments that the magnetic field generated by a laser-driven magnetic field can indeed radially confine relativistic electrons within a small spot [31], [32].

H. Morita, et al., High Energy Density Physics 37 1008742 (2020)

A numerical study on the pulse duration dependence of a magnetic field generated using a laser-driven capacitor-coil target

Toru Sasaki 🎗ª⊠, Kazumasa Takahashi ª, Takashi Kikuchi ª, Atsushi Sunahara <sup>b</sup>, Hideo Nagatomo <sup>e</sup>, Shinsuke Fujioka °

Show more 🗸



### Though kT numbers are often cited, it is only one measurement in a sea of experiments that have been performed

	Driving Laser					Coil Design			+	$[\tau_p \text{ (ns)}]$
Ref	E(kJ)	t (ns)	$\lambda$ (nm)	$I \ (10^{15} {\rm W/cm^2})$	d (mm)	Mat.	$r (\rm{mm})$	PRad	B-Dot	Faraday Rot
Daido (1986)	0.1	1.0	10000	0.13	0.3, 0.5	$\mathbf{Cu}$	1.0	_	15, 28	_
Daido (1986)	0.1	1.0	10000	0.13	0.7, 1.0	$\mathbf{Cu}$	1.0	_	41, 16	_
Santos(2015)	0.5	1.0	1057	100	0.9	Cu, Ni, Al	0.25	-, 95 [0.35], -	800, 600, 150	-, 450 [0.2], -
Fujioka (2013)} (	@(0.095, 0.27)	1.3	526	4, 0.8	0.78	Ni	0.25	-,-	$^{-,33}$	160, -
<sup>-</sup> ujioka (2013) 2	@(0.37, 0.5)	1.3	1053	16, 43	0.78	Ni	0.25	_	_	480, 1100
_aw (2016)	0.54	2.0	1053	21	0.5	Unknown	0.25	610 [2.3]	620 [2.0]	_
Courtois (2005)	0.3	1.0	1053	40	0.5	Cu	1.25	_	7.5	_
Courtois (2005)	0.022	0.08	526	40	0.5	$\mathbf{Cu}$	1.25	_	7.5	_
Tarifeño (2008)	0.0005	7.0	1064	0.0007	1.3	$\mathbf{Cu}$	1.0	_	0.001	_
Gao (2016)	2 @ 1.25	1.0	351	30	0.6	$\mathbf{Cu}$	0.3	45[4.1]	_	_
Goyon (2017)	2 @ 0.375	0.75	351	4	0.5	Au + CH	0.25	200[1.0]	_	_
Wang (2018)	0.01	$3 \times 10^{-5}$	800	1000	0.5	Ni	0.5	20 [0.05]	_	_
Zhu (2015)	0.18	2.0	527	0.14	Inf	$\mathbf{Cu}$	0.55	50 [unk.]	_	_
Zhu (2015) 🛛 $8  {@}$	(0.063, 0.125)	2.0	351	1.4, 2.8	Inf	$\mathbf{Cu}$	0.55	_	45, 73	_
Zhu (2015) 🛛 $8  {@}$	(0.222, 0.246)	2.0	351	5.0, 7.7	Inf	$\mathbf{Cu}$	0.55	_	102, 205	_
Matsuo (2017)	1.3	1.2	351	1.5	0.5	Ni	0.45	_	219	206
Korobkin (1979)	0.001	20	1053	_	0.1	$\mathbf{Cu}$	0.7	_	2	_
lvanov (2020)	0.022	2.8	532	8	0.5	$\mathbf{Cu}$	0.5	_	_	6.4
Ivanov (2020)	0.022	0.07	532	25	0.5	$\mathbf{Cu}$	0.5	_	_	12.4



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	]	Driving Laser				Coil Design		Field (T)	+	$[\tau_p \text{ (ns)}]$
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J. Peebles et al. Phys. Plas. 27 063109 (2020)

### The authors of the > 1 kT result declare it breaks energy conservation throwing suspicion on the Faraday results



ROCHESTER

### 1 kT is not feasible, what about the 600+ T measurements from other experiments?



Experiments by Law et al. and Santos et al. measured fields of > 600 T

For reference, MIFEDS (480 J stored energy) currently reaches 100 T, with coils made as small as possible (~2 mm<sup>3</sup> magnetized volume)



## Field energy density provides some insight into whether quoted field values and applications are reasonable



• Energy in the field is absolutely limited by the energy in the driving laser

Field Energy Density 
$$= \frac{|B|^2}{2\mu_0} + \frac{|E|^2}{2\varepsilon_0}$$
  
Energy in  $B = \frac{1}{2\mu_0} \int B^2 dV = \frac{LI^2}{2}$ 

- Lasers deliver comparable energy as pulsed power capacitors (~< 1 kJ)
- To be competitive with pulsed power, LDCs have one advantage for generating higher peak fields



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- Lasers deliver comparable energy as pulsed power capacitors (~< 1 kJ)
- To be competitive with pulsed power, LDCs have one advantage for generating higher peak fields
- Magnetized volume can be made much smaller in LDCs



## Field energy density provides some insight into whether quoted field values and applications are reasonable



- MIFEDS for example can typically have a efficiency of 60% (peak field energy/stored energy), due mainly to not being critically damped
- Finding the peak conversion efficiency of a laser driven coil is not so simple:
  - Fraction of laser energy converted to hot electrons (~ 30%)
  - Fraction of electrons that escape the plate
  - Fraction of electrons that escape backwards towards the undriven plate\*
  - Fraction of electrons are absorbed by the undriven plate\*
- With less energy available, <u>magnetizing large volumes (such as</u> <u>implosions) is not an option</u>

For this reason nearly all laser driven coil experiments have coils with a diameter < 1 mm



### Energy estimates indicate that some scrutiny should be applied to 600+ T results



I = 290 kA for 610 T in loop as described Magnetic Field Energy > 200 J Laser Energy < 540 J Conversion laser to field > 37%

K. F. F. Law et al., Appl Phys Lett., 108 091104 (2016)



I = 375 kA for 800 T in loop as described Magnetic Field Energy > 250 J Laser Energy < 500 J Conversion laser to field > 50%

J. J. Santos et al., New J. Phys., 17 083051 (2015)





## The range of results have large differences between similar experiments

 Experiments by Law et al.<sup>1</sup> and Courtois et al.<sup>2</sup>, give very different results using similar diagnostics and driving laser parameters (same Iλ<sup>2</sup>)



Courtois' loop was 5 times larger, but this does not account for nearly 100 times less field



# The range of results, even within a single experiment, indicate that they can't all be true

• Santos et al.<sup>3</sup> measured fields of 95, 450, 600 T for the same type of coil using radiography, Faraday rotation and b-dot probes respectively



A more thorough investigation and comparison of diagnostic techniques is required



LLE

#### **Outline**



- Laser driven coils (LDC) overview:
- How well do LDCs work?
- Electric and magnetic field diagnostics
  - B-dot probes and Faraday rotation
  - Transverse proton radiography
  - Axial proton radiography
- Calibrating and validating diagnostics
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- Probe diagnostics such as Faraday rotation and B-dot probes require material to be placed in the vicinity of the coil
- These diagnostics often fail due to blanking and EMI so peak fields can rarely be measured
- These tools also inadvertently measure fields from the laser interacting with the disk and other sources





- Placing the diagnostic far from the coil requires a significant degree of extrapolation and assumption on field geometry (5 orders of magnitude)



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- Placing the diagnostic far from the coil requires a significant degree of extrapolation and assumption on field geometry (5 orders of magnitude)
- It also assumes current is perfectly uniform along the wire and no other sources of voltage





- Protons traveling transverse to the coil should be deflected by the axial field and create a void
- There is ambiguity in what field causes the creation of a proton void, electric and magnetic fields can both duplicate features seen
- The void actually significantly decreases the information gained about conditions in the center of the loop



\*Santos *et al* NJP 17, 083051 (2015)

LLE

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Another way to view it is the transverse probe is too sensitive for 100s of Tesla

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LLE





#### An obvious experiment would be to reverse the current on the coil

• The transverse probe deflection is asymmetric, if the current is reversed by switching a wire the pattern should flip!

- This would demonstrate how much deflection is from the B field vs E field
- Nothing has been <u>published</u> attempting this simple experiment





### Magnetic fields typically require much more field energy to create proton voids than electric fields

- Some experiments, especially with high intensity lasers, produce very large voids
- Simulating a void around a coil with electric fields (blue) required 25 J of energy, with magnetic fields (magenta) 183 J, 7 times the energy!
- Attributing all the deflection to magnetic fields often produces an unreasonable energy conversion





### Axial proton radiography can address all of the concerns for transverse radiography for laser driven coils

- Electric fields along the coil will cause protons to be focused or defocused
- Weaker, radial magnetic fields should cause a rotation in the mesh
- Even when the two effects are combined, they can be decoupled
- Information is gained about conditions inside the center of the coil





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# Axial probing of a MIFEDS coil confirms functionality and provides information on the proton energy spectrum



Angle of rotation was calculated for the expected coil and current for each proton energy Reversing current with MIFEDS is trivial



### Radiographs of coil and inverted current measure anticipated rotation



Overall mesh rotation is the same for the same proton energies, simply inverted





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Overall mesh rotation is the same for the same proton energies, simply inverted



## Bulging feature is a consequence of coils not being perfect (requiring an ingress and egress)



LLE



## Bulging feature is a consequence of coils not being perfect (requiring an ingress and egress)







# Bulging feature is a consequence of coils not being perfect (requiring an ingress and egress)







### We can compare rotation changes based on proton energy to simulations



STER

Energy (MeV)	Angle
5	18.0 X
7	15.0 X
9	13.0 √
10.5	12.0 √
14.5	10.5 🗸
20	8.0 √

Axial radiography is accurate enough to tell us that the first 2 films are dominated by shadowing from the 3rd



## Faraday rotation has been tested in the lab as a technique to verify fields from MIFEDS





Coil designed for 9 T in the center

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Original TGG crystal performed well but blanks quickly in a laser environment



## A custom B-dot probe was designed, constructed and calibrated for multiple shots on OMEGA-EP





### A custom B-dot probe was designed, constructed and calibrated for multiple shots on OMEGA-EP





# A custom B-dot probe was designed, constructed and calibrated for multiple shots on OMEGA-EP



The B-dot is looking at an assumed time derivative of B field

An E field imparting a signal is not the derivative w.r.t. time (not an E-dot probe!). Time integrating the signal and attributing it to changing B field would lead to significant error

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## Initial experiments only tested the axial probe on different types of LDCs





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\*J. Peebles et al., PoP 27, 063109 (2020)

#### The lack of current for the double-plate coil is explained by expanding plasma from both plates measured by $4\omega$ angular filter refractometry\* (AFR)

- X-rays from the driven plate indirectly drive the other plate
- From a voltage perspective: equal charge displacement leads to no voltage difference and no current
- From a circuit perspective: any current would jump the gap between plates as a "short-circuit" since plasmas are conductors



\*D. Haberberger et al., CLEO Technical Digest, Paper ATu3M.3

UR



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#### New experimental goals as a consequence

- Verify axial proton probing results by using multiple diagnostics on the same shot
- Try to remove as many sharp corners and burrs from the target as possible to reduce E field enhancement
  - Previously laser cut from foil, move to smooth magnet wire and foils
  - Previously used a 5/6 loop with bends, switch to U shaped 1/2 loop to have no bends
  - Place the stalk of the target on the opposite side of the loop from the plate on the single plate targets
  - Measure current in both directions on a single shot



# Two configurations with two types of targets were used on the subsequent experimental campaign



Double radiography configuration



#### B-dot + Faraday rotation



# Transverse proton radiography on the single plate coils provided a very accurate current measurement

- Single plate coils are really just a coiled stalk in this experiment, the helix provides two half loops in one target for transverse probe purposes
- · Effectively we get to see the effect of current and if the current were reversed





Transverse

Axial (no mesh)



# Transverse proton radiography on the single plate coils provided a very accurate current measurement

- Single plate coils are really just a coiled stalk in this experiment, the helix provides two half loops in one target for transverse probe purposes
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Orientation of transverse proton probe





Orientation of transverse proton probe

20 MeV experimental radiograph





10 kA, 20 MeV synthetic radiograph

20 MeV experimental radiograph

UR





10 kA, 20 MeV synthetic radiograph

20 MeV experimental radiograph



# 20 kA is beyond the upper limit for current measured by proton radiography



20 kA, 20 MeV synthetic radiograph

20 MeV experimental radiograph

UR LLE

20 kA current bulges and focuses more than in experiment



### Estimating current passing through the target and stalk ("infinite wire") using the B-dot probe





#### If we attributed all signal close to t0 to dB/dt, what is the field?



#### Back of the envelope calculation

In the lab 2.5 V measured with probe 1.6 cm from 1 kA wire (should be ~0.0117 T)

Equivalent to 30-40 V at 20 dB (lab attenuation)

Coil is tilted ~45° relative to stalk

Assume some attenuation due to system bandwidth (30%)

(40/2.5)\*sqrt(2)\*1.4\*0.0117 T ≈ 0.4 T

Using B =  $\mu_0 I/2\pi$ ,  $I_{stalk} < 30 \text{ kA}$ 

# Adding the short pulses for proton radiography introduces significant signal that cannot be magnetic field

On shots with the short pulses the B-dot probe measured significantly higher spikes in signal ~1 ns after the long pulse

It is unlikely that this is due to magnetic field since 300 kA through a stalk would be more energy than the short pulse laser



# Faraday rotation measured no significant rotation in fused silica media compared to other regions in the probe



Light is split into || and  $\perp$  polarizations

The Faraday effect in a medium with a refractive index causes polarization rotation



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The Faraday effect in a medium with a refractive index causes polarization rotation







Polarization 2 Summed image



# If no significant rotation is measured, how sensitive is the diagnostic?





Edges of Faraday medium showed changes of  $\pm 2\%$ 

If attributed to field this corresponds to a rotation of  $\pm 1.5^{\circ}$  or  $\pm 1.23$  T integrated over the 800 µm Faraday medium



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Faraday estimates at most 7.5 kA, with significant assumptions





#### Proton radiography: **I** = 10 ± 2.5 kA

• Still could use more work on introducing E fields, but opposing loops restricts this condition

B-Dot probe: I < 30 kA (unknown how much is from electric field)

 Nothing is currently known about electric field coupling on ns timescales to the probe so this can only impose an upper limit

#### Faraday rotation: I = 0 - 7.5 kA (for lower energy shots)

• Depending on where you measure in the Faraday glass there is either no rotation or very slight rotation, likely due to noise



# Double plate proton radiography does not make a convincing argument for high current





Two bulges at the loop can be seen in the 36 MeV transverse radiograph, though bulge size is consistent with ~ 5 kA current



Rotation should be apparent with > 10 kA current No rotation measured



#### Transverse radiographs still require more analysis

E < 15 MeV protons are focused by net negative charge (E ~10<sup>7</sup> V/m), obscuring the coil E = 20-30 MeV protons have a void at the tip, caused primarily by small E and B field E > 30 MeV protons are not significantly affected by either field around the coil





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#### Transverse radiographs still require more analysis

E < 15 MeV protons are focused by net negative charge (E ~10<sup>7</sup> V/m), obscuring the coil E = 20-30 MeV protons have a void at the tip, caused primarily by small E and B field E > 30 MeV protons are not significantly affected by either field around the coil





## B-dot probe measurements of the double plate showed similar results to the single plate, with less energy





### Integration makes the B-dot probe look a little more convincing than raw data



LLE





#### Proton radiography: **I** = **1.5** ± **5 kA**

 Axial probe restricts this to < 10 kA, transverse probe suggests 1.5 kA, but the shape mismatch suggests more electric field involvement

#### B-Dot probe: I < 30 kA (unknown how much is from electric field)

A similar voltage was measured to the single plate though with less drive energy

Faraday rotation: N/A

• Faraday glass blanked on shots with 1250 J



#### Outline



- Laser driven coils (LDC) overview:
- How well do LDCs work?
- Electric and magnetic field diagnostics
- Calibrating and validating diagnostics
- Overview of initial and follow up experimental results
- Discussion and conclusions



### There are a few key differences between these experiments and others



 $I\lambda^2$  (W/cm<sup>2</sup>)·( $\mu$ m<sup>2</sup>) is lower for these experiments than those with very high field measurements

Comparing  $I\lambda^2$  across most experiments there appears to be little correlation with current or field







### Very recent experiments on the Vulcan laser corroborated our result on OMEGA-EP

- P. Bradford *et al.* performed experiments with a similar dual proton probe geometry and intensity but with **1 μm** light
- Axial radiographs found no rotation and a limited current was inferred from transverse radiography







P. Bradford et al., High Power Laser Science and Engineering, Vol. 8, e11, (2020)

# Target construction and target stalk placement vary across experiments

Initial motivation stemmed from drawing a current through the target stalk, what happened to target stalks for these experiments? Most designs place it on the driven plate (acting as a current divider), or do not specify

Targets have varying degrees of smoothness in their construction geometry



### Finally, the ultimate problem with LDCs: magnetizing an experiment

- The laser driven coil is not independent from the experiment it is magnetizing
- Driving the LDC creates a huge x-ray source that ٠ irradiates the target
- To keep fields high and inductance low, experiments ٠ must be placed close to the driving plate





### Finally, the ultimate problem with LDCs: magnetizing an experiment

- <u>The laser driven coil is not independent from the</u> <u>experiment it is magnetizing</u>
- Driving the LDC creates a huge x-ray source that irradiates the target
- To keep fields high and inductance low, experiments must be placed close to the driving plate
- Shielding the experiment is completely ineffective as it provides a clear short circuit path if any current was present
- Should be noted that any experiments that magnetize a target with an LDC likely significantly preheat the target





# Even if LDCs functioned as well as the most optimistic publication suggests, there are limitations

- Since laser driven coils have a rise time on the order of the laser pulse, dB/dt is extremely high, leading to a high EMF
- Conductive targets will generate opposing currents very quickly to stop field penetration
- These induced currents coupled with resistive heating will cause the target to blow up before an experiment can be performed







#### Let's revisit the question: How well do laser driven coils work?



#### Not very well, and you have to jump through these hoops:

- Have very efficient (>30%) coupling to electrons
- Make a <u>single</u> plate LDC small enough so that inductance is kept low
- Since the inductance must be kept low, the main experiment must be shielded from x-rays generated by a kJ laser 1-2 mm away
- Make sure your shielding is < 1 mm in size and doesn't become part of the circuit
- Make sure your main experimental target is < 1 mm in size and non-conductive
- Have an experiment where you don't mind if the field is non-uniform and don't mind if the experiment is preheated







#### **Extra Slides**





- Measure the current and electric field from short pulse driven stalks using a mesh fiducial (has been done in the past but somewhat inconclusive)
- Faraday rotation using 4ω polarimetry for a MIFEDS coil
- Measure EMI from different material targets and stalk designs



# Future experiments will take lessons learned to better understand proton radiography and diagnosing fast rising current





### B-dot and Faraday rotation diagnostics will be compared concurrently on the same shot





# Synthetic radiographs were constructed with a leap-frog particle pusher in Matlab

- Current and charge are explicitly placed to allow for convoluted current geometries
- A 3D field map for the entire simulation box (6x6x10 mm) is explicitly calculated using Coulomb and Biot-Savart laws

- Particles are uniformly sent from the proton source with a given energy towards the target
- Particles that overlap with a reference mesh grid, with the coil material or leave the box are removed
- Particles that reach the end of the box have their trajectory extrapolated to the film
- Small magnification changes are made afterwards due to varying location of film in the RCF stack



# Revisiting assumptions: does the assumption of Ohm's law and steady state current apply to LDCs?

- A non-uniform current leads to charge build up in the coil
- Ohm's law is frequently used to justify that current in the coil must be uniform, but our experiment data indicates otherwise
- Ohm's law is an empirical law, which is not necessarily true in all circumstances
- The <u>Drude model</u> of electron transport indicates how Ohm's law comes from electron-ion collisions in conductors
- In an AC spatially uniform field the Drude model shows material tends toward plasma like behavior (response time is that of an electron plasma period)



LLE

Graphic from Wikipedia showing how electron inertia impacts current development in a fast rising E field



#### What does the Drude model say about the LDC environment?

- An electric field is generated at the plate which propagates at the speed of light
- Electrons in the wire material will be accelerated by the electric field
- Inertia causes the electron response to have a rise time, rather than be instantaneous
  - The rise time is on the time scale of an electron plasma period
- Electrons accelerated collide accelerate collide, until an equilibrium is reached, several plasma periods long
- These current transients in our system should generally travel several mm/ns
- The electric field is not steady temporally or spatially in the LDC, further complicating the establishment of a steady state current

It seems likely that LDCs are transient dominated systems which would be difficult to model due to system size



# The single plate system had anomalous, asymmetric current which appeared to rise <u>after</u> the laser turned off

- The dipole electric field established by the laser affects both wires due to coil geometry leading to at least 2 current transients
- Turning the laser off may change the conditions seen by the wires leading to a change in current flow and new transients forming

