

Igniting, by most definitions

Achieving a burning plasma on the National Ignition Facility Laser

HEDS Seminar Series (reprise of APS-DPP 2021 talk)

Debra A. Callahan (she/her)
Associate Division Leader for HED-ICF Design

On behalf of the ICF program

December 16, 2021

LLNL-PRES-828821

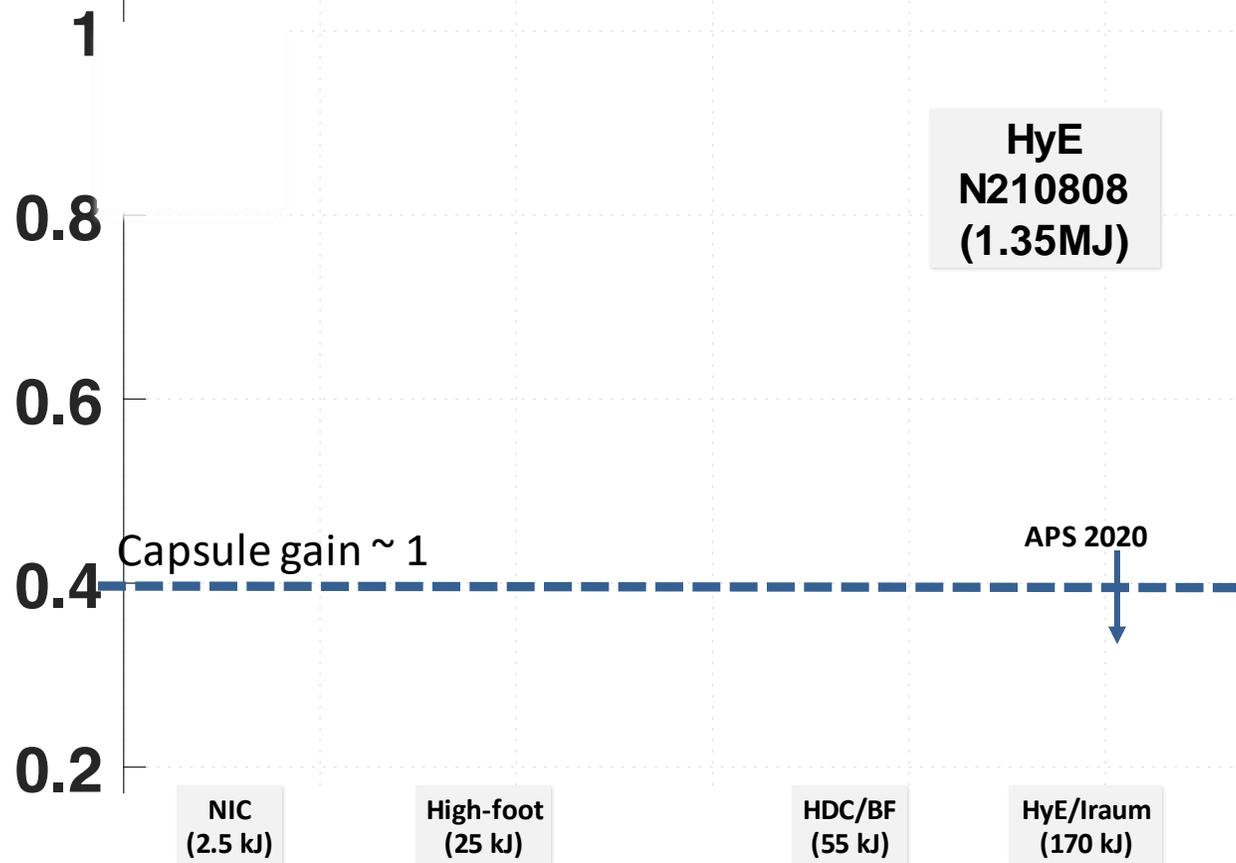
This work was performed under the auspices of the U.S. Department of Energy by Lawrence Livermore National Laboratory under contract DE-AC52-07NA27344. Lawrence Livermore National Security, LLC

This work builds on decades of research by an incredible team across LLNL and the wider community!

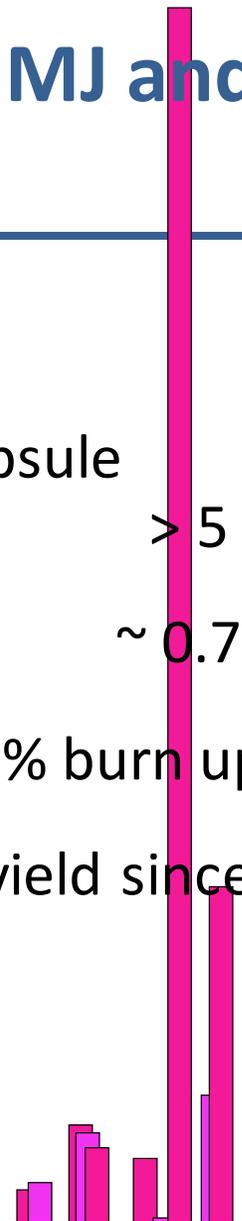


...and many more

The August 8th NIF shot (N210808) yielded more than 1.3 MJ and marks a significant advance in ICF research



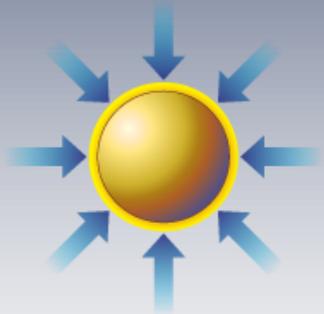
- Capsule gain (yield/capsule absorbed energy) > 5
- Laser gain ~ 0.7
- Burn propagation ~ 2% burn up
- 13x increase in fusion yield since APS 2020



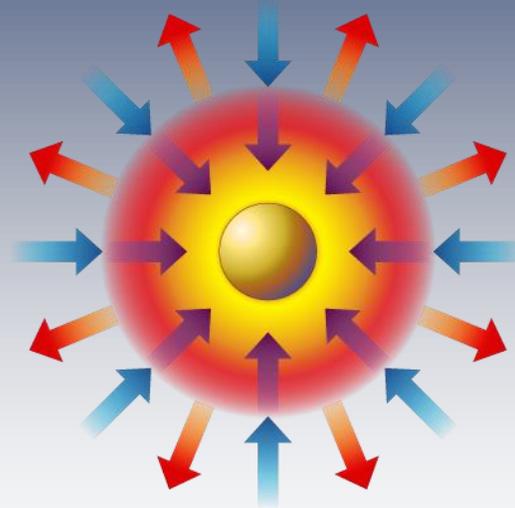
Our field is in a place that we have never been in before!

Inertial Confinement Fusion (ICF) can be achieved by using high power lasers to drive a spherical implosion

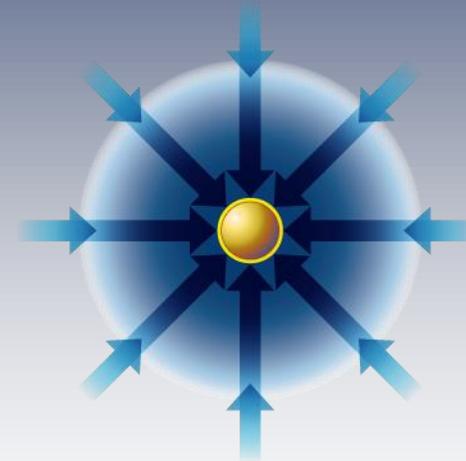
Lasers produced X-rays rapidly heat the surface of a capsule containing deuterium-tritium (DT) fuel



The blowoff plasma accelerates the DT fuel inwards in a rocket-like reaction



The fuel stagnates creating a hot central core, surrounded by a dense confining shell



The core ignites and fusion burn propagates into the dense shell, yielding many times the input energy



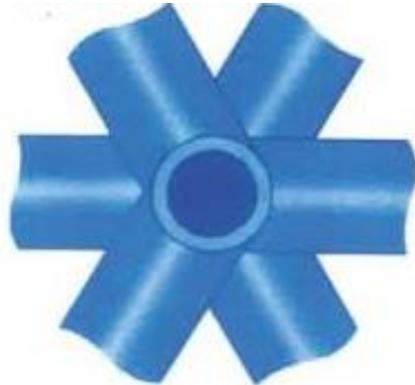
-  Radiation
-  Blowoff
-  Inward Accelerated Shell

The U.S. is pursuing several complementary approaches to ignition in ICF research

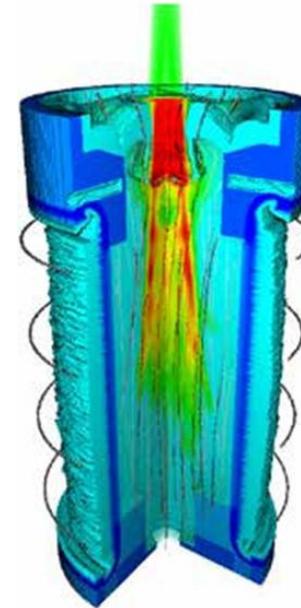
Laser indirect drive



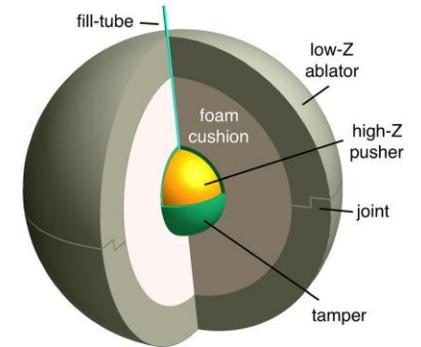
Laser Direct Drive



Magnetic Drive



Double shells



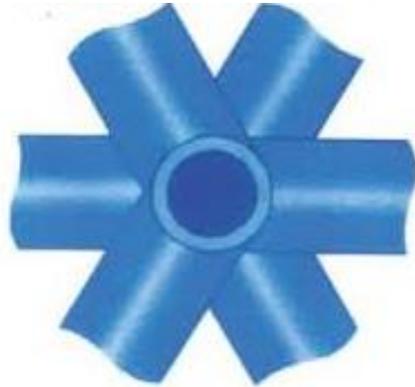
Achieving ignition in the laboratory is a Scientific Grand Challenge nearly 60 years in the making

This talk will focus on laser indirect drive

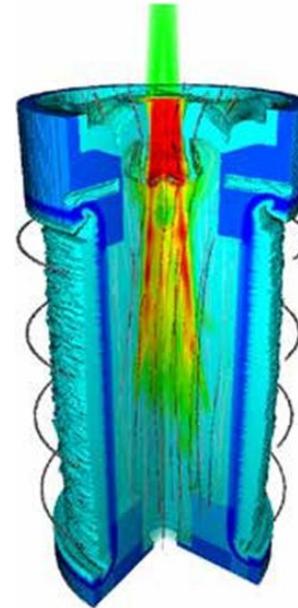
Laser indirect drive



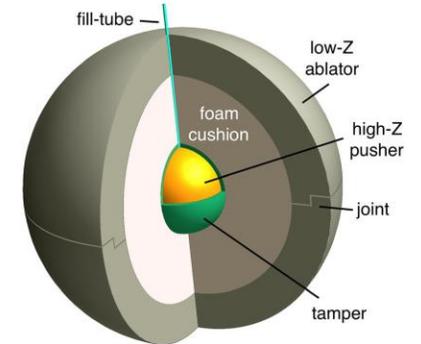
Laser Direct Drive



Magnetic Drive



Double shells



Invited sessions: BI01, KI02, QI02, WI02, ZI02

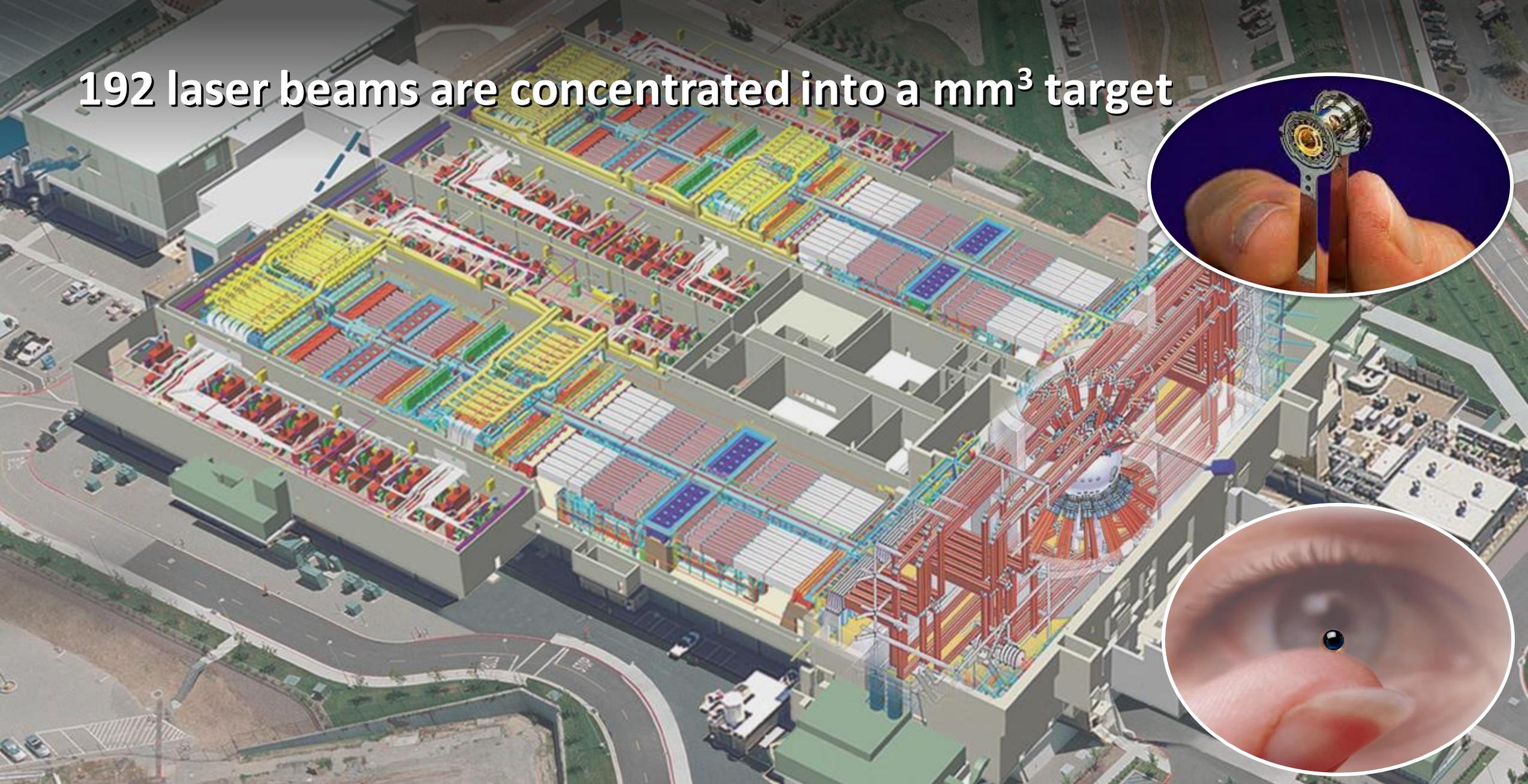
Achieving ignition in the laboratory is a Scientific Grand Challenge nearly 60 years in the making

NIF is the world's most energetic laser enabling the study of extreme conditions for Stockpile Stewardship



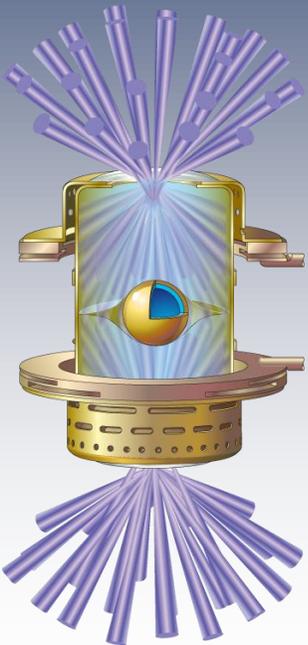
- 192 Beams, 1.9 MJ Energy, 500 TW Power
- Matter temperature $>10^8$ K
- Radiation temperature $>3.5 \times 10^6$ K
- Densities $>10^2$ g/cm³
- Pressures $>10^{11}$ atm
- Number of Diagnostics >120

192 laser beams are concentrated into a mm³ target

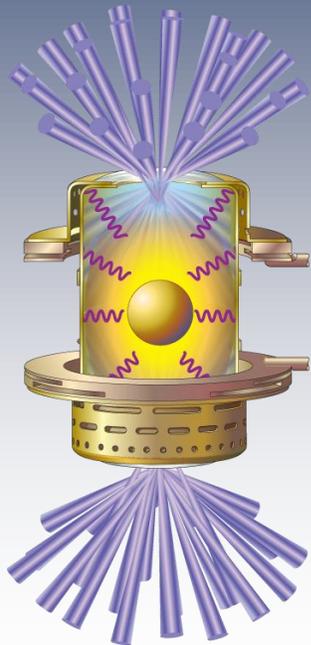


Indirect drive uses a laser driven hohlraum to compress a fuel pellet to the conditions needed for ignition

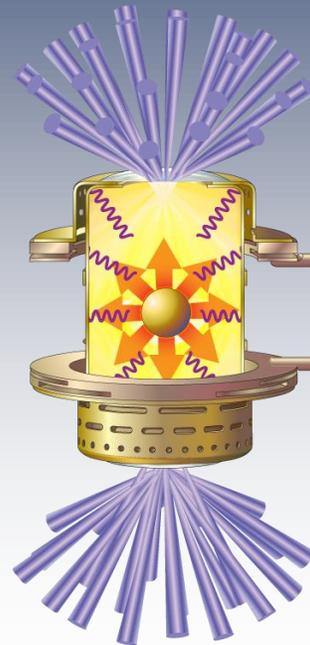
Each of the 192 laser beams are focused onto the inner wall of the hohlraum



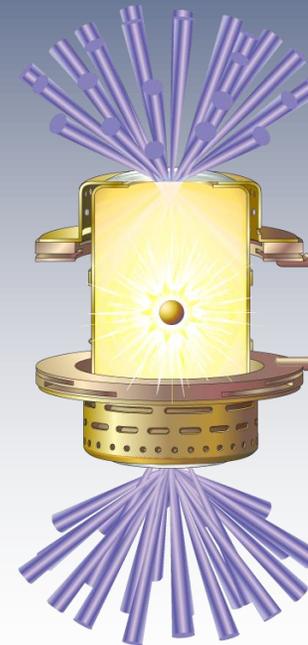
Laser beams rapidly heat the inside surface of the hohlraum creating x-rays



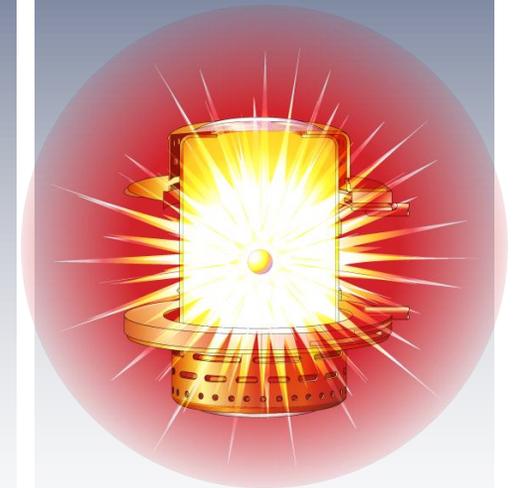
The x-rays ablate the capsule, accelerating the fuel inward to ~ 400 km/s



The fuel core reaches 500-1000 g/cc and ignites at temperatures >5 keV



Fusion burn spreads rapidly through the compressed fuel, yielding many times the input energy



Ignition (a thermal instability), the plasma must overcome all energy losses for a duration of time

Cold DT shell

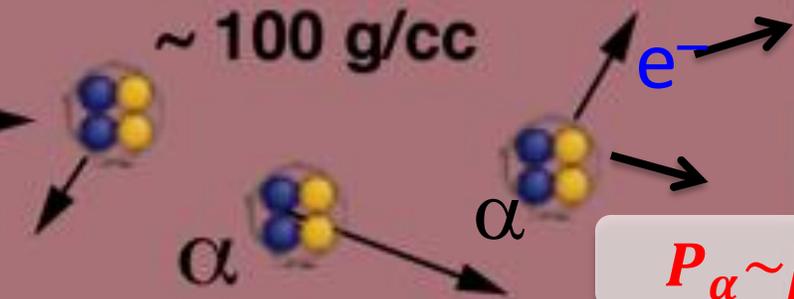
~ 1000 g/cc

Compressed DT fuel with hot central core

Fusion starts in DT hotspot

~ 50 million degrees

~ 100 g/cc



particles deposit energy
- Bootstrap heating

Brems x-ray loss

$$P_B \sim \rho \sqrt{T}$$

Spitzer thermal conduction

$$P_e \sim T^{7/2} / (\rho R^2)$$

$$P_\alpha \sim \rho T^{3.6}$$

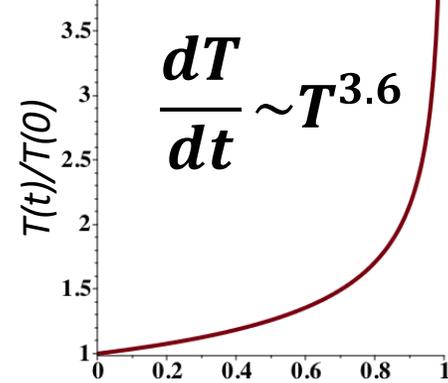
Alpha-heating

Time dependent heat balance (power/mass):

$$c_{DT} \frac{dT}{dt} = f_\alpha P_\alpha - f_B P_B - P_e - \frac{1}{m} p \frac{dV}{dt}$$

Ignition when these terms dominate

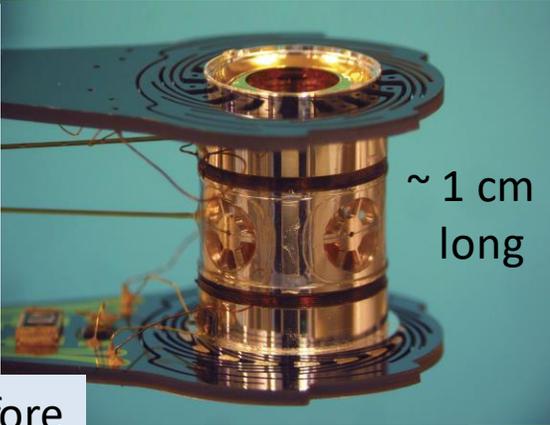
Thermonuclear instability



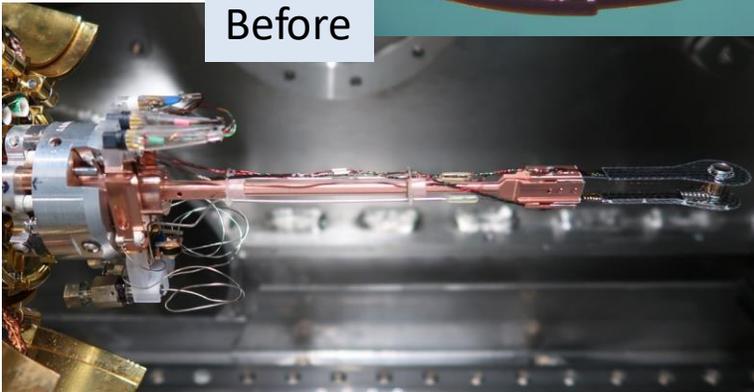
$t/\tau \quad \tau \sim 10's \text{ of picoseconds}$

We are trying to engineer a situation where heating dominates over losses

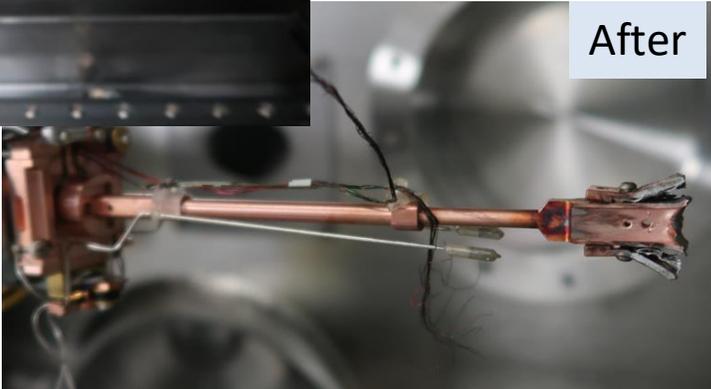
Targets are engineering marvels



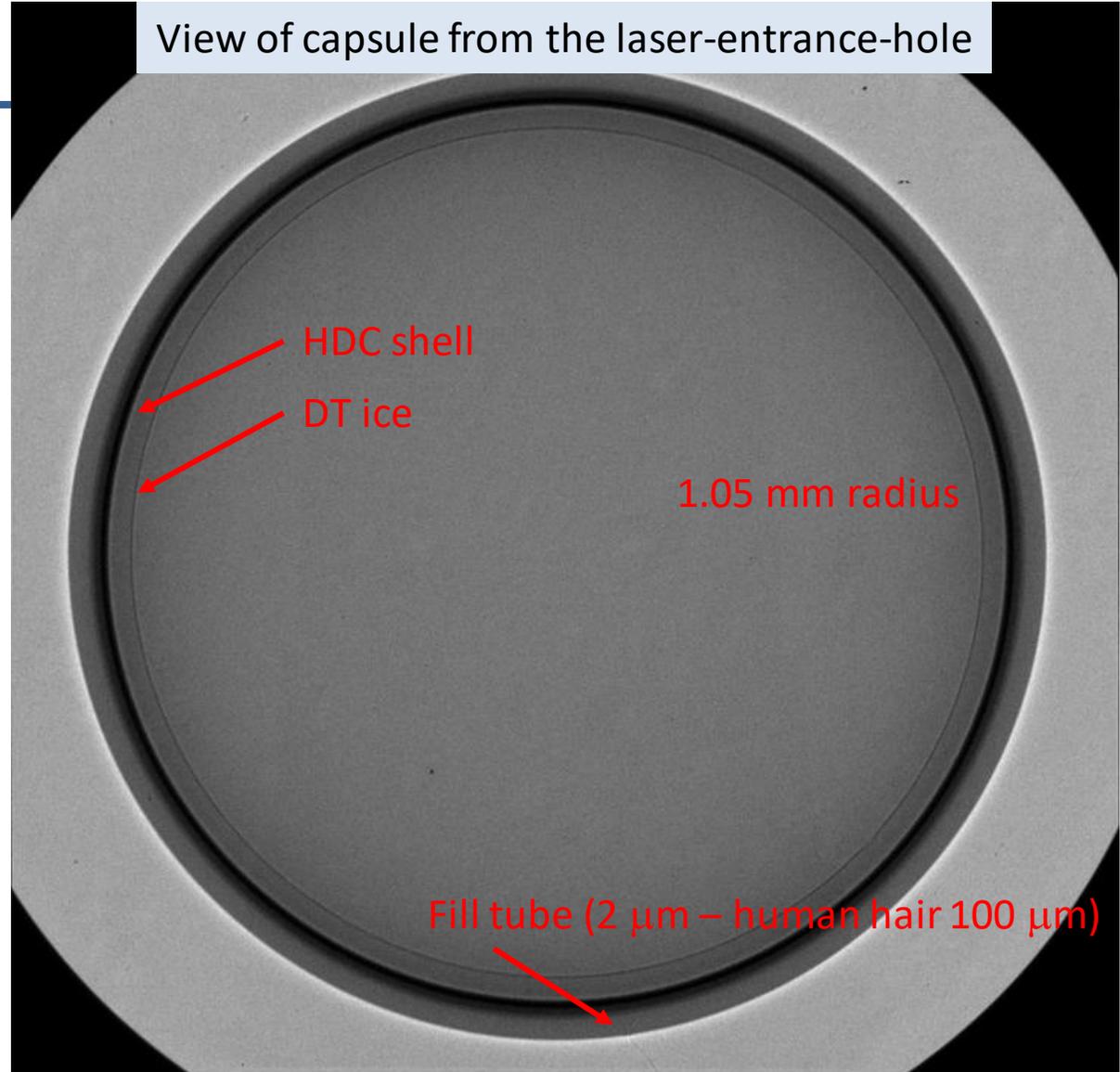
Before



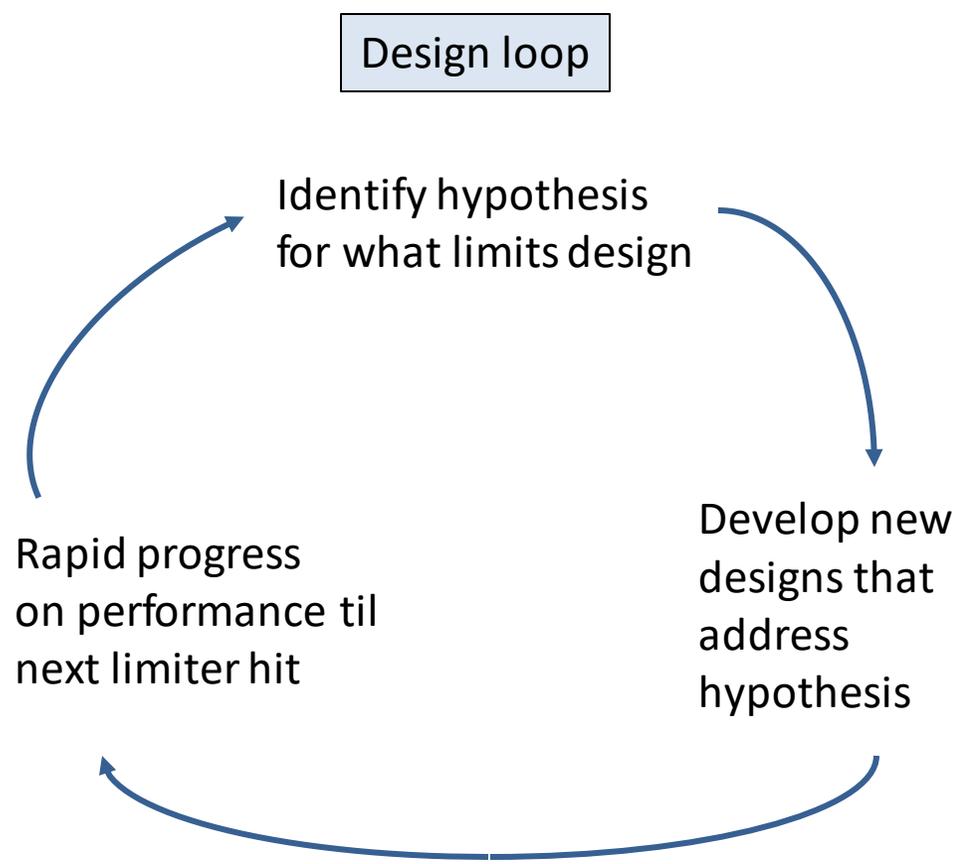
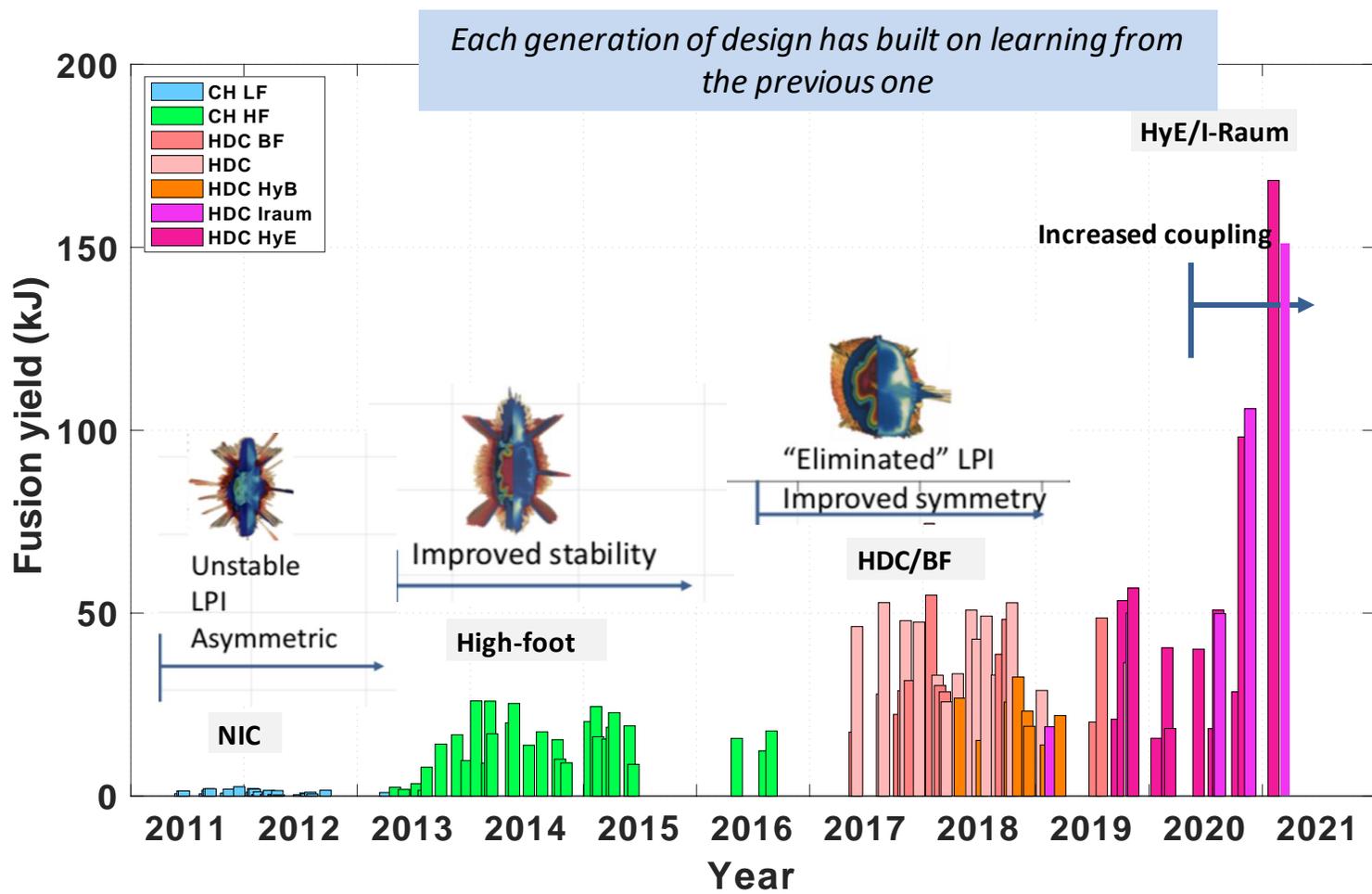
After



View of capsule from the laser-entrance-hole



We have made progress towards ignition in steps – learning what limits the implosion and the redesigning based that learning

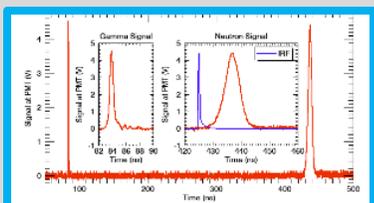


Simulations: D. Clark

NIF diagnostics have provided key insight into our experiments and built understanding, here are some examples

DT Ion temperature, hot spot velocity, fuel density, yield

- Five Neutron Time of Flight (nToF)'s and the Magnetic Recoil Spectrometer (MRS)

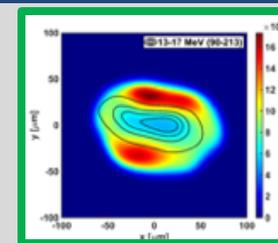


DT Neutron yield



- Zirconium/Copper Nuclear activation

Hot spot and Fuel Shape from Neutron Imagers



- 3 Neutron Imaging (NIS) Lines of sight for 3D reconstruction of neutron hot-spot
- 2 NIS down-scatter lines of sight for fuel shape

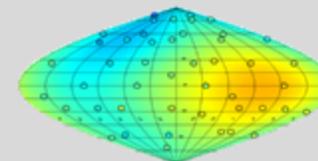
Burn width, Bang Time, DT neutron yield



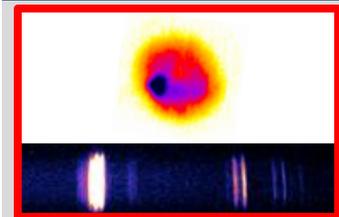
- Gamma Reaction History

DT Yield Map /Fuel uniformity

- 48 Real-Time Nuclear Activation (NAD)'s



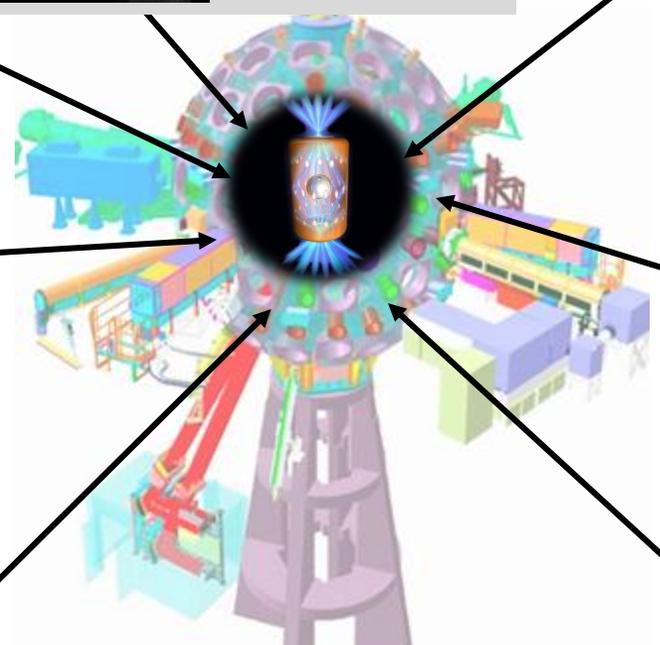
X-ray Imaging & Spectroscopy



- 3 x-ray imaging lines of sight
- X-ray spectroscopy to characterize material mixed into the hot spot

DT Fuel uniformity: Compton Radiography

- ~100keV x-rays produced by Advanced Radiography Source provide radiographs of DT fuel



See sessions GO07, ZO04, KI02

This is the best diagnosed HED plasma on the planet! -> Developed over decades by the whole HED community

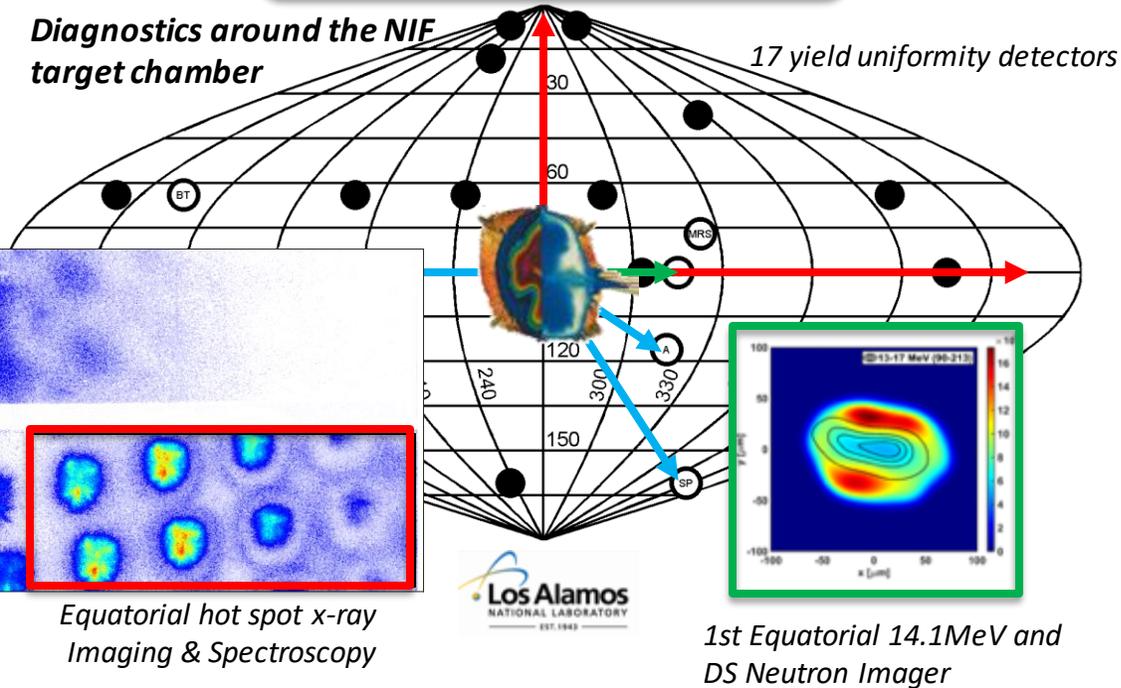
Our ability to diagnose the 3D aspects of the implosion has increased significantly from 2016 to today

2016

Polar hot spot x-ray Imaging
+ limited Spectroscopy

Diagnostics around the NIF target chamber

17 yield uniformity detectors



Equatorial hot spot x-ray Imaging & Spectroscopy

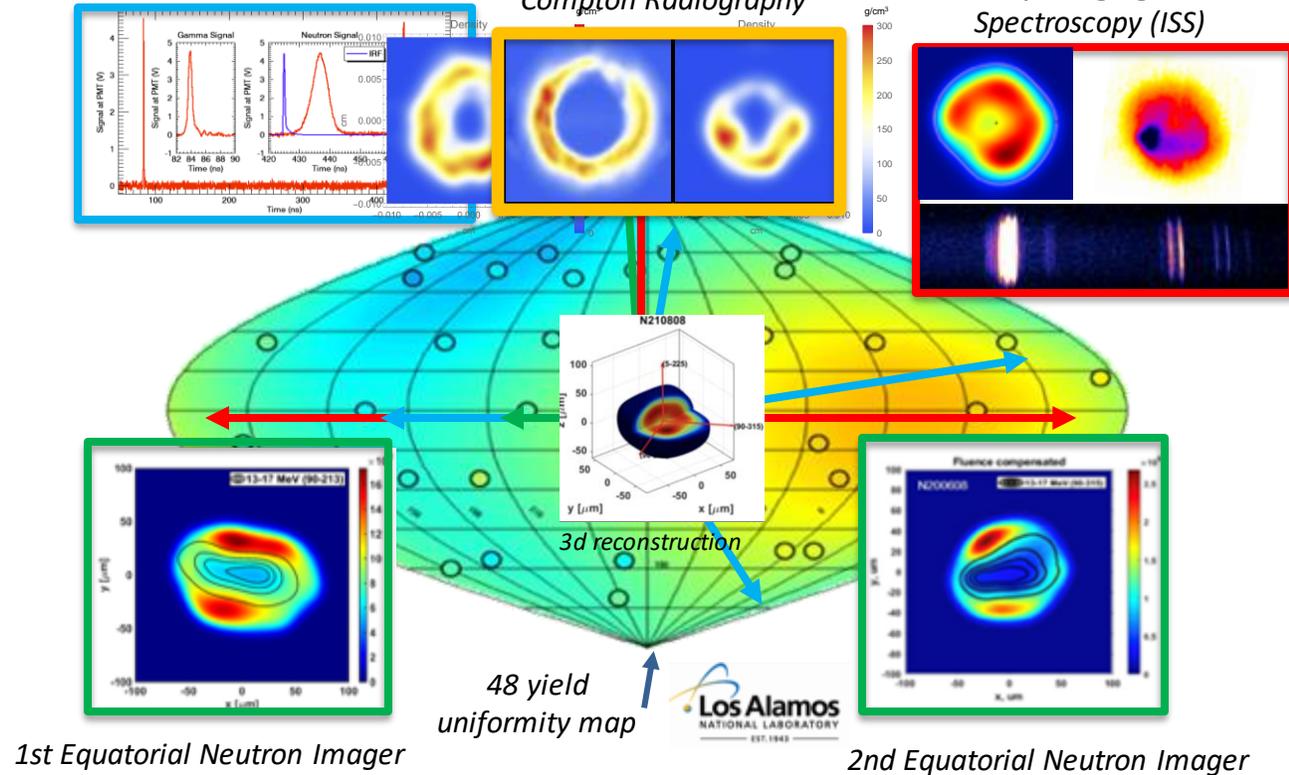
1st Equatorial 14.1MeV and DS Neutron Imager

Today

5 Quartz Cherenkov nToF lines of sight

Compton Radiography

Polar Neutron, x-ray Imaging & Spectroscopy (ISS)



1st Equatorial Neutron Imager

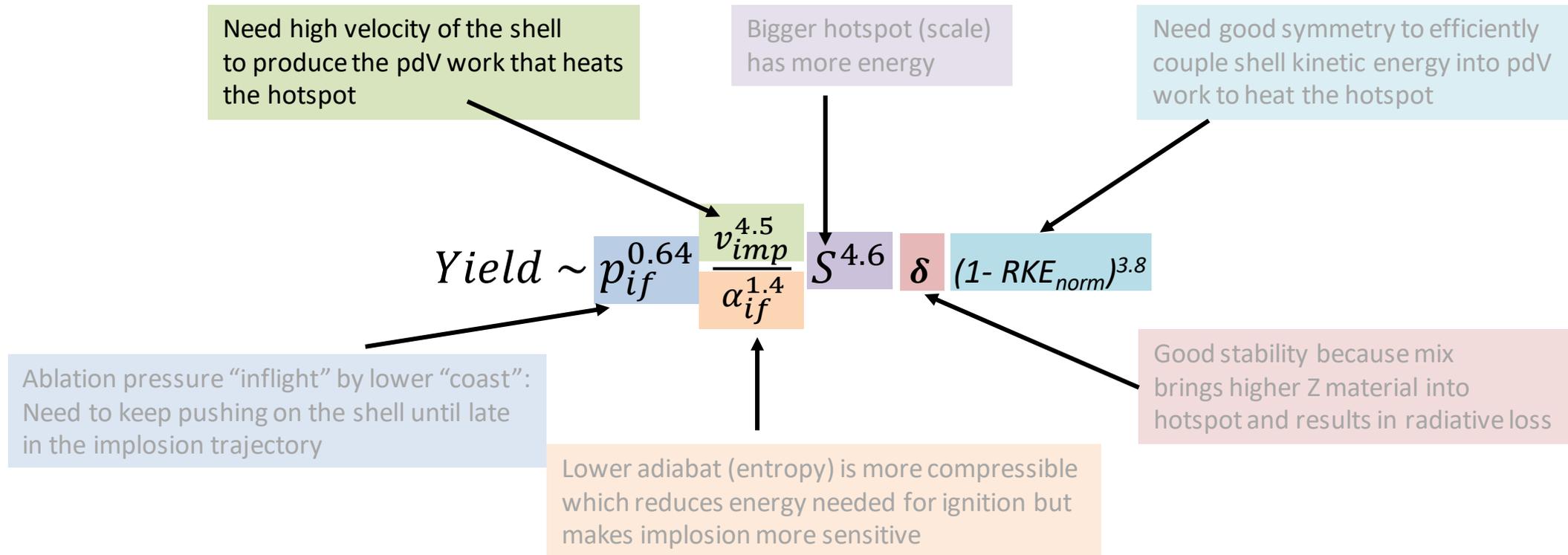
48 yield uniformity map

2nd Equatorial Neutron Imager

Improved diagnostics, theory, and simulations are key to developing our understanding

Neutron yield is a function of six parameters that we need to control to get ignition

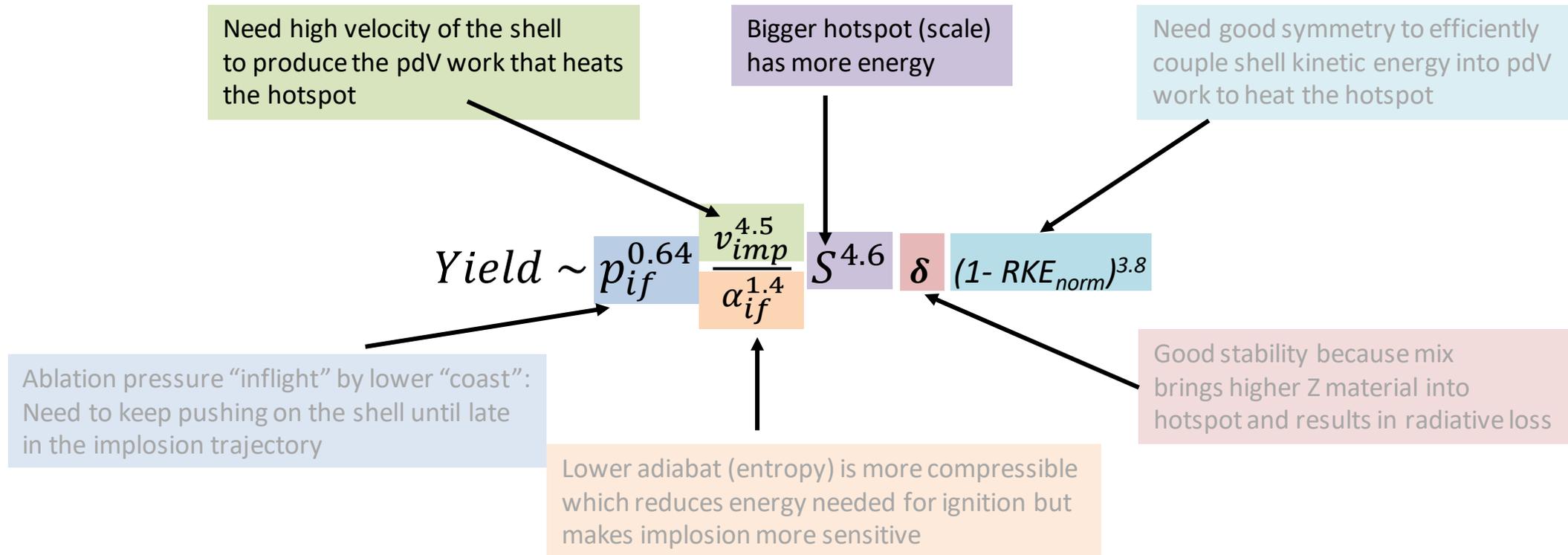
Scaling law for yield amplification ~ 2 from theory, which guides our design choices but needs to be tested experimentally



1: O. Hurricane et al, APS-DPP, PO7.00001 (2017); PPCF 61, 014033 (2019); PoP 26, 052704 (2019)

Neutron yield is a function of six parameters that we must control to get ignition

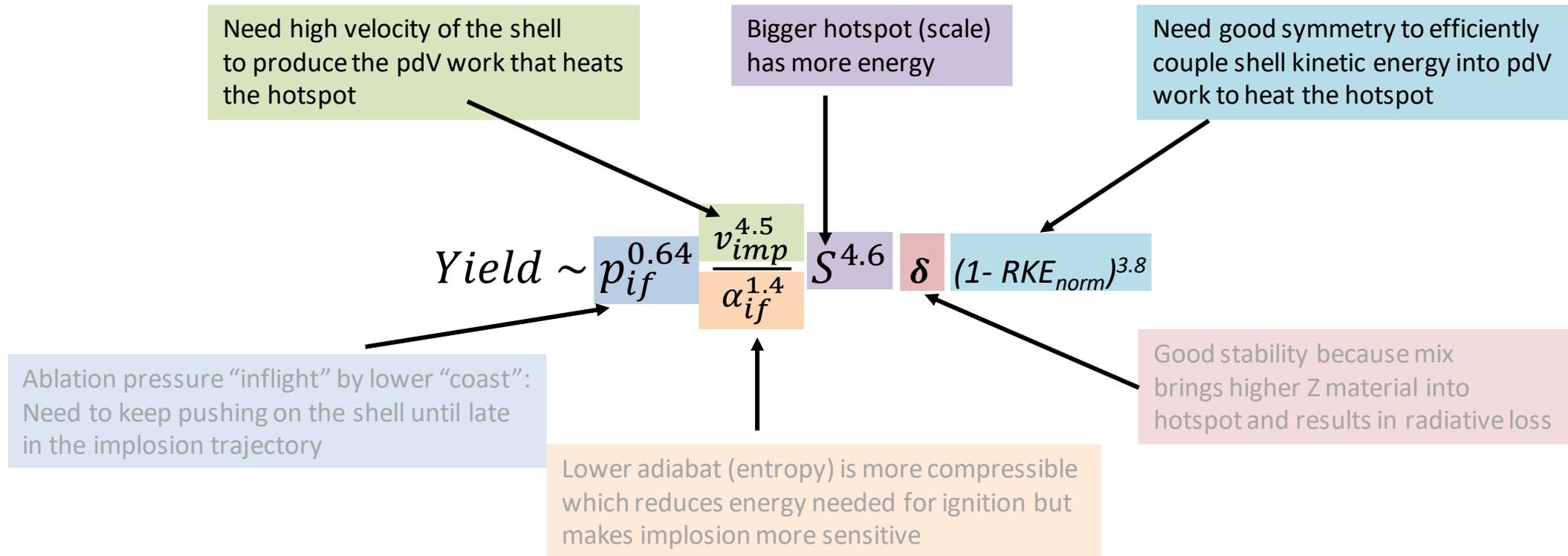
Scaling law for yield amplification ~ 2 from theory, which guides our design choices but needs to be tested experimentally



1: O. Hurricane et al, APS-DPP, PO7.00001 (2017); PPCF 61, 014033 (2019); PoP 26, 052704 (2019)

Neutron yield is a function of six parameters that we must control to get ignition

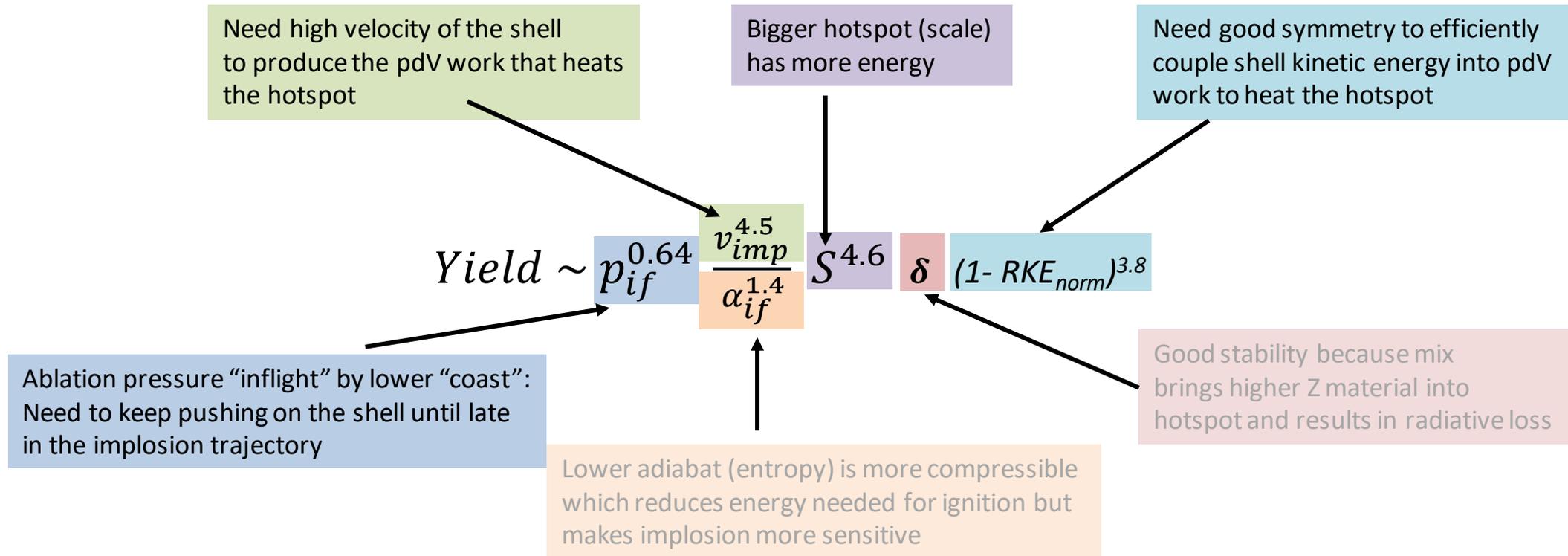
Scaling law for yield amplification ~ 2 from theory, which guides our design choices but needs to be tested experimentally



1: O. Hurricane et al, APS-DPP, PO7.00001 (2017); PPCF 61, 014033 (2019); PoP 26, 052704 (2019)

Neutron yield is a function of six parameters that we must control to get ignition

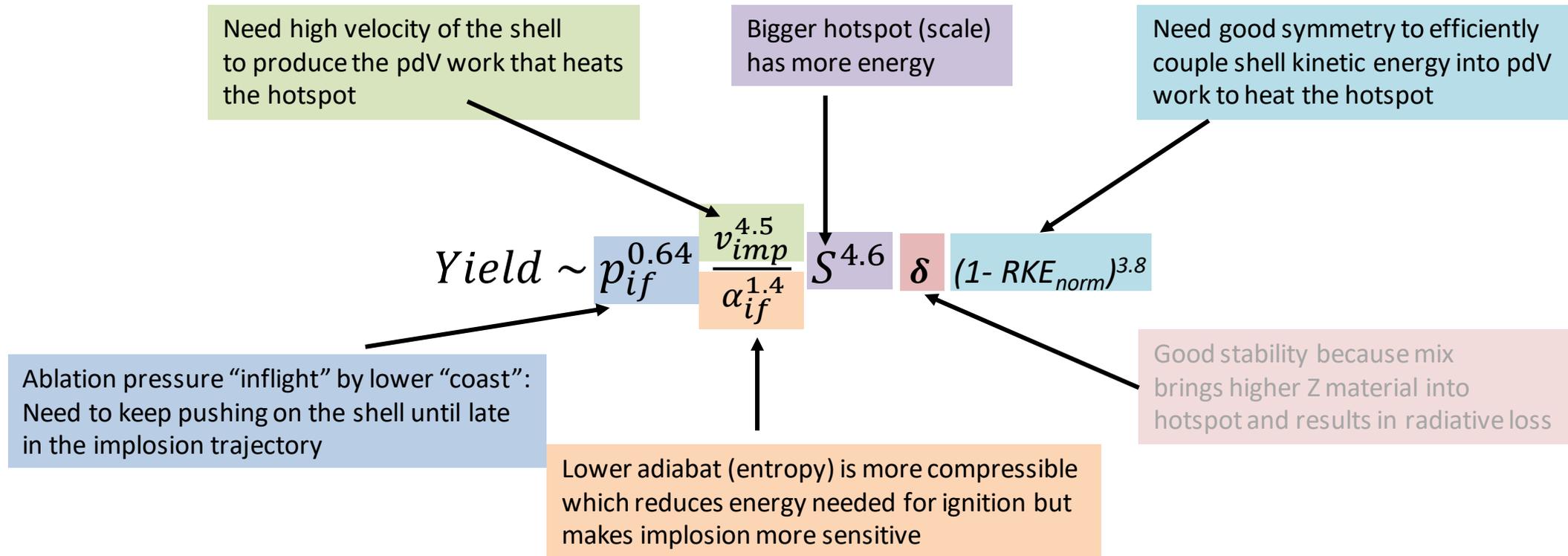
Scaling law for yield amplification ~ 2 from theory, which guides our design choices but needs to be tested experimentally



1: O. Hurricane et al, APS-DPP, PO7.00001 (2017); PPCF 61, 014033 (2019); PoP 26, 052704 (2019)

Neutron yield is a function of six parameters that we must control to get ignition

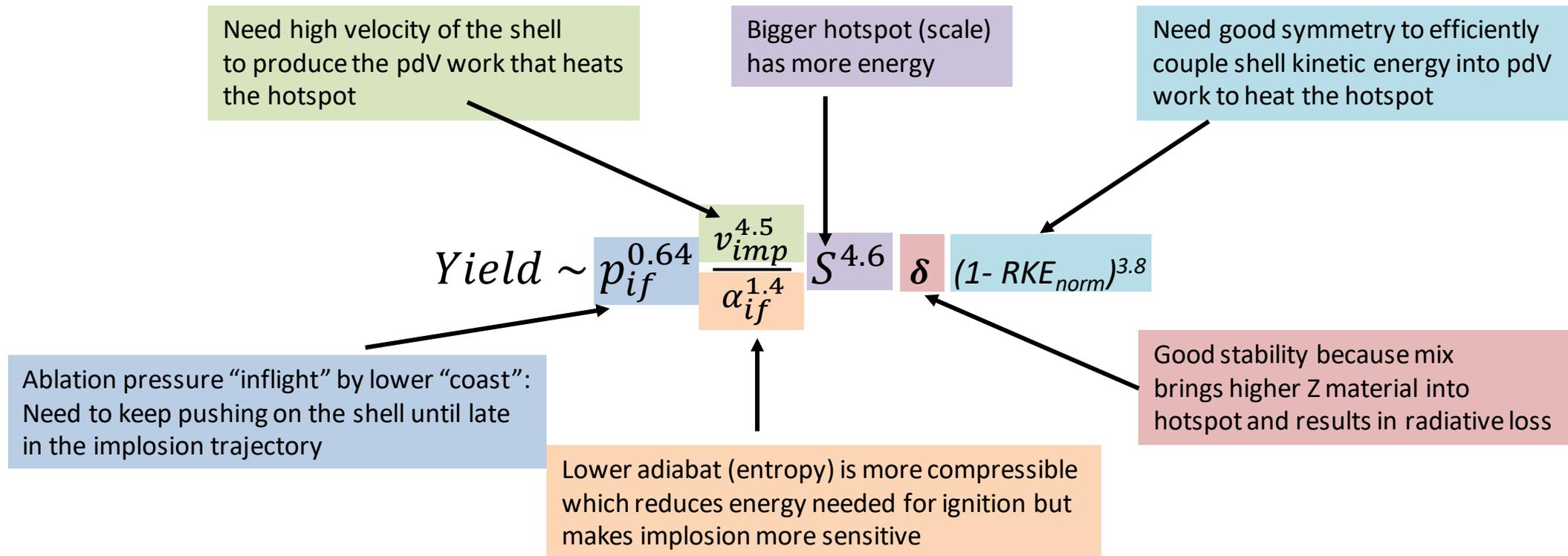
Scaling law for yield amplification ~ 2 from theory, which guides our design choices but needs to be tested experimentally



1: O. Hurricane et al, APS-DPP, PO7.00001 (2017); PPCF 61, 014033 (2019); PoP 26, 052704 (2019)

Neutron yield is a function of six parameters that we must control to get ignition

Scaling law for yield amplification ~ 2 from theory, which guides our design choices but needs to be tested experimentally

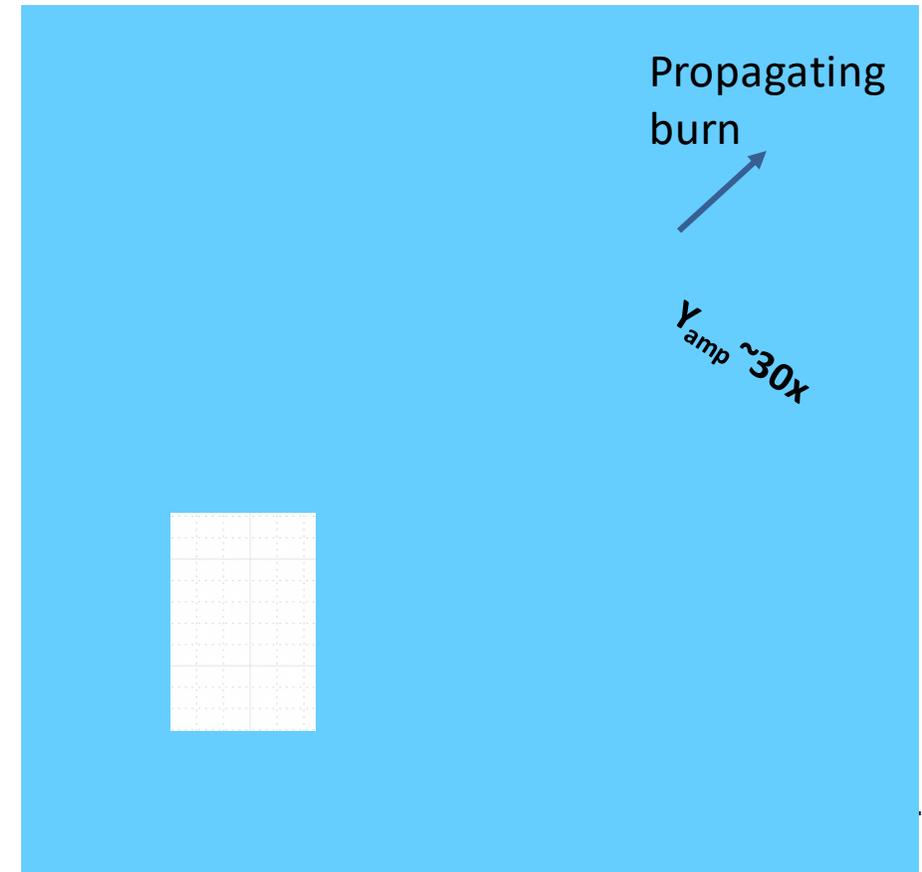


1: O. Hurricane et al, APS-DPP, PO7.00001 (2017); PPCF 61, 014033 (2019); PoP 26, 052704 (2019)

We have steadily advanced our physics understanding and the technology over the last decade to improve performance

- $E_{HS}P_{HS}^2$ [Patel] related to ITFX and Generalized Lawson Criteria [Betti]
 - This metric uses “no burn” quantities
 - A model is used to convert data from burn on to burn off
 - “Dudded fuel” implosions check that model (more of these expts upcoming)
 - Simulations indicate that ignition corresponds to $Y_{amp} \sim 15\text{-}30x$
 - Boundary is uncertain (based on simulations)

Ignition figure of merit $\sim (\rho R)^3 T^3 \sim E_{HS} P_{HS}^2$



Walk through the various designs from the last several years

P. Patel, Phys. Plasmas 27, 050901 (2020)

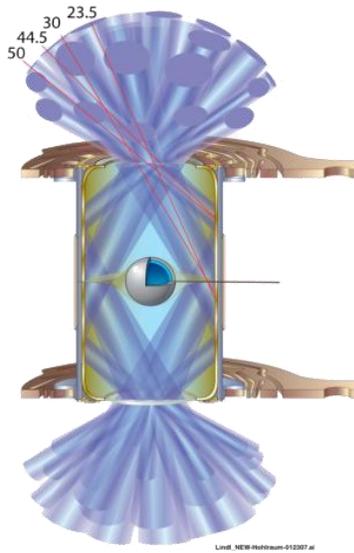
R. Betti, et al, Phys. Plasmas, 17, 058102 (2010)

A. R. Christopherson, et al, PHYSICAL REVIEW E 99, 021201(R) (2019) 4:

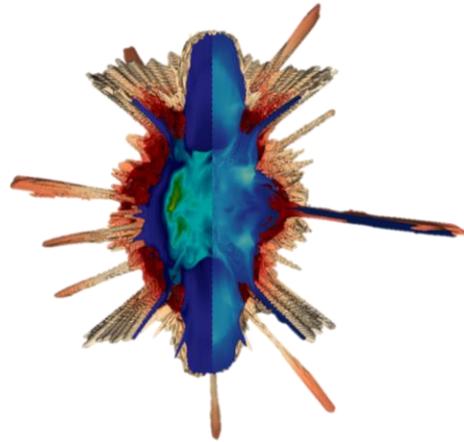
J. D. Lindl, et al, Phys. Plasmas, 25, 122704 (2018)

Initial designs were low adiabat, plastic capsules in high-gas filled hohlraums and had low hot spot pressure and energy

“Low foot¹” $Y = 2 \text{ kJ}$

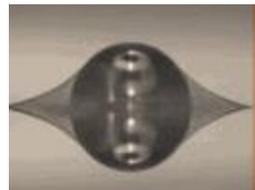


High gas filled hohlraum
Plastic capsule



Clark, et al IWPCTM, 2018

- > High LPI (reduces drive and velocity)
- > Symmetry swings
(Cross beam energy transfer)
- > Mix (tent)



Tent

Ignition figure of merit $\sim (\rho R)^3 T^3 \sim E_{HS} P_{HS}^2$

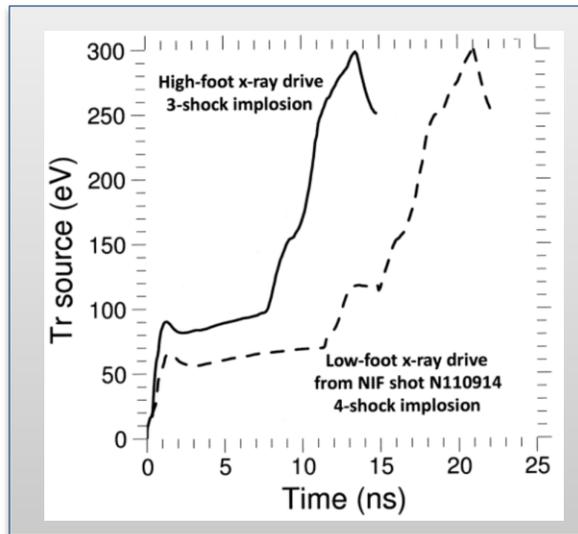


It was later found that the capsule support membrane (“tent”) was a major factor disrupting the hot spot

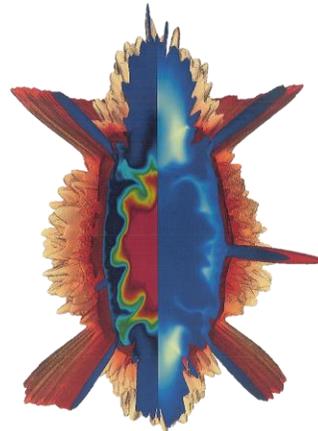
The high foot design increased the adiabat to reduce capsule instability and convergence – performance improved but plateaued

“High foot¹⁻³” $Y = 27$ kJ

Higher adiabat¹⁻³

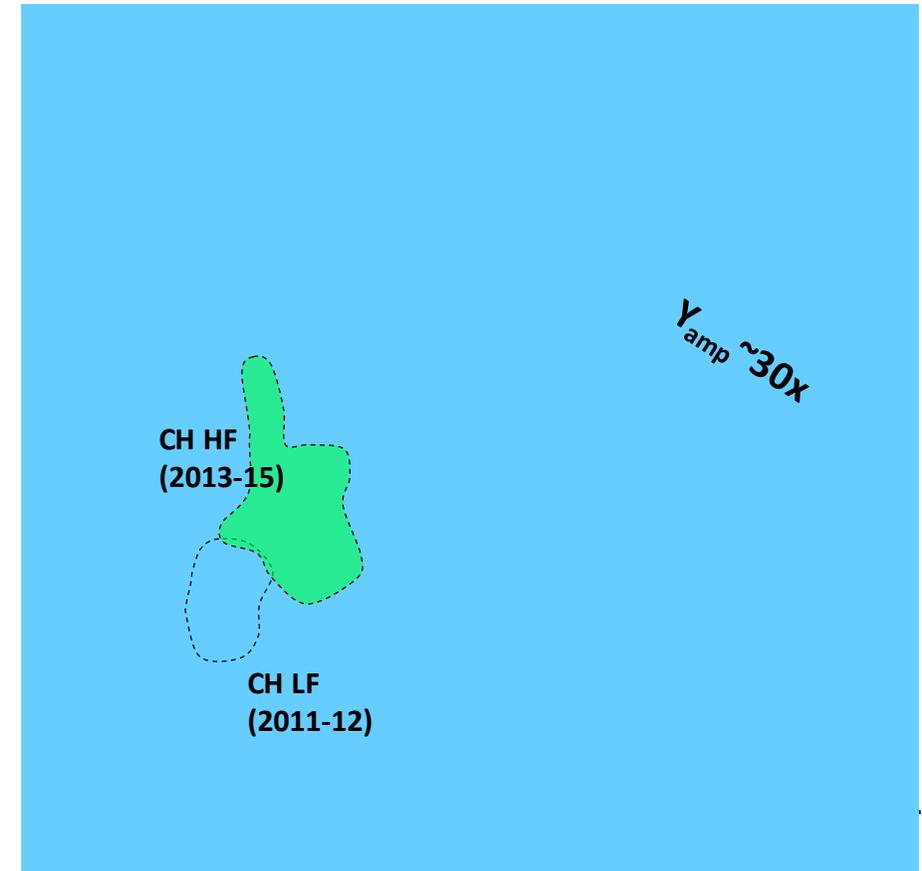


High gas filled hohlraum
Plastic capsule



Clark, et al IWPCTM, 2018

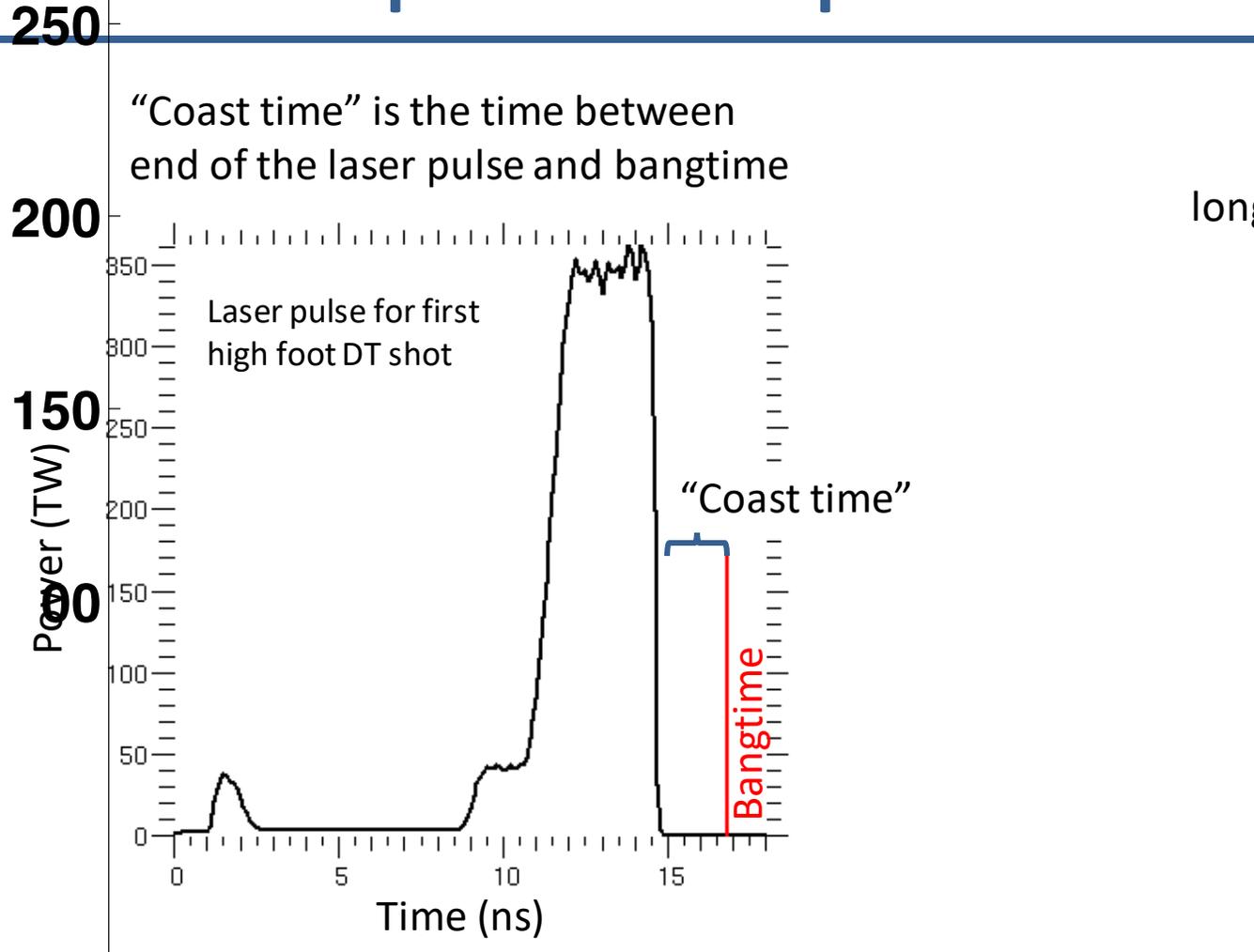
$$\text{Ignition figure of merit} \sim (\rho R)^3 T^3 \sim E_{HS} P_{HS}^2$$



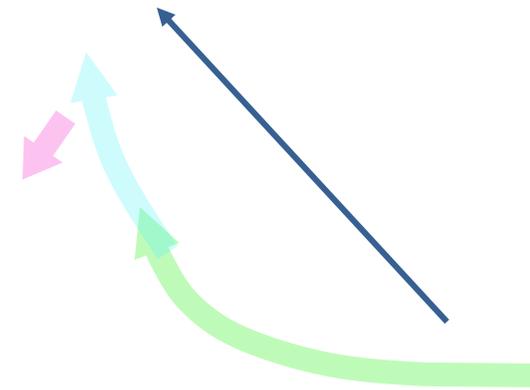
- > Improved stability
- > Higher velocity, less mix
- > LPI and shape swings persist

Tent caused the plateau. LPI remained a major problem -> energy losses, large symmetry swings

High foot showed that maintaining high ablation pressure until late in the pulse was important – control by “coast time”



Initially, we saw empirically that reducing coast time improved performance -- the longer laser pulse makes symmetry more difficult



We now have a better understanding of the impact of coast time -- see O. Hurricane’s talk in the next session

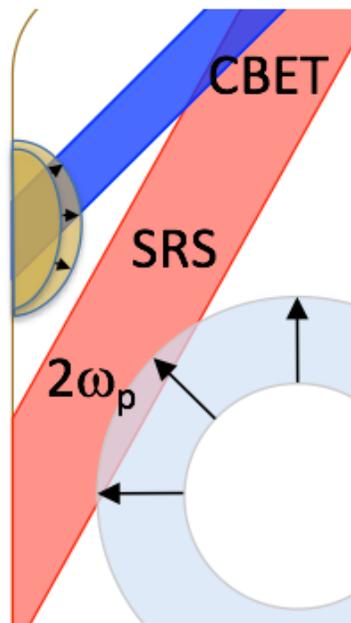
Moved to low gas-filled hohlraums to reduce laser-plasma-interactions (LPI) that had been present in high and low foot

High gas fill – LPI dominated
(~ 4-5% critical density)

Cross beam energy transfer

Stimulated Raman Scatter

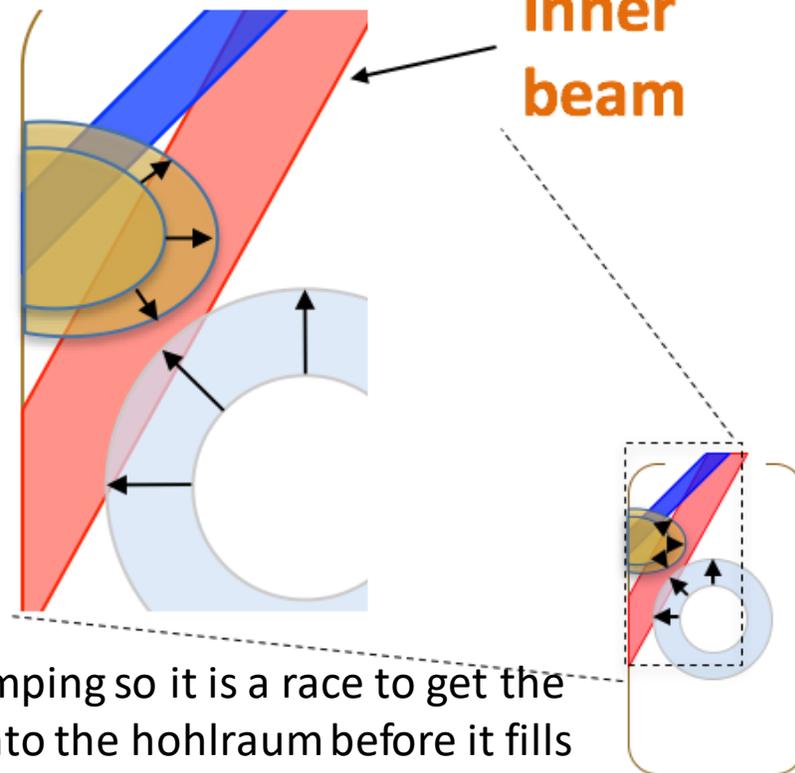
$2\omega_p$ instability



Hohlraum fill tamps the expansion of the wall and keeps the hohlraum from filling with gold plasma but has high LPI

Low gas fill – Radiation hydrodynamics dominated
(< 2.5% critical density)

Inner beam



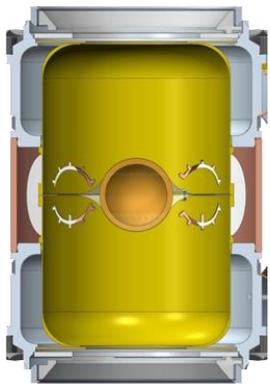
Less tamping so it is a race to get the pulse into the hohlraum before it fills with plasma

Low gasfill hohlraums are a good match to diamond (HDC) ablators because the high density leads to shorter pulses

Diamond (HDC) capsules resisted the tent and used shorter laser pulses – performance improved

“HDC¹/BigFoot²” Y = 55 kJ

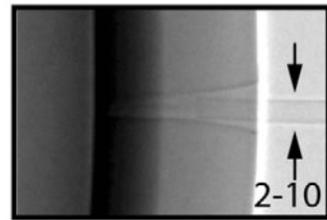
Low gas filled hohlraum



Diamond capsule



10μm -> 5μm fill tube³



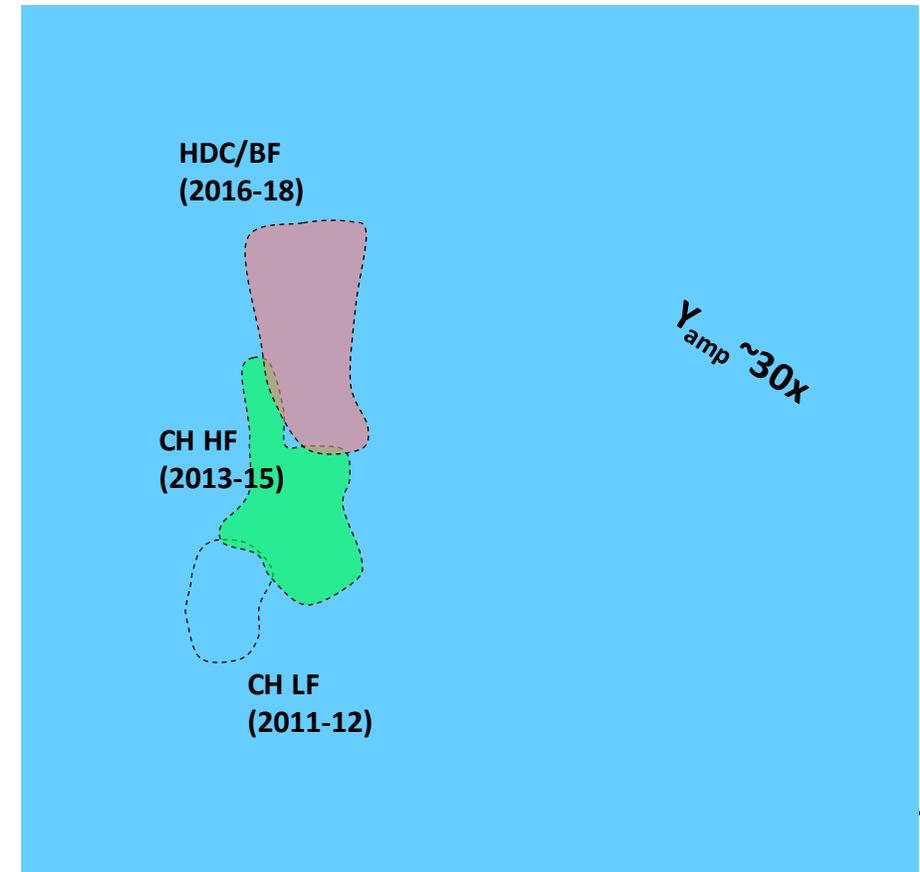
Fill tube

- > Reduced LPI
- > Better symmetry

- > Tent resistance
- > Higher velocity

- > Reduced mix
- > Less radiative loss

Ignition figure of merit $\sim (\rho R)^3 T^3 \sim E_{HS} P_{HS}^2$



Energy still too low – need a bigger capsule, same implosion pressure, but no more laser energy – major challenge

HYBRID-E Challenge: make capsule bigger! But keep similar adiabat, stability, velocity, “coast time”, and *symmetry* with fixed laser energy

Much bigger capsule in to slightly larger hohlraum



HDC (BigFoot)

Lead designer: L. Berzak Hopkins, C. Thomas
Lead expt: S. Le Pape, D. Casey



HYBRID-E²

Lead designer: A. Kritcher
Lead expt: A. Zylstra

$$Y \sim p_{if}^{0.64} \frac{v_{imp}^{4.5}}{\alpha_{if}^{1.4}} S^{4.6} \delta (1 - RKE_{norm})^{3.8}$$

Yield

Ablation pressure “inflight” by lower “coast”

Implosion velocity

Scale

Symmetry

Adiabat

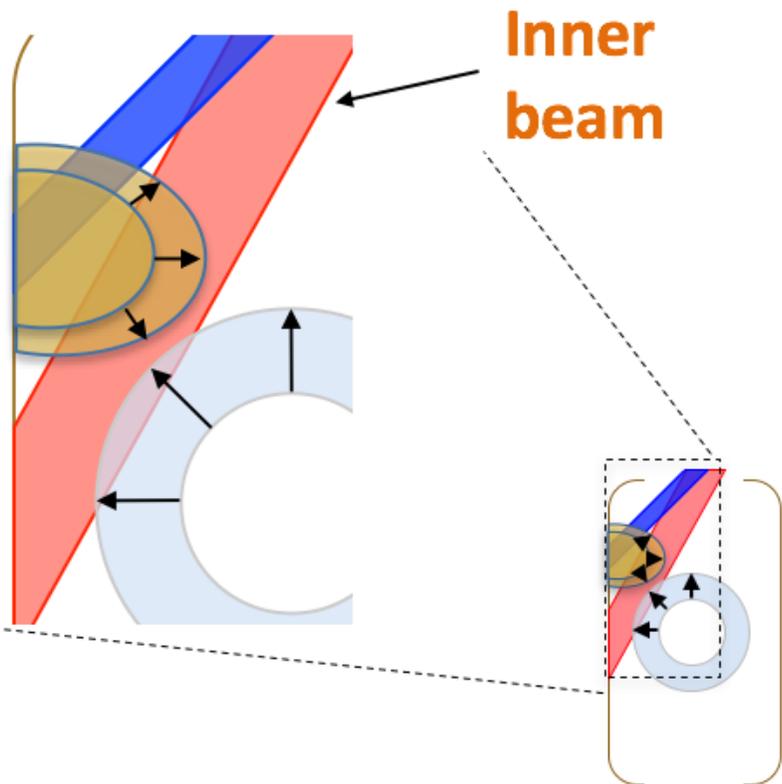
Stability

High Yield Big Radius Implosion Design (HYBRID) strategy¹

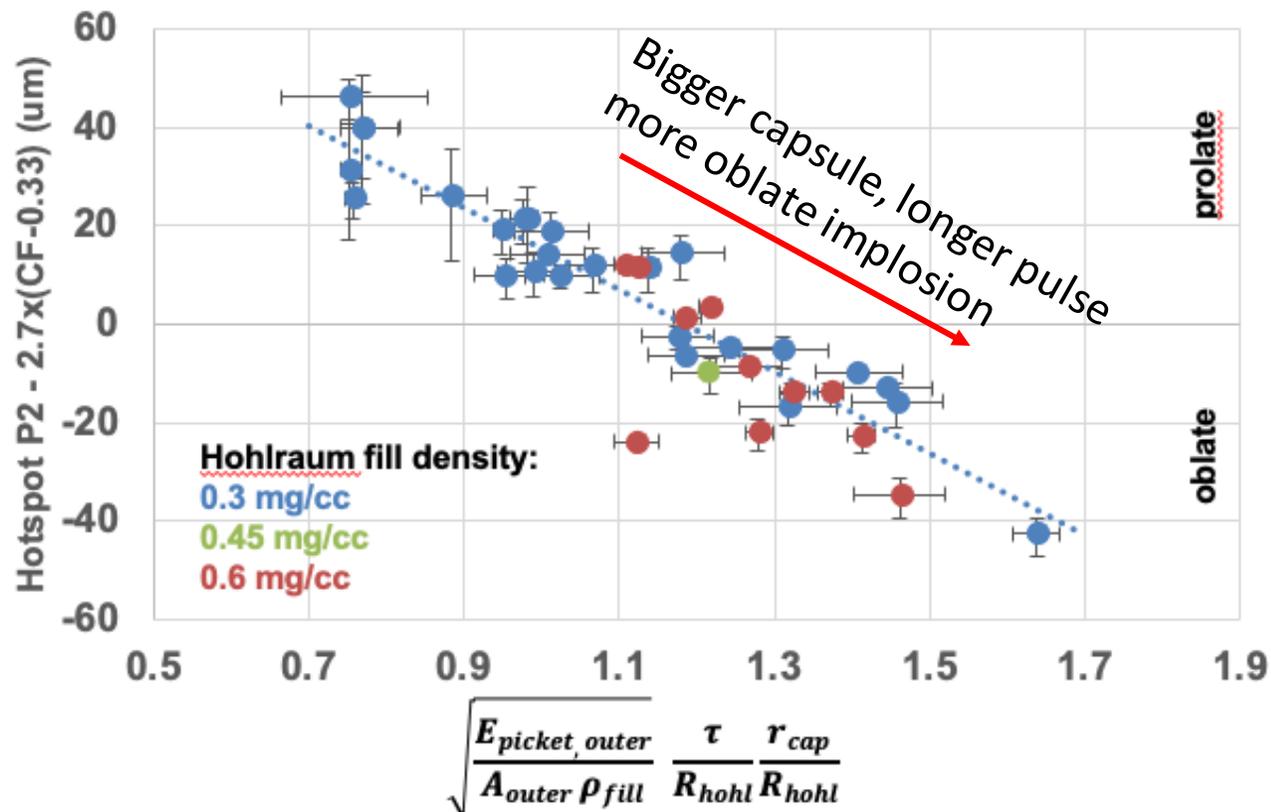
1: O. Hurricane et al, APS-DPP, PO7.00001 (2017); PPCF 61, 014033 (2019); PoP 26, 052704 (2019)
2: A.B. Zylstra et al., PRL 126, 025001 (2021); A.L. Kritcher et al., PoP 28, 072706 (2021)
3: D.A. Callahan et al., PoP 25, 056305 (2018); J. Ralph, et al., PoP, 25, 082701 (2018)
4: A. L. Kritcher, et al Phys. Rev. E 98, 053206 (2018), L. Pickworth, et al, PoP (2020)

Need to maintain a symmetric implosion with the larger capsule and low coast to effectively couple kinetic energy to the hotspot

Hohlraum symmetry dominated by the inner beams being stopped by the gold “bubble” and expanding ablator



Identified the parameters important to asymmetry in 2017/2018

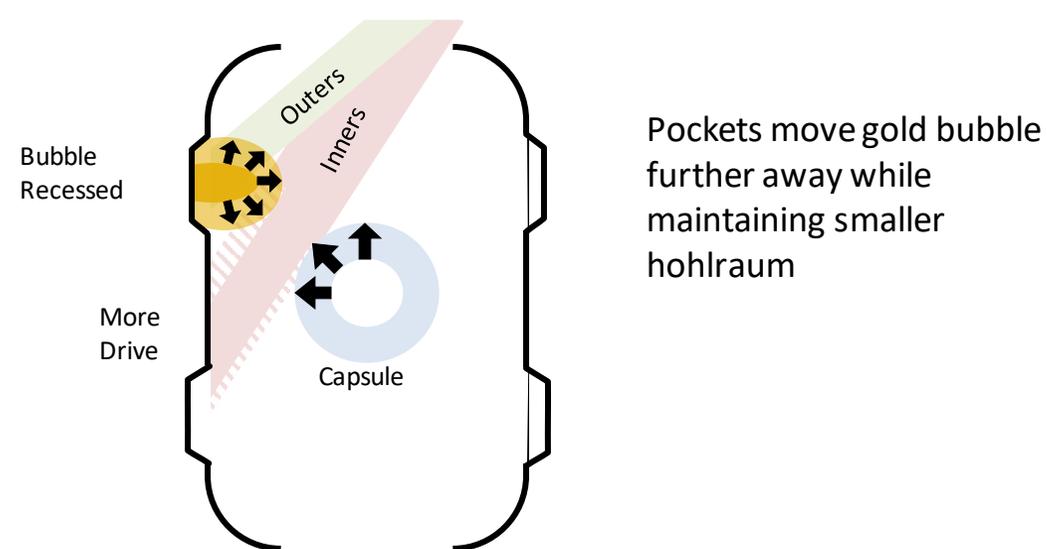
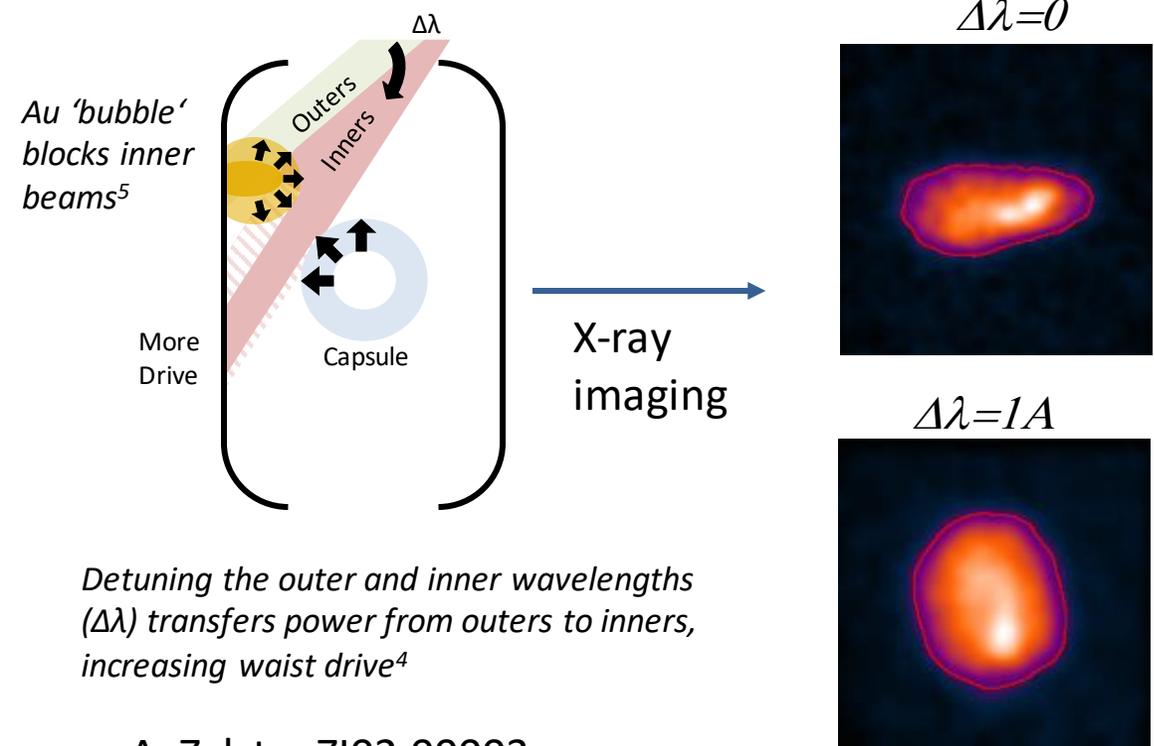


Needed some additional symmetry techniques to drive a round implosion with the larger capsule and short coast time

Developed two additional techniques for symmetry control: cross beam energy transfer and “Iraum” geometry

Hybrid E
Cylindrical hohlraum with cross beam energy transfer

“Iraum”
Shaped hohlraum with cross beam energy transfer



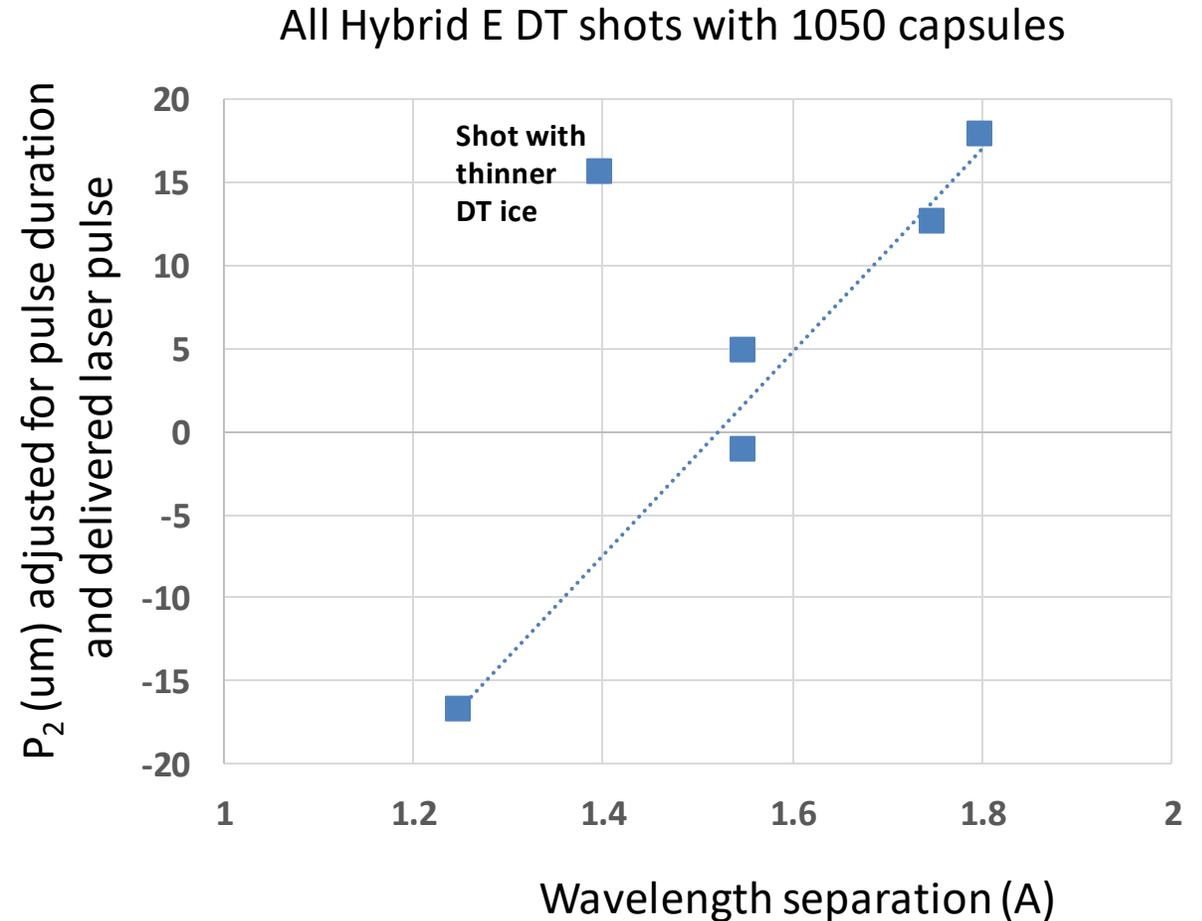
S. Ross BI01.00002
C. Young ZI02.00001

A. Zylstra ZI02.00003
A. Kritcher GO04.00002

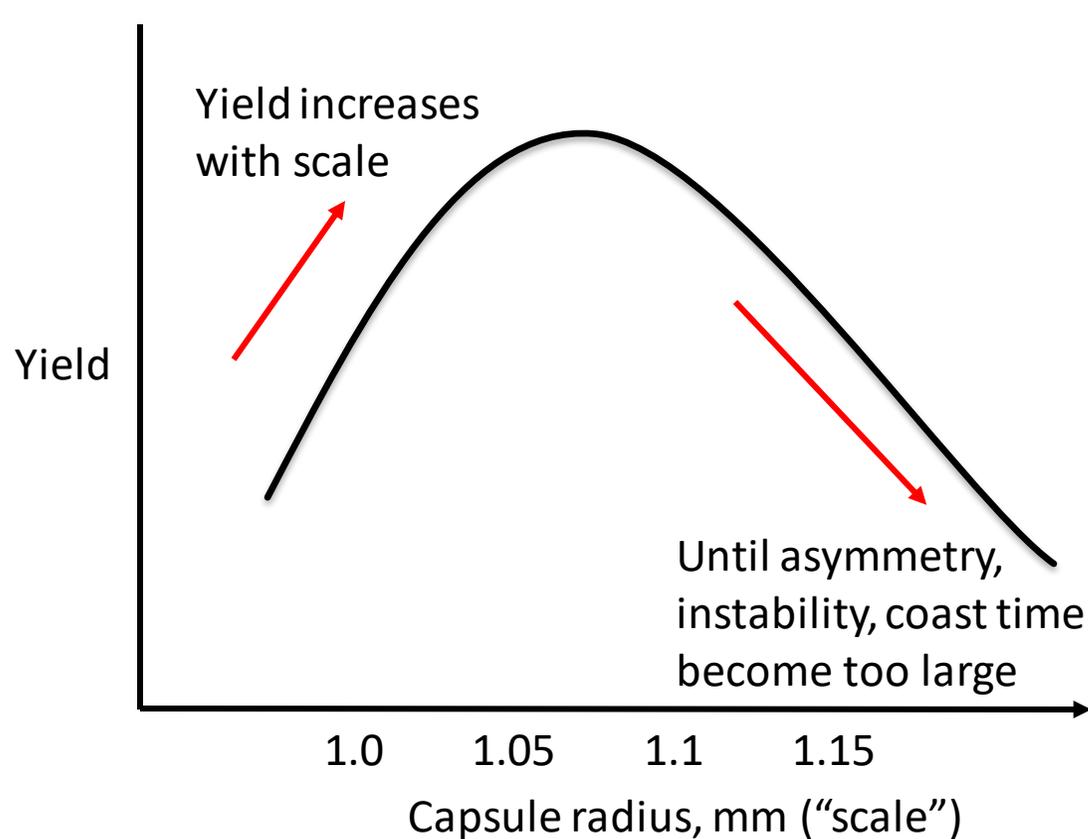
Both designs have produced a “burning plasma” in the last year!

Cross beam energy transfer (CBET) is a powerful tool for symmetry control in low gasfill hohlraums

- Low fill hohlraums -- smaller wavelength separation than high gasfill (1-2 Å vs 8-9 Å)
 - High gasfill hohlraums -- very large transfer in the foot ($\sim 10x$). Led to swings in symmetry which were difficult to control/predict
 - Experiments show consistent scalings of P_2 with CBET but can have occasional surprises when other large changes are made (but can recover)



Because there is a tradeoff between capsule size, symmetry, and “coast time,” we expected there to be an optimum capsule size



$$Y \sim p_{if}^{0.64} \frac{v_{imp}^{4.5}}{\alpha_{if}^{1.4}} S^{4.6} \delta (1 - RKE_{norm})^{3.8}$$

Yield

Ablation pressure “inflight” by lower “coast”

Implosion velocity

Scale

Symmetry

Adiabat

Stability

Plan was to use different capsule sizes and find the optimum experimentally

Initial HYBRID-E used the largest capsule and increased hotspot energy, but pressure dropped, capsule quality degraded, mix

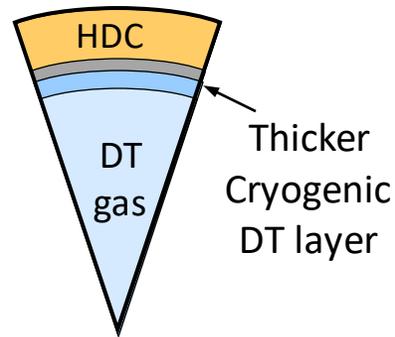
“HYBRID-E¹” 1100 μm size Y = 55 kJ

910μm → 1100μm capsule



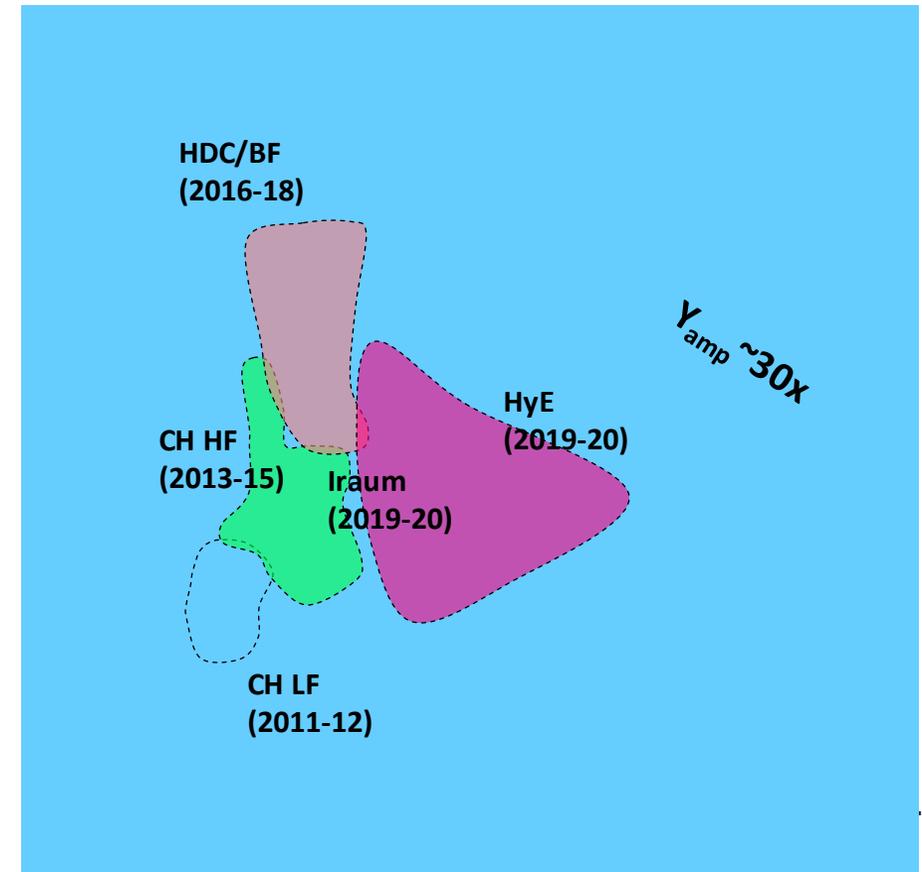
- > More coupled energy via reduced CCR (+maintain shape & velocity)
- > Longer coast
- > Capsule pits, voids
- > Mix
- > Lower pressure

Capsules and ice not hydro-scaled



- > Protection from defects
- > Better stability at low velocity but not enough

Ignition figure of merit $\sim (\rho R)^3 T^3 \sim E_{HS} P_{HS}^2$

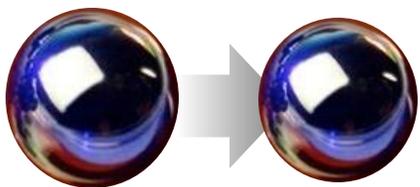


Recovering pressure (and temperature) would require lower coast and less mix

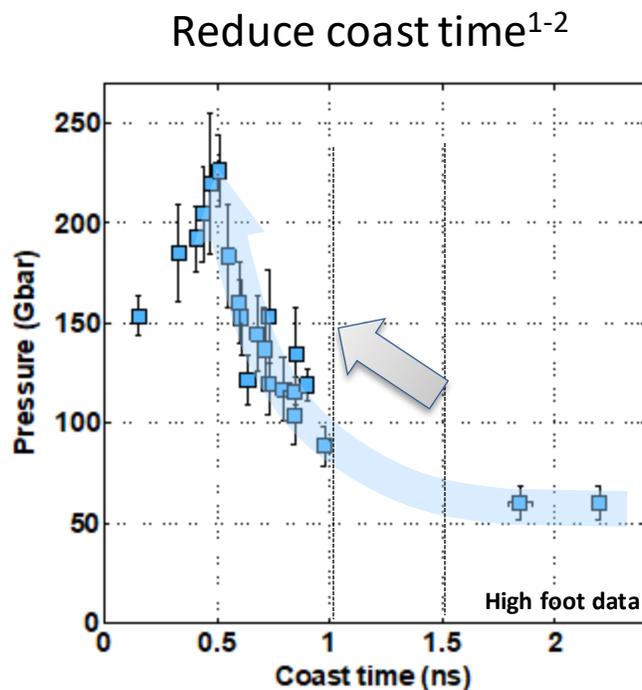
Slightly smaller capsules reduced coast time, and some capsule quality improvement – performance improved markedly

“HYBRID-E” 1050 μm size Y = 170 kJ

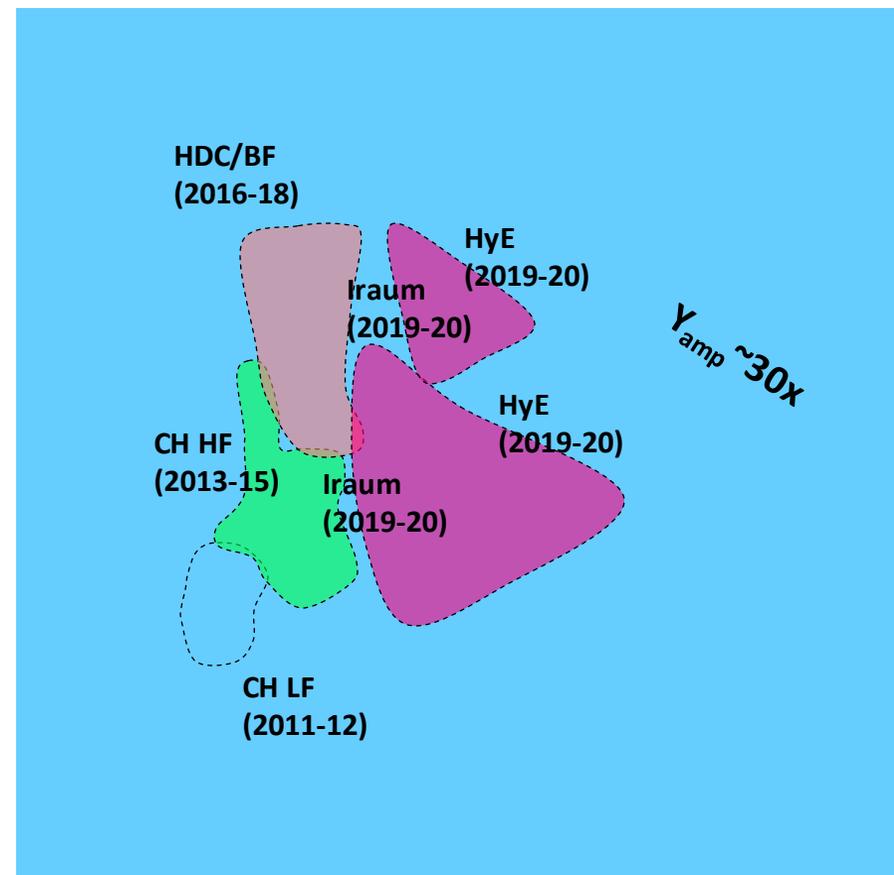
1100μm -> 1050μm capsule



- > Reduced coast
- > Higher pressure
- > Better capsules, and higher W, less mix and higher temperature



$$\text{Ignition figure of merit} \sim (\rho R)^3 T^3 \sim E_{HS} P_{HS}^2$$

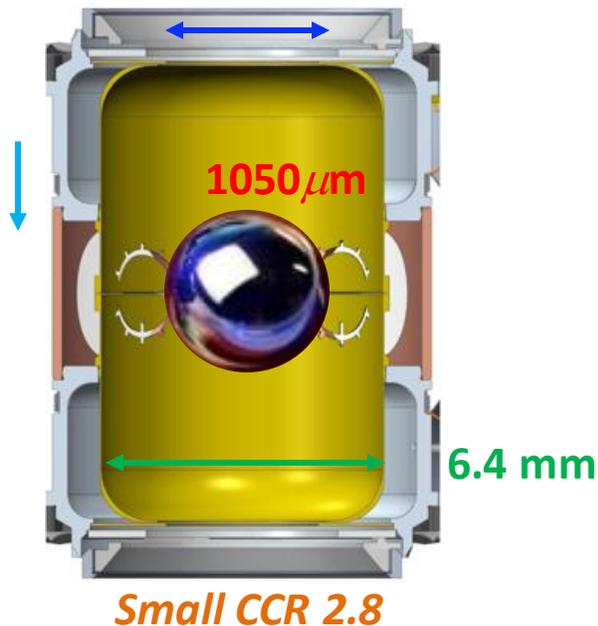


While not confirmed, analysis suggested coast was still sub-optimal, motivating test of yet lower coast

To reduce coast time further, we had to make the hohlraum even more efficient by reducing the laser-entrance-hole size

HYBRID-E design modified
with smaller LEH³

3.65 → 3.1 mm LEH (27% less area)



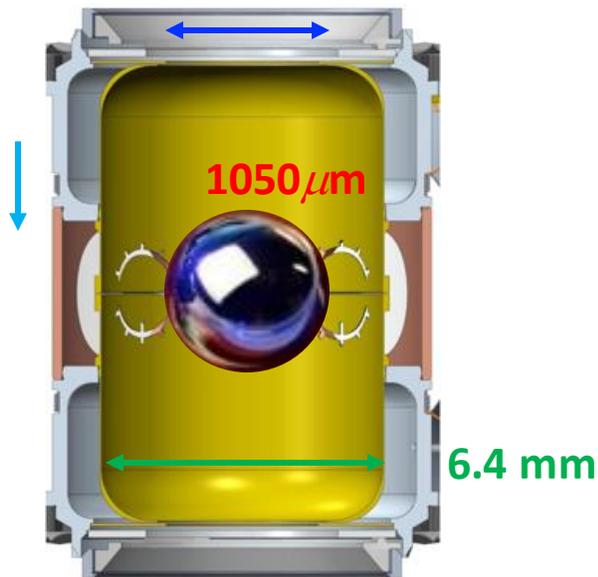
- Reducing the laser-entrance-hole size reduces the radiation losses

J. Ralph, GO04 Tuesday

The smaller laser-entrance-hole allows shorter coast time while maintaining implosion velocity

HYBRID-E design modified with smaller LEH³

3.65 → 3.1 mm LEH (27% less area)

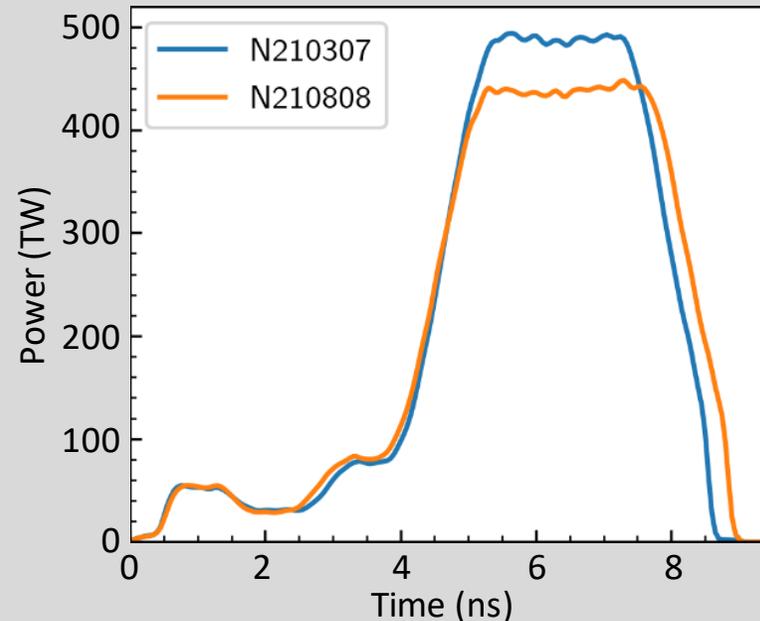


Small CCR 2.8

Used simulations and data driven models to re-tune symmetry for smaller LEH on first DT!

J. Ralph, GO04 Tuesday

Reduced radiation loss means can get same drive (Trad) with lower power

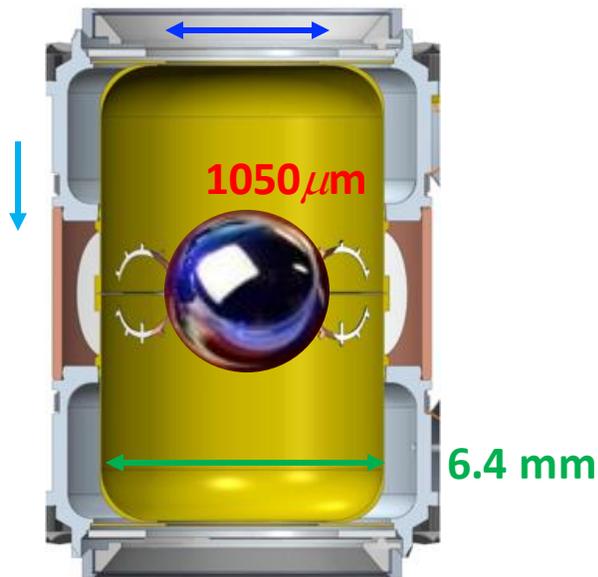


- Add the energy to the back end of the pulse to reduce coast time!

This allowed us to reduce the coast time another 350 ps – into the steep part of the curve

HYBRID-E design modified with smaller LEH³

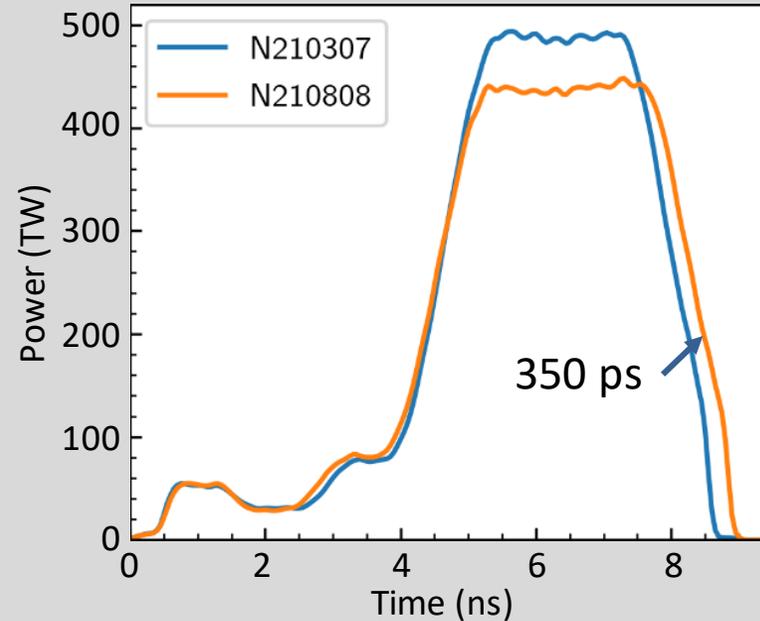
3.65 → 3.1 mm LEH (27% less area)



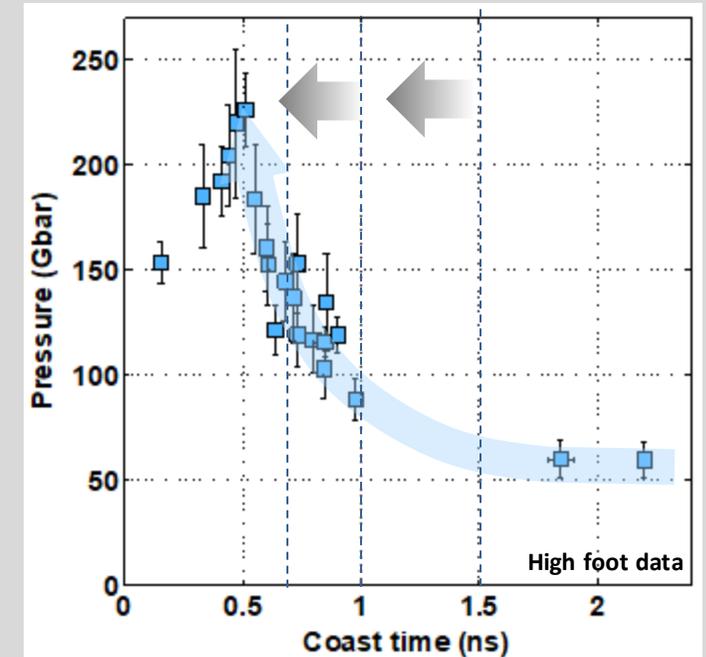
Small CCR 2.8

Used simulations and data driven models to re-tune symmetry for smaller LEH on first DT!

Reduced radiation loss means can get same drive (Trad) with lower power



Allowed us to reduce coast another 350 ps in steep part of the curve

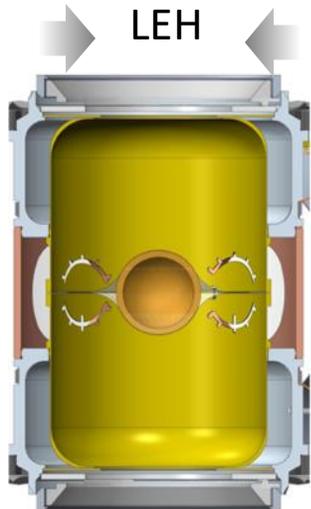


There may be room to make further improvements on this design

N210808 produced 1.35 MJ and the highest yield amplification of any shot to date

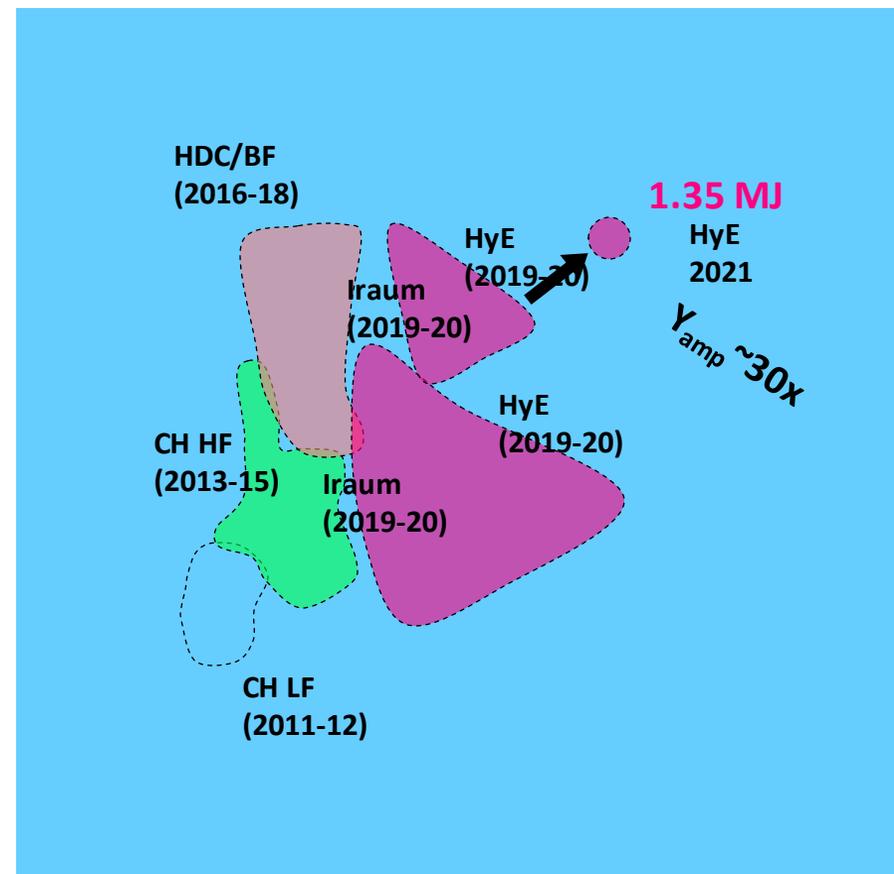
“HYBRID-E” Y = 1.3 MJ

3.6mm -> 3.1 mm



- > Reduced coast
- > Higher pressure
- > Better stability

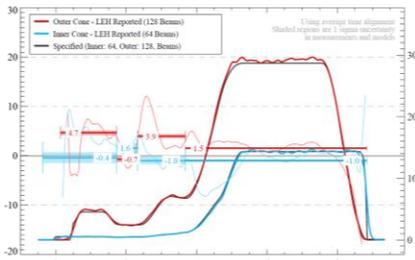
$$\text{Ignition figure of merit} \sim (\rho R)^3 T^3 \sim E_{HS} P_{HS}^2$$



There were also other important improvements on this shot

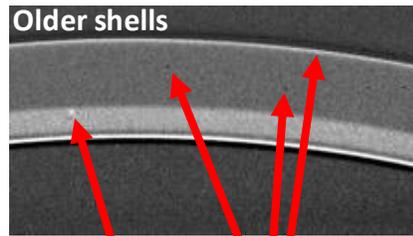
“HYBRID-E” Y = 1.3 MJ

Accumulated laser improvements



-> Precision energy delivery

Highest quality diamond capsule ever shot



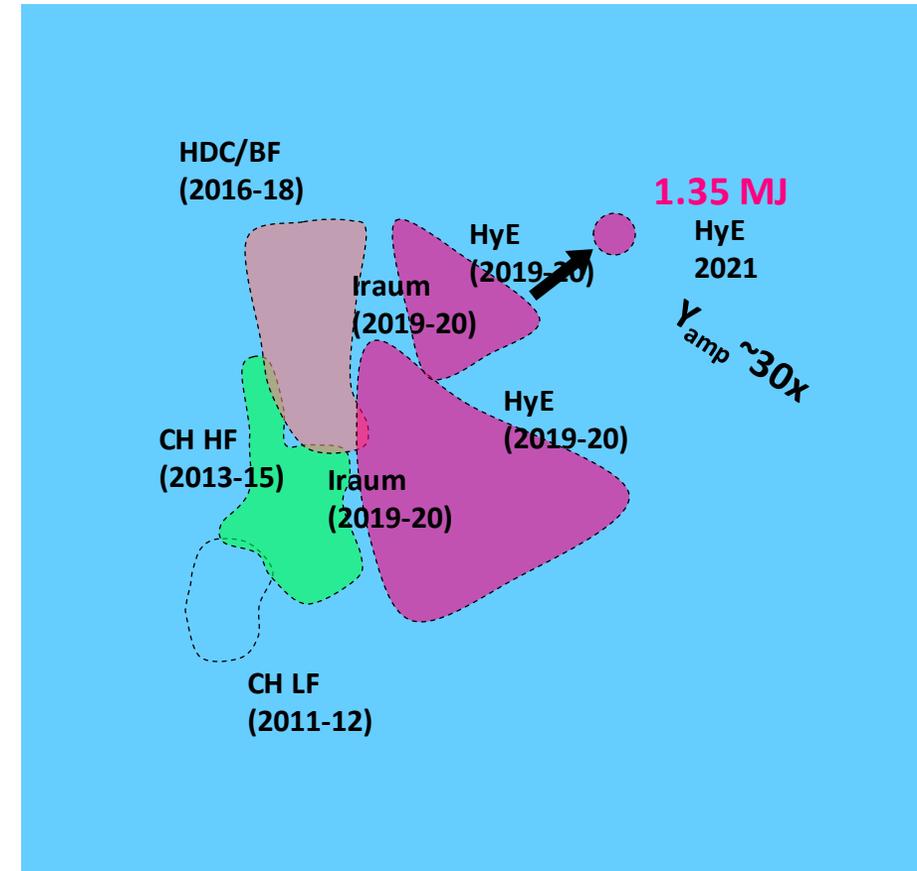
High-Z particle Voids
GA, Diamond Materials, LLNL

-> Less mix
-> Higher temperature

5µm -> 2µm fill tube

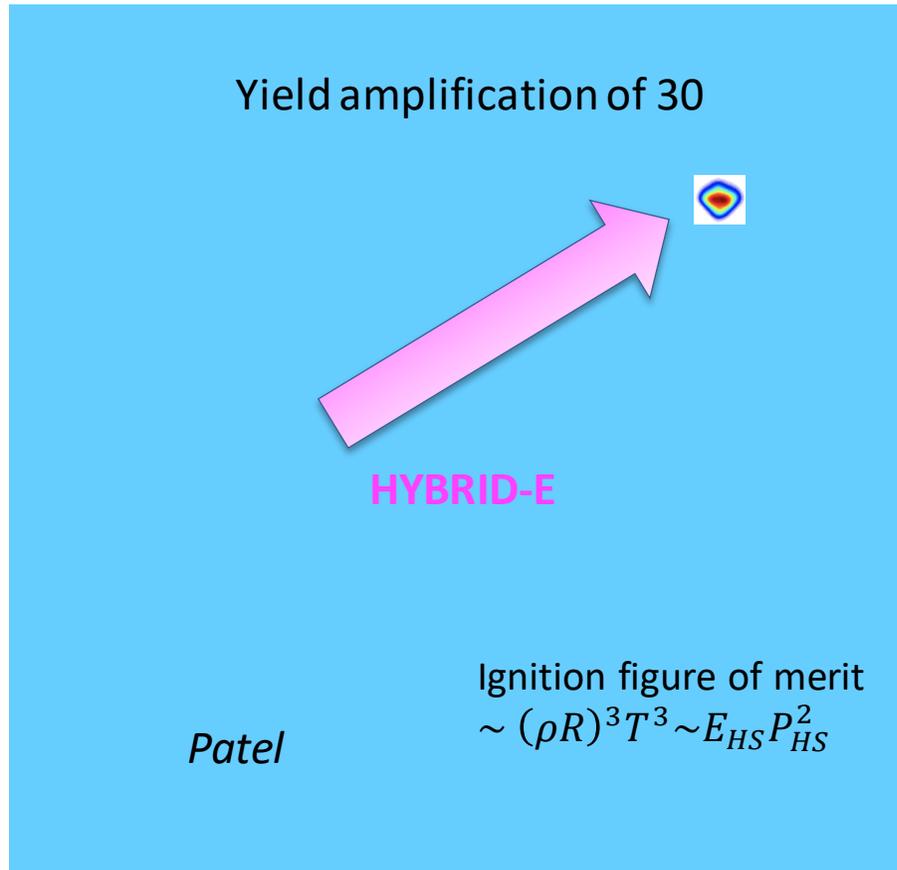


$$\text{Ignition figure of merit} \sim (\rho R)^3 T^3 \sim E_{HS} P_{HS}^2$$

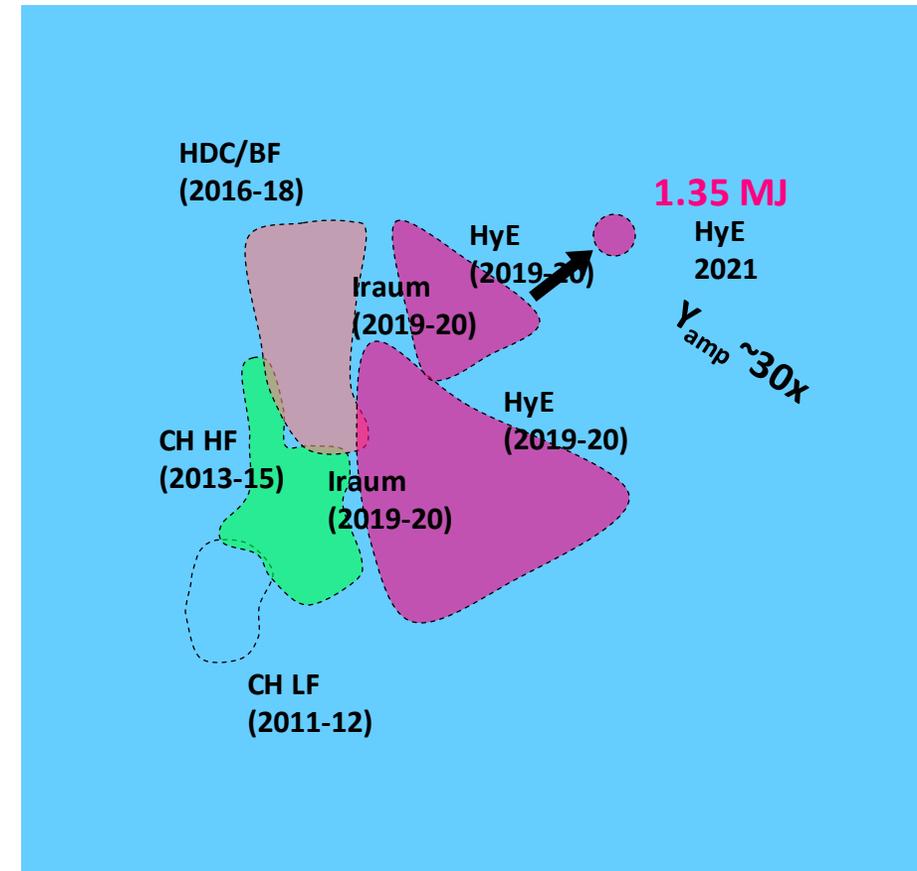


Future experiments will help understand the relative importance of each change

At the threshold of ignition, small changes in “no-alpha” space can lead to large changes in real “alpha-on” space

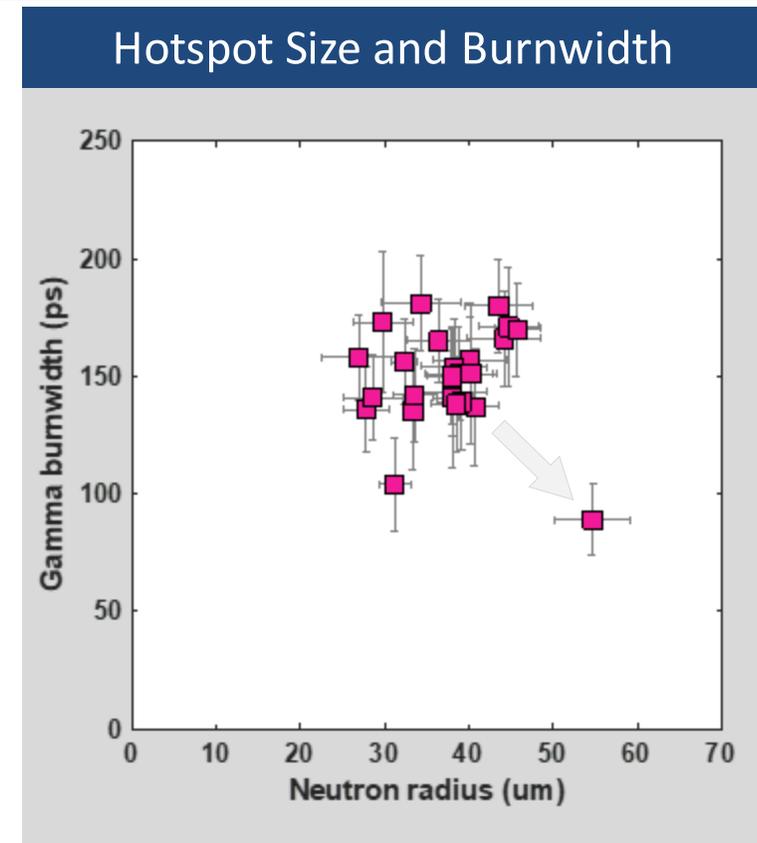
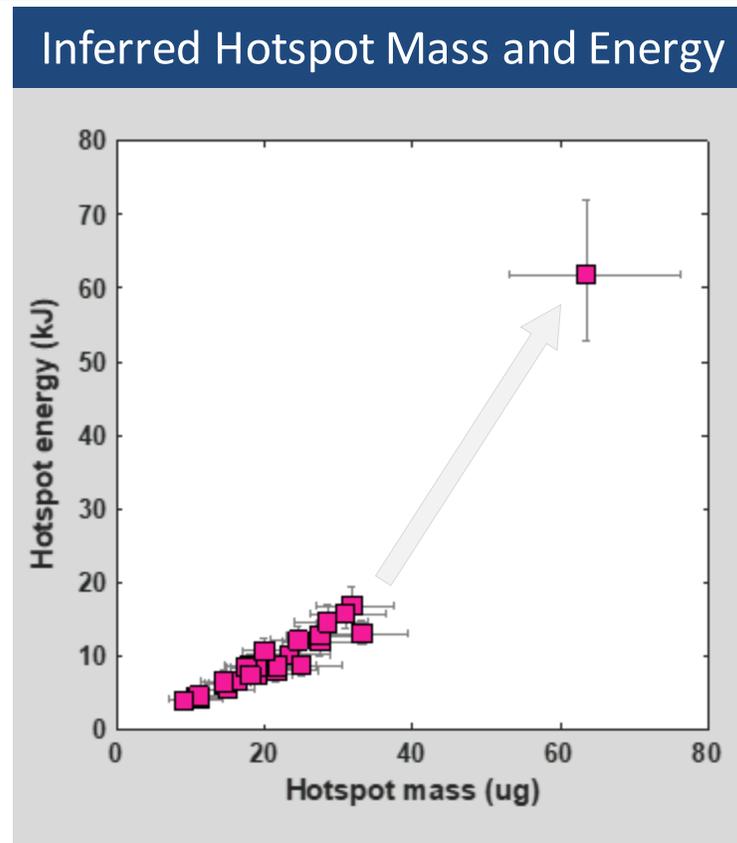
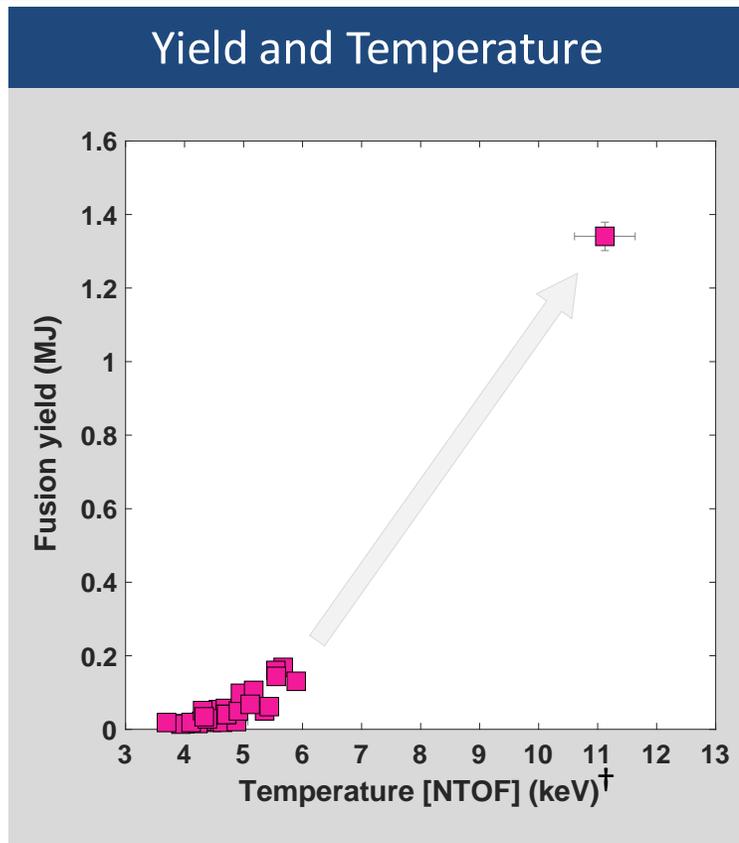


Ignition figure of merit $\sim (\rho R)^3 T^3 \sim E_{HS} P_{HS}^2$



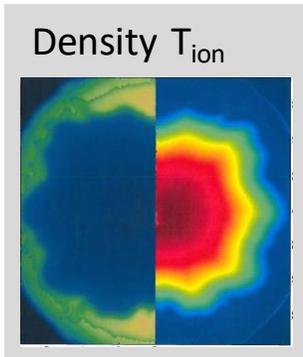
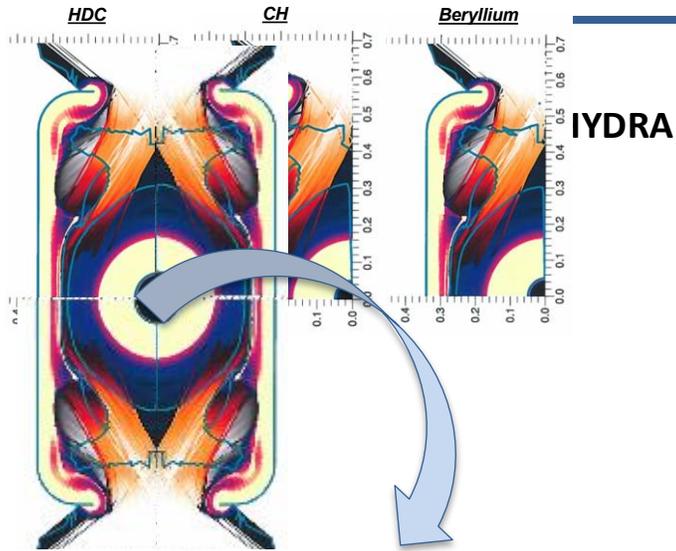
Much higher pressure and energy with small changes to input conditions indicates new regime

Many of the key experimental measurements are pointing to this implosion being in a fundamentally new regime

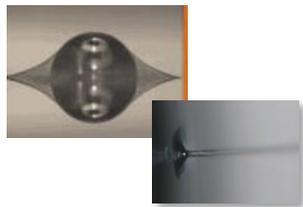


These are all signatures of a hotspot undergoing rapid self-heating and beginning to propagate burn into the surrounding dense shell

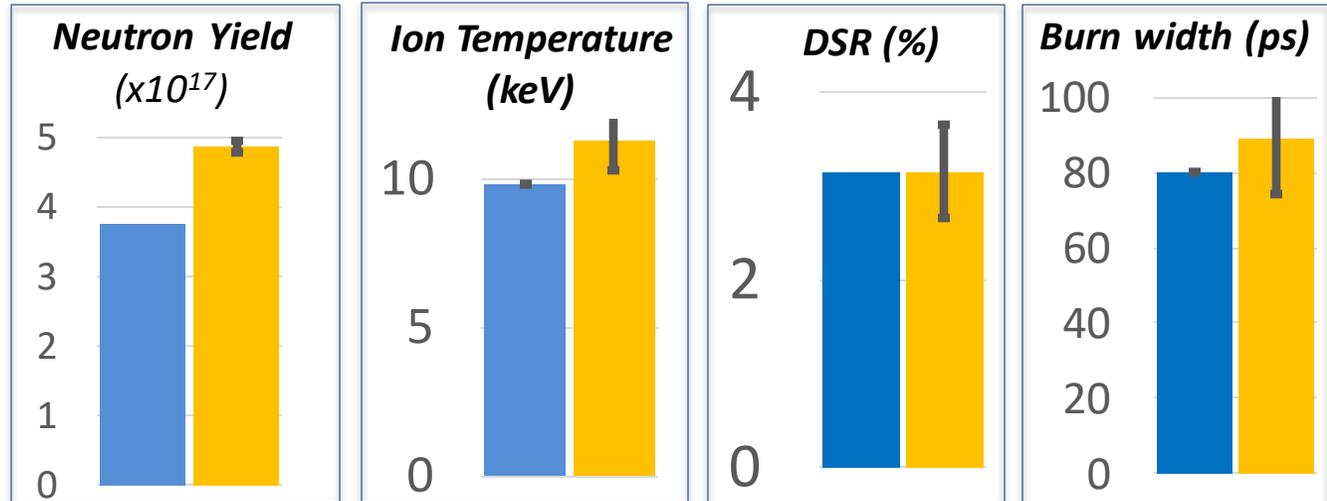
2D post-shot simulations capture many of the important implosion performance metrics in this new regime for this design



Higher resolution capsule simulations



Model for degradations is benchmarked against predecessor shots



Simulated Data

- Consistent with high temperature, large burning hot spot
- Preshot predicted increase of 3x in neutron yield but below data (7.9x)
 - Postshot, including as delivered laser, 2 um fill tube, observed asymmetry, agrees to 20% in yield

A. Kritcher GO04.00002

Ignition (a thermal instability), the plasma must overcome all energy losses for a duration of time

Cold DT shell

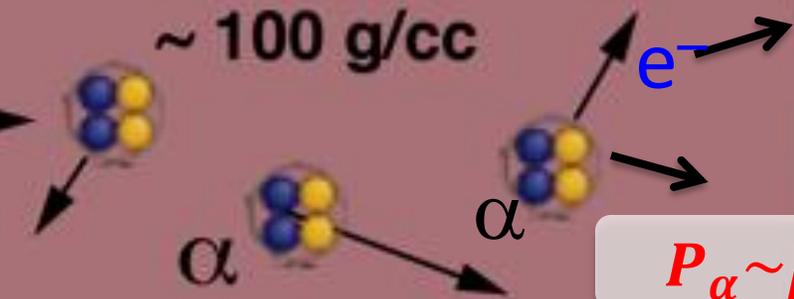
~ 1000 g/cc

Compressed DT fuel with hot central core

Fusion starts in DT hotspot

~ 50 million degrees

~ 100 g/cc



particles deposit energy
Bootstrap heating

Brems x-ray loss

$$P_B \sim \rho \sqrt{T}$$

Spitzer thermal conduction

$$P_e \sim T^{7/2} / (\rho R^2)$$

$$P_\alpha \sim \rho T^{3.6}$$

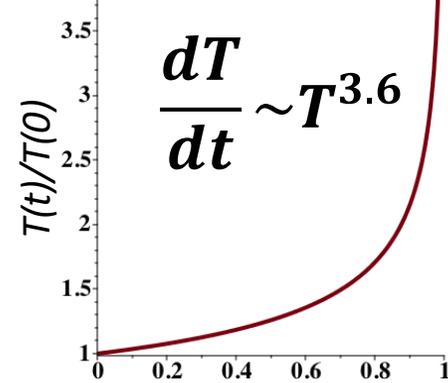
Alpha-heating

Time dependent heat balance (power/mass):

$$c_{DT} \frac{dT}{dt} = f_\alpha P_\alpha - f_B P_B - P_e - \frac{1}{m} p \frac{dV}{dt}$$

Ignition when these terms dominate

Thermonuclear instability

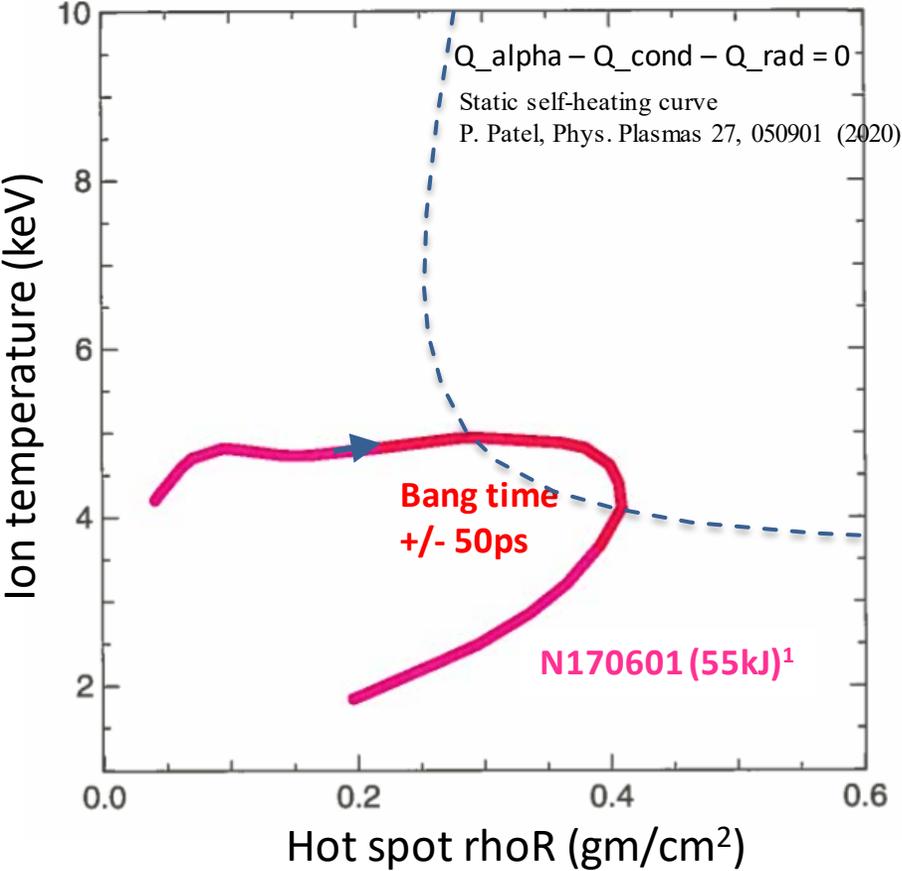


$t/\tau \quad \tau \sim 10's \text{ of picoseconds}$

Quantities are functions of ρ , R , and T so it is interesting to look at implosions in $\rho R / T$ space

Previous 55kJ experiments had alpha-particle heating but not enough to overcome losses

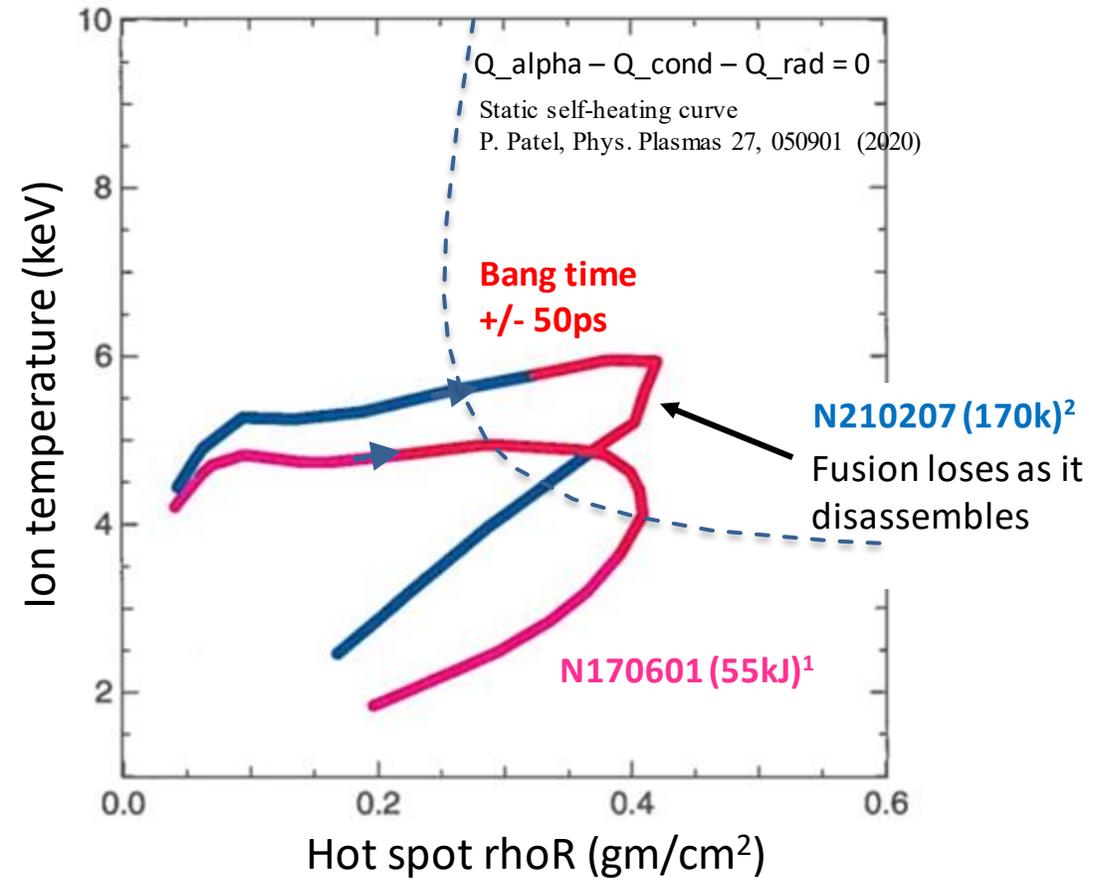
- 55kJ :
 - Alpha particle heating and hydro roughly balance radiation losses



1: Le Pape, et al, Phys. Rev. Lett. 120, 245003

170kJ shot -- alpha heating dominated implosions showed an increase in temperature but still succumb to losses

- 55kJ :
 - Alpha particle heating and hydro roughly balance radiation losses
- 170 kJ:
 - Higher initial hot spot temperature, alpha starting to win over losses but not there long enough for significant bootstrap heating

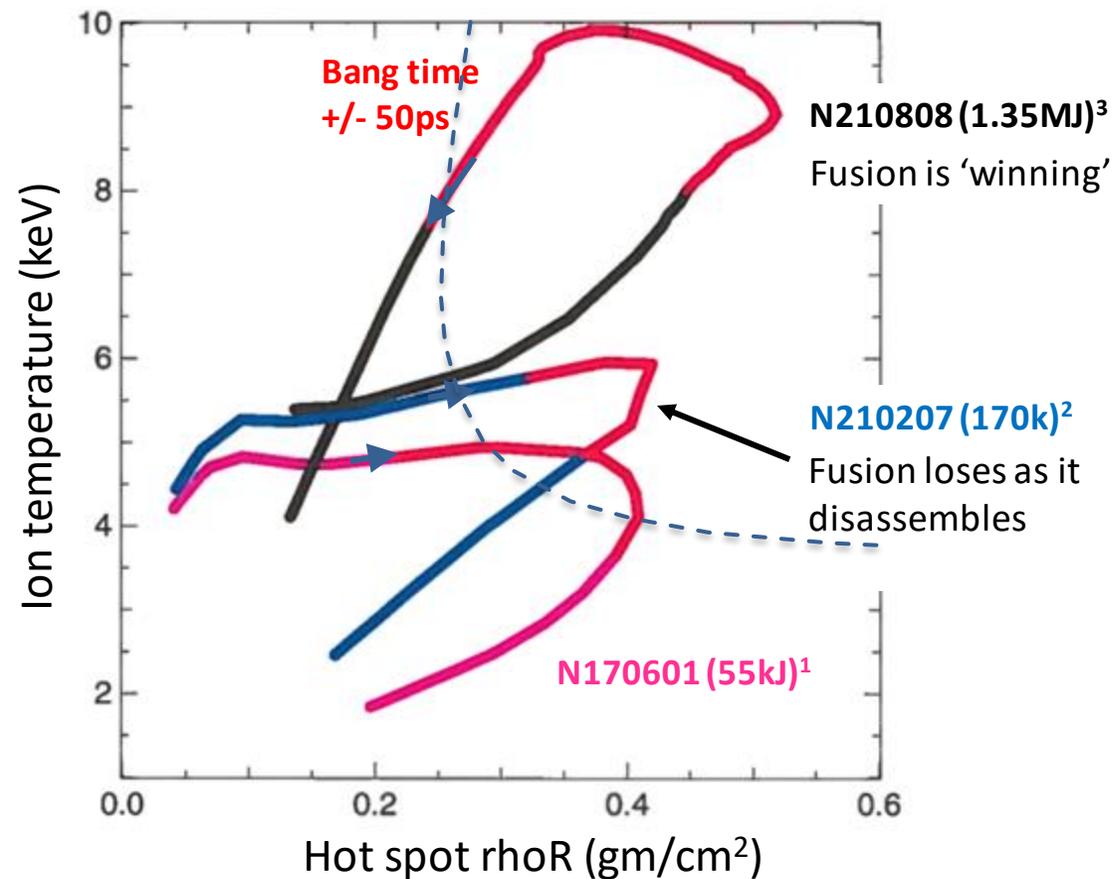


1: Le Pape, et al, Phys. Rev. Lett. 120, 245003

2: Zylstra, Hurricane, et al, in preparation; Kritcher, Young, Robey, et al, in preparation; Ross, Ralph, Zylstra et al, in preparation

N210808 shows substantial increases in calculated temperature and ρR and reversal of $\rho R - T_{ion}$ trajectory

- 55kJ :
 - Alpha particle heating and hydro roughly balance radiation losses
- 170 kJ:
 - Higher initial hot spot temperature, alpha starting to win over losses but not there long enough for significant bootstrap heating
- 1.35 MJ:
 - Similar hot spot formation but better confinement due to lower coast time allows alpha heating to outweigh losses and take off

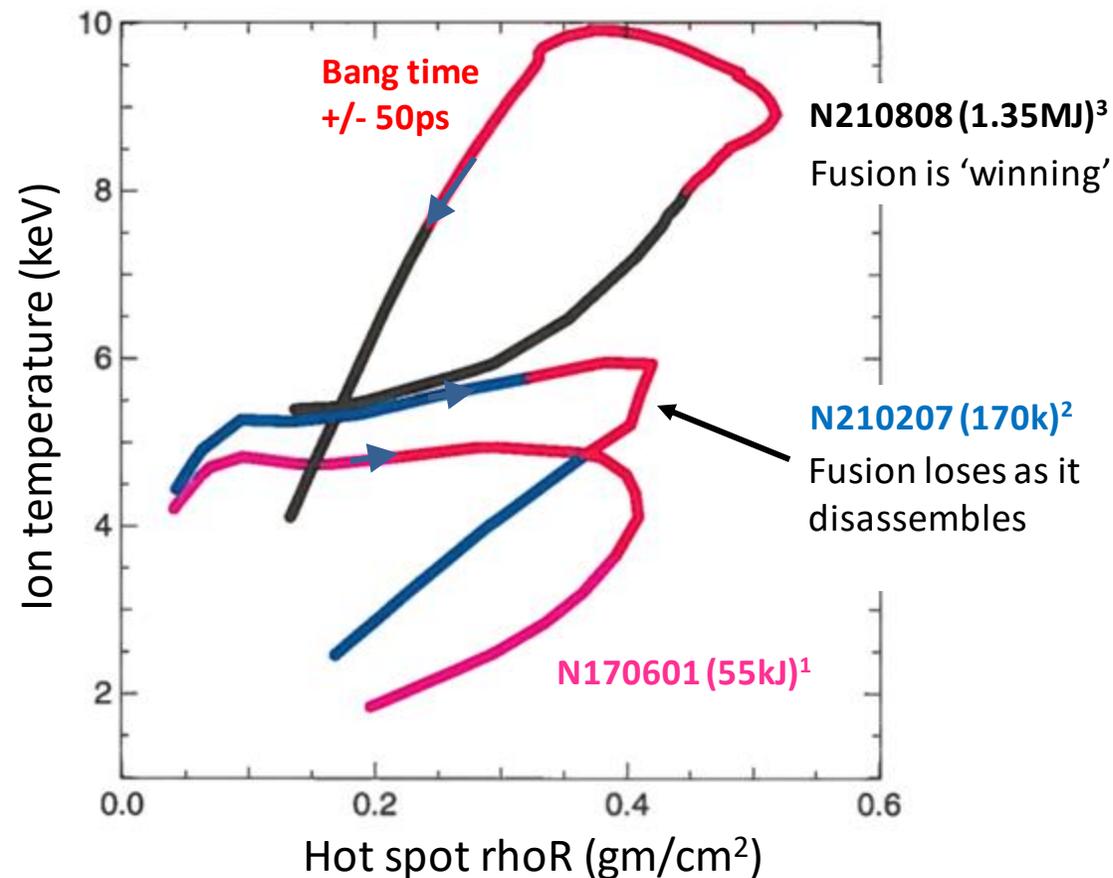


1: Le Pape, et al, Phys. Rev. Lett. 120, 245003

2: Zylstra, Hurricane, et al, in preparation; Kritcher, Young, Robey, et al, in preparation; Ross, Ralph, Zylstra et al, in preparation; 3: in preparation

We are now in a new regime where alpha energy is completely dominant

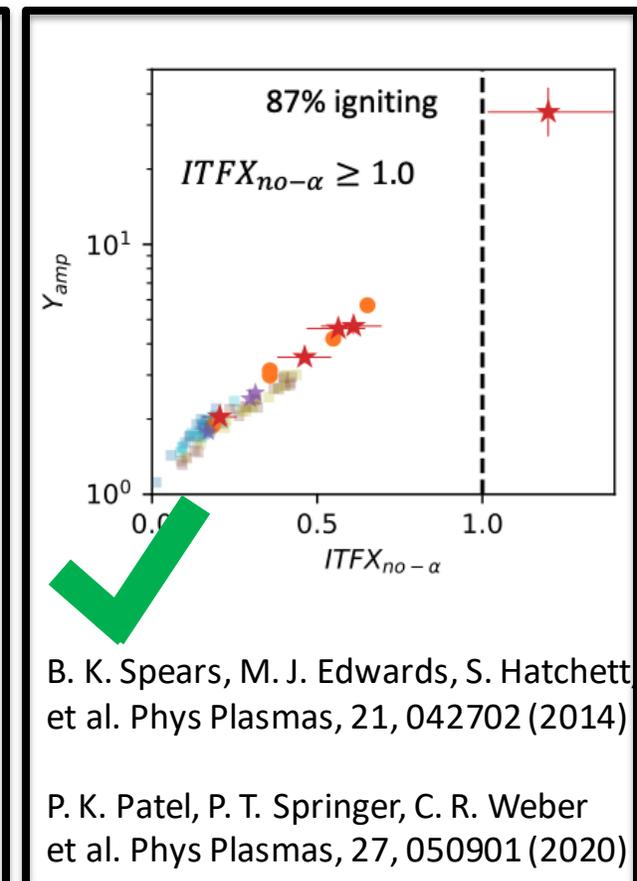
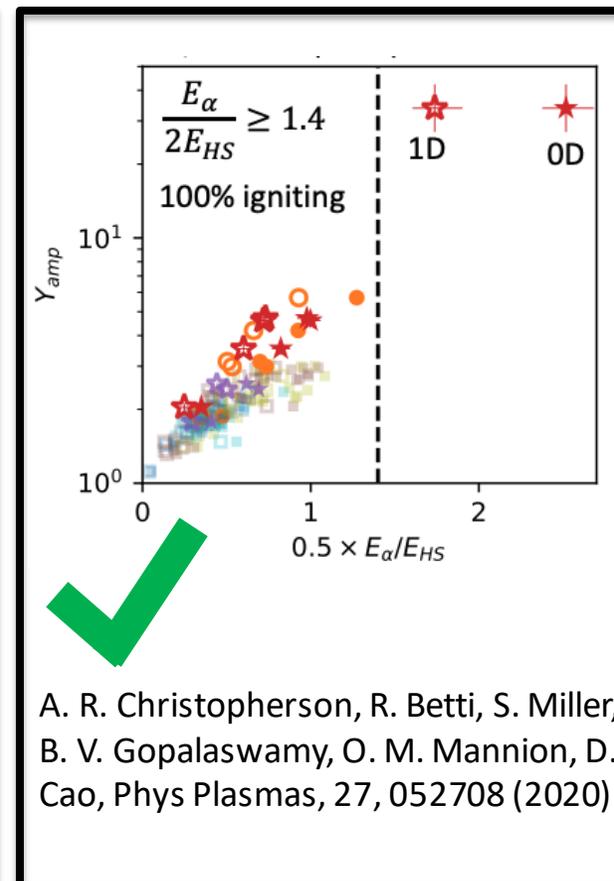
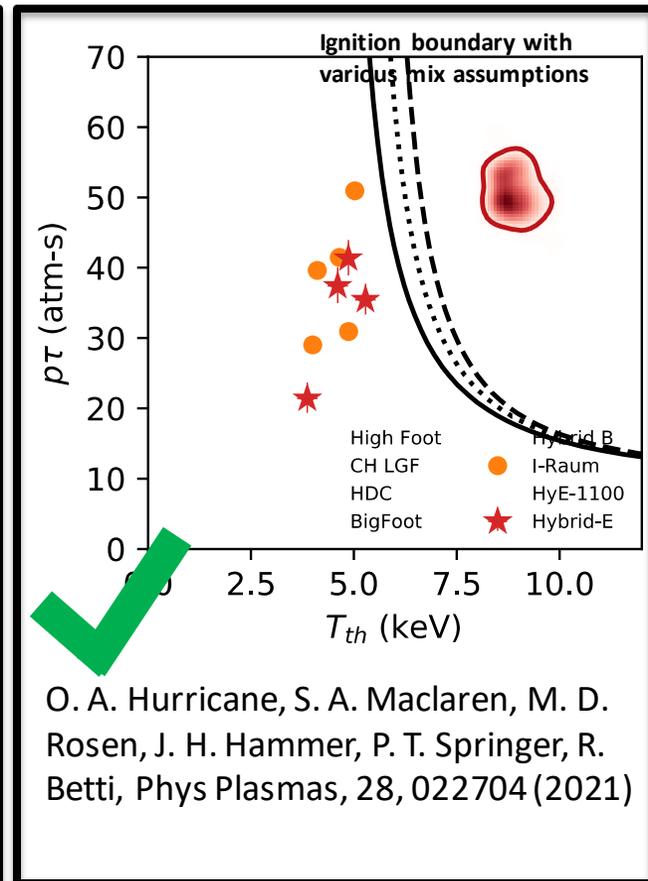
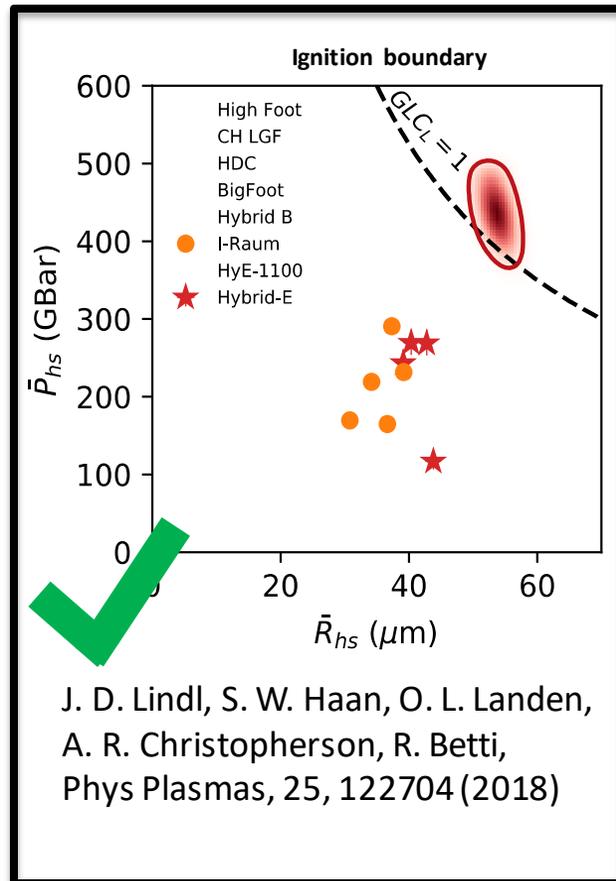
- Alpha energy >250 kJ
- PdV Work Done ~ 20 kJ
- Radiation loss ~ 60 kJ
- Total fusion energy and power: 1.35 MJ, 15 PW (quadrillion) for ~90 ps !!



1: Le Pape, et al, Phys. Rev. Lett. 120, 245003

2: Zylstra, Hurricane, et al, in preparation; Kritcher, Young, Robey, et al, in preparation; Ross, Ralph, Zylstra et al, in preparation; 3: in preparation

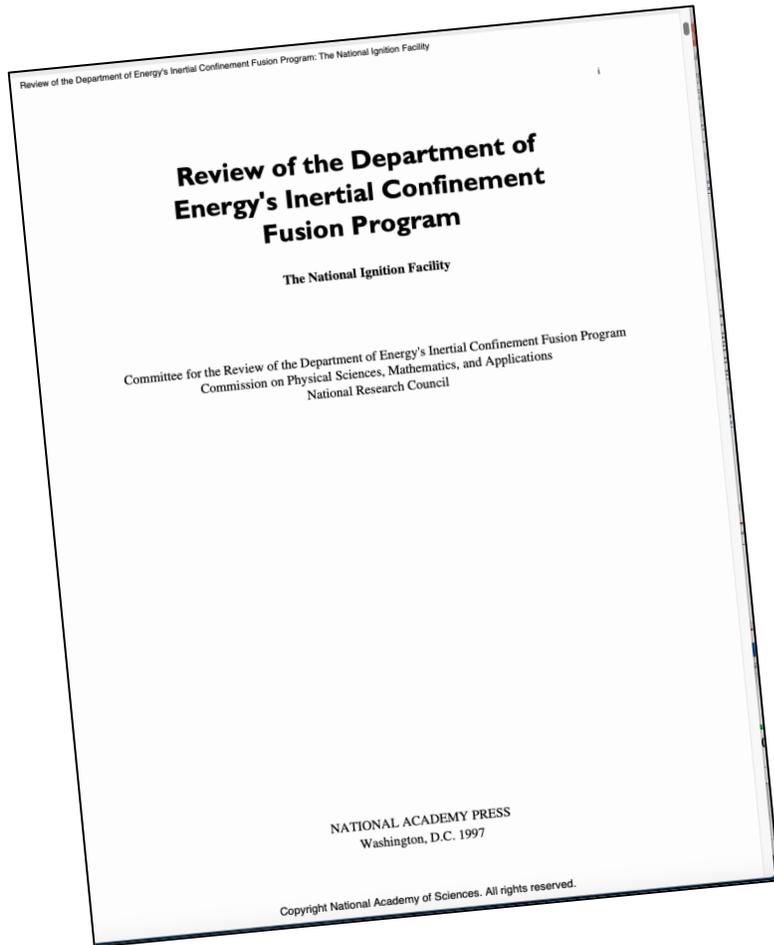
There are several published metrics for inferring ignition – we are evaluating this shot against them



See talks: A. Christopherson CO04:00005
A. Zylstra QI02.00001

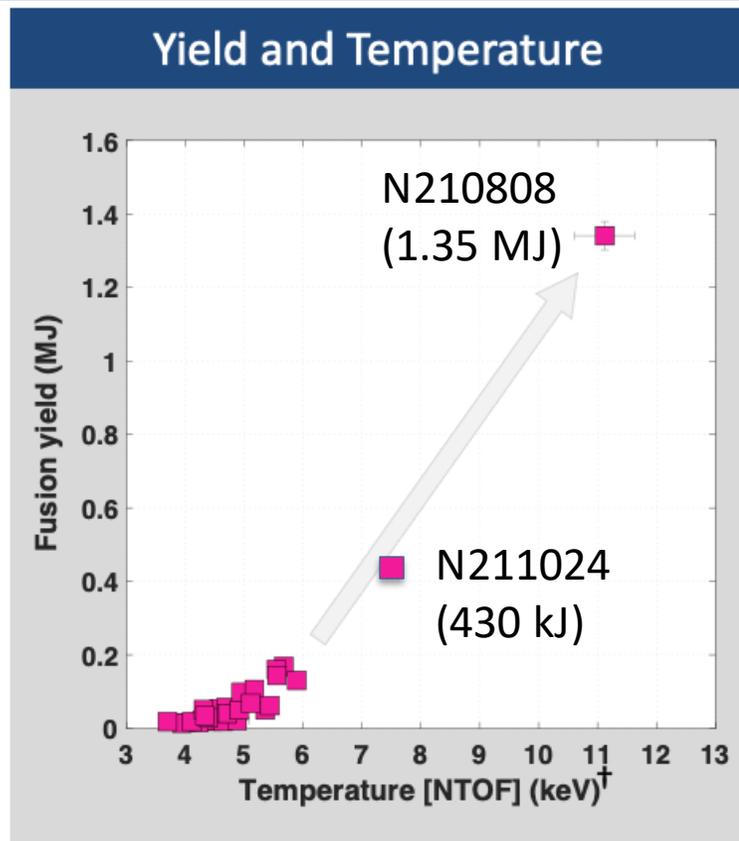
Analysis: A. Zylstra

The 1997 National Academy of Science review of the ICF program defined “ignition” as “gain = 1”

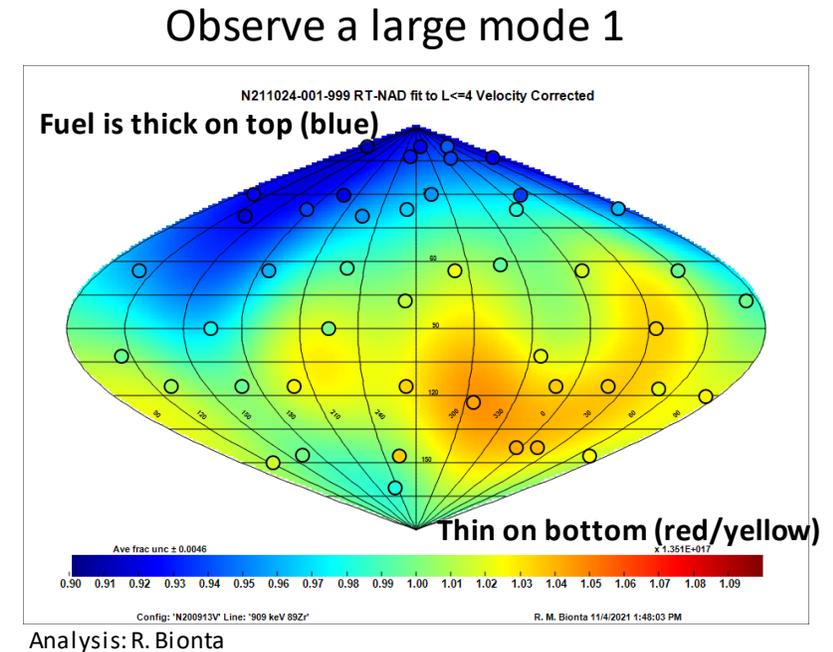


- Ignition is a statement about the power balance in the hotspot of the implosion
- NAS decided to use a definition that is not about power balance but about target gain
 - Easy to evaluate
- NAS defined “ignition” as target gain=1, where fusion yield is equal to laser energy
- This shot is target gain 0.7

First shot to assess variability on Oct 24 gave ~ 430 kJ of yield



- Second shot with capsule gain > 1 and yield > 1e17 neutrons!
- Data still being analyzed
- Simulations in progress
- Expect large variability near ignition cliff

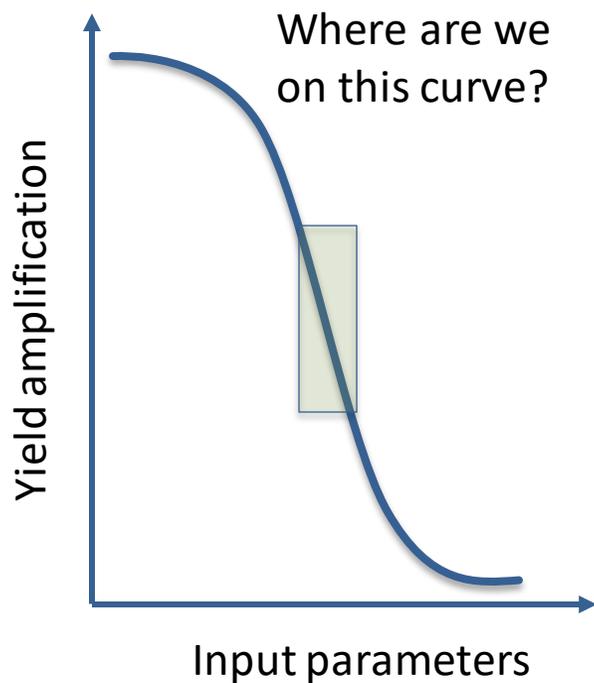


130 km/s velocity downward
(N210808 was 68 km/s)

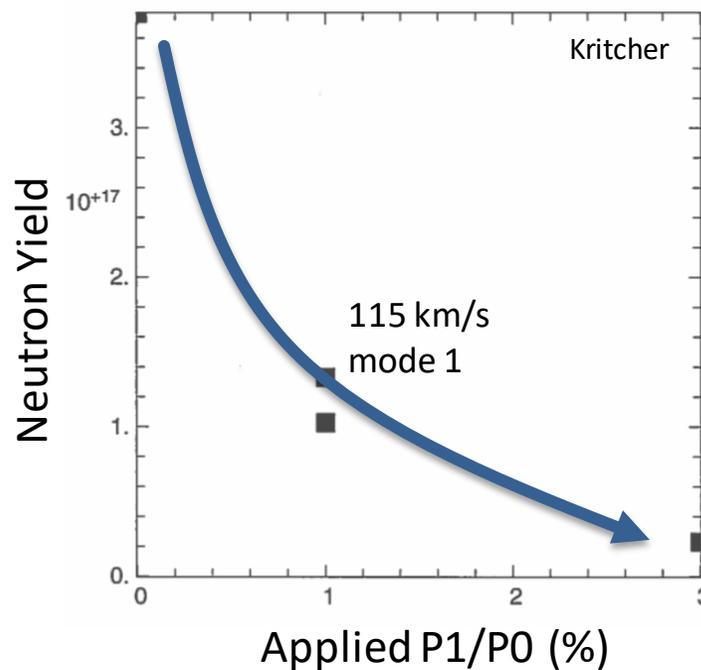
Mode 1: Schlossberg NO04.00010

Near ignition cliff, expect variability -- the designs are sensitive to input parameters and sensitivity increases with yield amplification

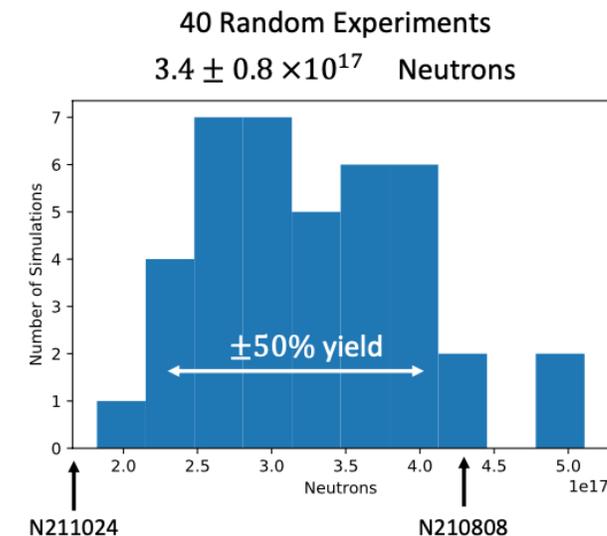
Because ignition is a threshold process we expect sensitivity to increase with higher yield amplification



Example: 2d simulations of N210808 show strong sensitivity to mode 1 (~ 2x between N210808 and repeat)

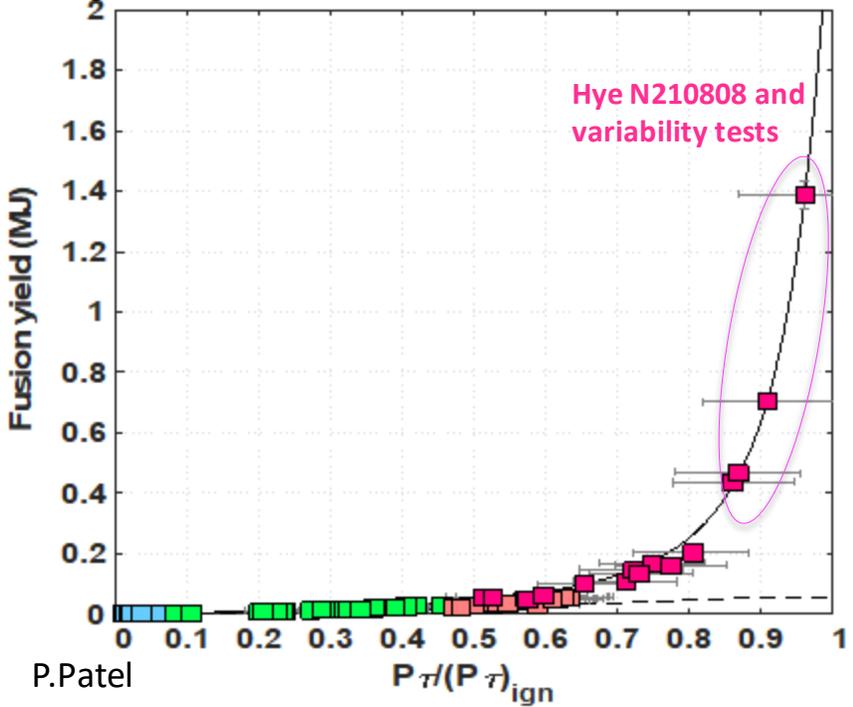
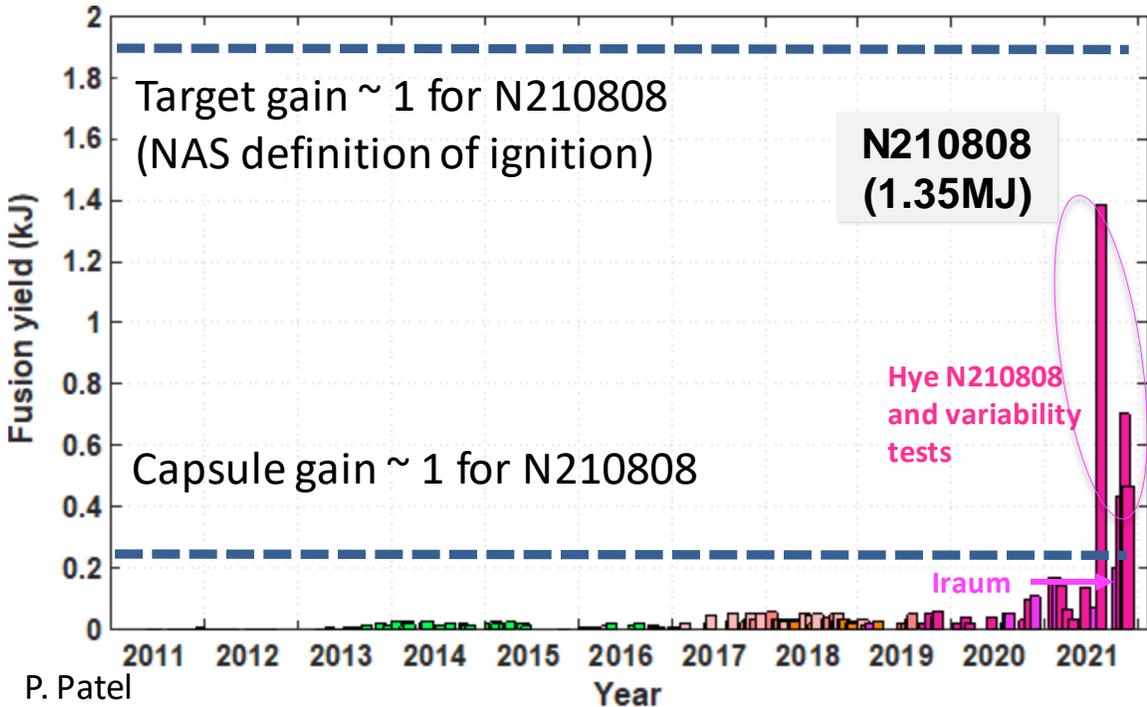


For N210808, 1d simulations show $\pm 25\%$ variability in yield from laser and target variations

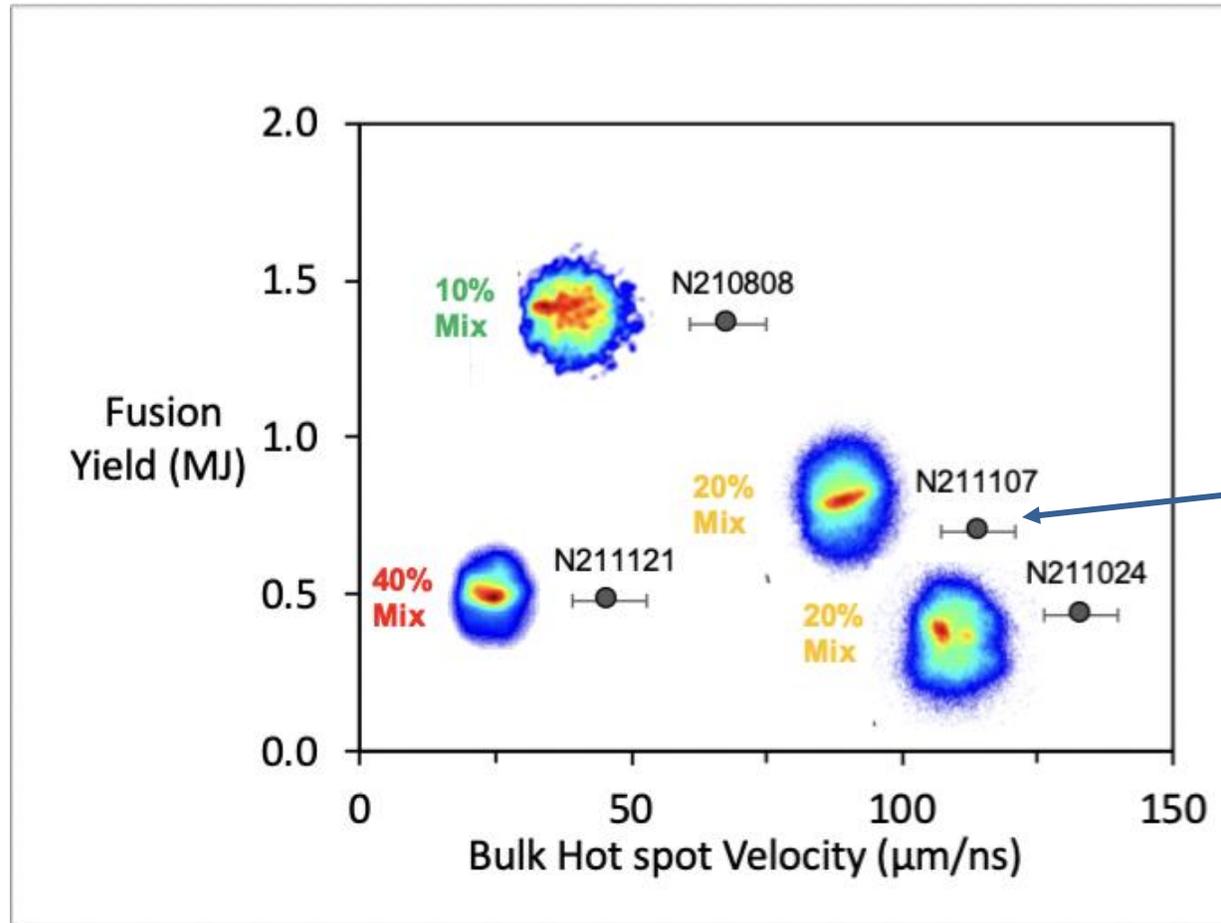


See L. Peterson TO07.00009

Three experiments testing N210808 variability have been performed: all reached $>1e17$ and capsule gain >2 , well beyond early 2021 shots



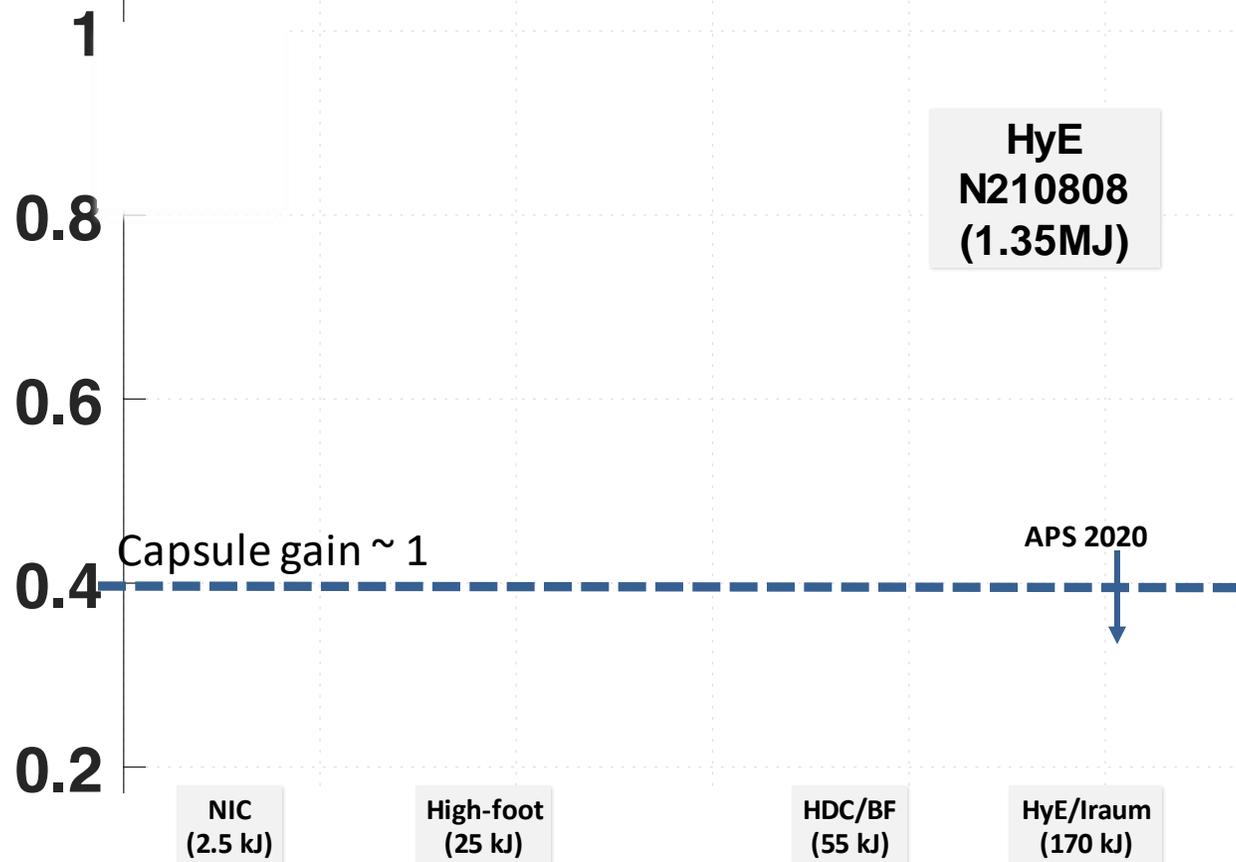
Variability mainly attributed to larger unintentional odd mode implosion asymmetry or more mix



5 μm fill tube

Working on understanding and mitigating these degradations to improve robustness

The August 8th NIF shot (N210808) yielded more than 1.3 MJ and marks a significant advance in ICF research



- Capsule gain (yield/capsule absorbed energy) > 5
- Laser gain ~ 0.7
- Burn propagation ~ 2% burn up
- 13x increase in fusion yield since APS 2020

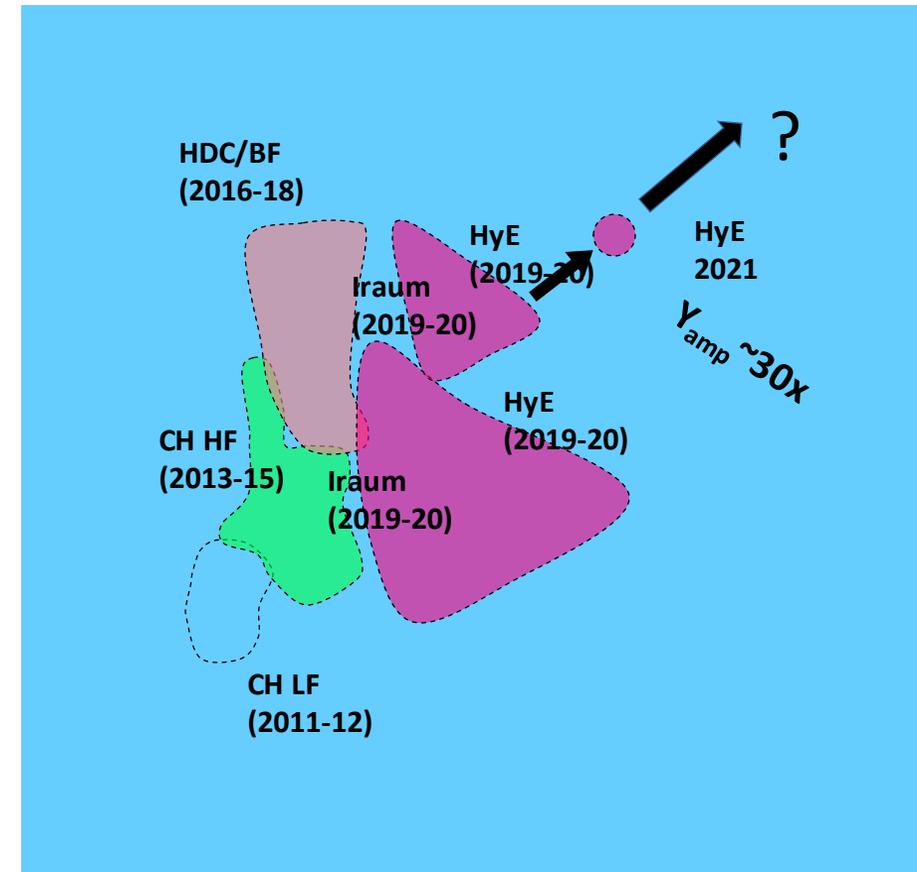
We've been waiting for this since Nuckolls 1972 Nature paper!

We have made much progress in the last few years but much work remains to be done

- Understand where we are
 - Assess variability
 - Assess sensitivity to input parameters
 - “Dudded” fuel to validate models in this Y_{amp} range
- Push to higher performance
 - Higher compression
 - Higher laser energy
 - Further improvements to hohlraum
- Use the current design for science experiments
 - Output from this implosion is most powerful/energetic driver that we have for HED
- Improve our simulation capabilities

Kritcher GO04.00002, Clark ZI02.00002, Casey CO04.00006

Ignition figure of merit $\sim (\rho R)^3 T^3 \sim E_{\text{HS}} P_{\text{HS}}^2$



This is a very exciting time for our field!

“Don’t stop believing”
-- Hybrid E theme song



Disclaimer

This document was prepared as an account of work sponsored by an agency of the United States government. Neither the United States government nor Lawrence Livermore National Security, LLC, nor any of their employees makes any warranty, expressed or implied, or assumes any legal liability or responsibility for the accuracy, completeness, or usefulness of any information, apparatus, product, or process disclosed, or represents that its use would not infringe privately owned rights. Reference herein to any specific commercial product, process, or service by trade name, trademark, manufacturer, or otherwise does not necessarily constitute or imply its endorsement, recommendation, or favoring by the United States government or Lawrence Livermore National Security, LLC. The views and opinions of authors expressed herein do not necessarily state or reflect those of the United States government or Lawrence Livermore National Security, LLC, and shall not be used for advertising or product endorsement purposes.