

### **Improving our models of nonlocal transport**

**C.P RIDGERS** 

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# What can you do if you dont have an ignition-scale facility?

• Large-scale ICF experiments inherently multi-scale – necessitates simple models for kinetic processes

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• Mean-free path of heat carrying electrons long



J. Brodrick PhD thesis (2018)

- Mean-free path of heat carrying electrons long
- Can't use diffusive LTE transport



J. Brodrick PhD thesis (2018)

 $T_e$ 

- Mean-free path of heat carrying electrons long
- Can't use diffusive LTE transport
- Solve kinetic equation for 'right' answer – can be slow
- Reduced non-local models must be used



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- Flux limiter arbitrarily limit heat flux when unphysical
- Reduced kinetic models also available (e.g. SNB)

• Ab-intio calculation – Vlasov-Fokker-Planck (VFP) equation

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \frac{\partial f}{\partial \mathbf{r}} - \frac{e}{m_e} (\mathbf{E} + \mathbf{v} \times \mathbf{B}) \cdot \frac{\partial f}{\partial \mathbf{v}} = \left(\frac{\partial f}{\partial t}\right)_{coll}$$

• Ab-intio calculation – Vlasov-Fokker-Planck (VFP) equation

$$\begin{array}{c} \overbrace{\partial f} \\ \hline \partial t \end{array} \mathbf{v} \cdot \frac{\partial f}{\partial \mathbf{r}} - \frac{e}{m_e} (\mathbf{E} + \mathbf{v} \times \mathbf{B}) \cdot \frac{\partial f}{\partial \mathbf{v}} = \left(\frac{\partial f}{\partial t}\right)_{col} \\ \end{array}$$
Solve for distribution

**function** – **non-LTE** 

• Ab-intio calculation – Vlasov-Fokker-Planck (VFP) equation

$$\begin{array}{c} \overrightarrow{\partial f} \\ \overrightarrow{\partial t} \\ \overrightarrow{\partial t} \\ \overrightarrow{\partial r} \\ \overrightarrow{\partial r}$$

• Ab-intio calculation – Vlasov-Fokker-Planck (VFP) equation



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R.J. Kingham & A.R. Bell, J. Comp. Phys., 194, 1 (2004)

• Ab-intio calculation – Vlasov-Fokker-Planck (VFP) equation



• Flux limiter – arbitrarily limit heat flux when unphysical

$$\frac{1}{q} = \frac{1}{q_{LTE}} + \frac{1}{fq_{FS}}$$

• Ab-intio calculation – Vlasov-Fokker-Planck (VFP) equation



LTE (Braginskii) heat flow

• Flux limiter – arbitrarily limit heat flux when unphysical



• Ab-intio calculation – Vlasov-Fokker-Planck (VFP) equation



• Flux limiter – arbitrarily limit heat flux when unphysical

 $\frac{1}{q} = \underbrace{\frac{1}{q_{LTE}} fq_{FS}}_{\text{LTE}} + \underbrace{fTE}_{\text{FS}} + \underbrace{\text{Braginskii}}_{\text{Hax heat flow}} + \underbrace{\text{free streaming limit (}q_{FS})}_{\text{Arbitrary limiter}} = f$ 

• Ab-intio calculation – Vlasov-Fokker-Planck (VFP) equation

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \frac{\partial f}{\partial \mathbf{r}} - \frac{e}{m_e} (\mathbf{E} + \mathbf{v} \times \mathbf{B}) \cdot \frac{\partial f}{\partial \mathbf{v}} = \left(\frac{\partial f}{\partial t}\right)_{coll}$$

• Reduced kinetic models (e.g. SNB)

**Quickly reach steady state heat flow** 

• Ab-intio calculation – Vlasov-Fokker-Planck (VFP) equation

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \frac{\partial f}{\partial \mathbf{r}} - \frac{\mathbf{e}}{\mathbf{w}_{e}} (\mathbf{E} + \mathbf{v} \times \mathbf{B}) \cdot \frac{\partial f}{\partial \mathbf{v}} = \left(\frac{\partial f}{\partial t}\right)_{coll}$$

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**Electric field LTE, most ignore B-fields** 

• Ab-intio calculation – Vlasov-Fokker-Planck (VFP) equation



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#### **Use Krook operator**

• Ab-intio calculation – Vlasov-Fokker-Planck (VFP) equation



• Reduced kinetic models (e.g. SNB)

Quickly reach steady state heat flow

Electric field LTE, most ignore B-fields

**Use Krook operator** 
$$\left(\frac{\partial f}{\partial t}\right)_{coll} = -v_c(f-f_M)$$

• Ab-intio calculation – Vlasov-Fokker-Planck (VFP) equation



• Reduced kinetic models (e.g. SNB)

Quickly reach steady state heat flow

Electric field LTE, most ignore B-fields

Use Krook operator  $\left(\frac{\partial f}{\partial t}\right)_{coll} = -v_c(f-f_M)$ Make f0+f1 approximation

• Ab-intio calculation – Vlasov-Fokker-Planck (VFP) equation

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \frac{\partial f}{\partial \mathbf{r}} - \underbrace{\left( \mathbf{E} + \mathbf{v} \times \mathbf{B} \right)}_{\mathcal{O}} \frac{\partial f}{\partial \mathbf{v}} = \underbrace{\left( \frac{\partial f}{\partial t} \right)}_{coll}_{coll}$$
 (Classic' SNE

• Reduced kinetic models (e.g. SNB) Quickly reach steady state heat flow Electric field LTE, most ignore B-fields Use Krook operator  $\left(\frac{\partial f}{\partial t}\right)_{coll} = -v_c(f - f_M)$ Make f0+f1 approximation

$$\left[\frac{1}{\lambda_g} - \nabla \cdot \left(\frac{\lambda_g}{3}\nabla\right)\right] H_g = \nabla \cdot \boldsymbol{U}_g$$

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 $\begin{bmatrix} 1 & & \lambda_g \\ \lambda_g & & \lambda_g \\ \end{pmatrix} H_g = \nabla \cdot U_g$ Mfp for velocity group

• Ab-intio calculation – Vlasov-Fokker-Planck (VFP) equation

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \frac{\partial f}{\partial \mathbf{r}} - \underbrace{\left( \mathbf{E} + \mathbf{v} \times \mathbf{B} \right)}_{\mathcal{B}} \frac{\partial f}{\partial \mathbf{v}} = \underbrace{\left( \frac{\partial f}{\partial t} \right)}_{coll} \text{ (Classic' SNB)}$$

• Reduced kinetic models (e.g. SNB) Quickly reach steady state heat flow Electric field LTE, most ignore B-fields Use Krook operator  $\left(\frac{\partial f}{\partial t}\right)_{coll} = -v_c(f-f_M)$ Make f0+f1 approximation • Reduced kinetic models (e.g. SNB)  $\left[\frac{1}{\lambda_g} - \nabla \cdot \begin{pmatrix}\lambda_g \\ 3 \\ 3 \end{pmatrix}\right] H_g = \underbrace{\nabla \cdot U_g}_{coll}$ 

• Ab-intio calculation – Vlasov-Fokker-Planck (VFP) equation

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term

• Ab-intio calculation – Vlasov-Fokker-Planck (VFP) equation

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \frac{\partial f}{\partial \mathbf{r}} - \underbrace{\left( \mathbf{E} + \mathbf{v} \times \mathbf{B} \right)}_{\mathcal{B}} \frac{\partial f}{\partial \mathbf{v}} = \underbrace{\left( \frac{\partial f}{\partial t} \right)}_{coll}_{coll}$$
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$$\frac{1}{\lambda_g} - \nabla \cdot \left(\frac{\lambda_g}{3}\nabla\right) = \nabla \cdot \boldsymbol{U}_g$$

$$\boldsymbol{q} = \boldsymbol{q}_{LTE} - \frac{m_e}{2} \int \frac{\lambda_g}{3} \nabla H_g dv$$

### High flux model



M.D. Rosen et al., HEDP, 7, 180 (2011)

# High flux model

• XSN + f=0.05 first model for NIC



M.D. Rosen et al., HEDP, 7, 180 (2011)

# High flux model

- XSN + f=0.05 first model for NIC
- DCA + Nonlocal

Better NLTE physics but slow

M.D. Rosen et al., HEDP, 7, 180 (2011)





• B-fields grow via Biermann battery



- B-fields grow via Biermann battery
- Modify heat tranport



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- B-fields grow via Biermann battery
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• Heat transport determines B-field evolution
- Heat transport determines B-field evolution
- Braginskii's Ohm's law + Faraday

$$e n_e (\mathbf{E} + \mathbf{C}_i \times \mathbf{B}) = -\nabla P_e + \frac{\underline{\alpha} \cdot \mathbf{j}}{e n_e} + n_e \underline{\beta} \cdot \nabla T_e \qquad \frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E}$$

- Heat transport determines B-field evolution
- Braginskii's Ohm's law + Faraday

$$e n_e (\mathbf{E} + \mathbf{C}_i \times \mathbf{B}) = -\nabla P_e + \frac{\underline{\alpha} \cdot \mathbf{j}}{e n_e} + n_e \underline{\beta} \cdot \nabla T_e \qquad \frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E}$$
  
Frozen-in flow

- Heat transport determines B-field evolution
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$$e n_e (\mathbf{E} + \underbrace{\mathbf{C}_i \times \mathbf{B}}_{i}) = \underbrace{\nabla P}_{e} + \frac{\underline{\alpha} \cdot \mathbf{j}}{e n_e} + n_e \underline{\beta} \cdot \nabla T_e \qquad \frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E}$$
  
Frozen-in flow **Biermann**

- Heat transport determines B-field evolution
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$$e n_{e} (\mathbf{E} + (\mathbf{C}_{i} \times \mathbf{B})) = (\nabla P_{e} + (\underline{\underline{a} \cdot j})) + (\underline{\underline{a} \cdot j}) + (\underline{\underline{a} \cdot j$$

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W.A. Farmer et al., Phys. Plasmas, 24, 052703 (2017)



W.A. Farmer et al., Phys. Plasmas, 24, 052703 (2017)



W.A. Farmer et al., Phys. Plasmas, 24, 052703 (2017)

• Magnetize hohlraum to imporove capsule yield



J. Moody et al., J. Fusion Energy, 41, 7 (2022)

- Magnetize hohlraum to imporove capsule yield
- Simulations indicate Nernst effect crucial in magnetized hohlraums



J. Moody et al., J. Fusion Energy, 41, 7 (2022)

- Magnetize hohlraum to imporove capsule yield
- Simulations indicate Nernst effect crucial in magnetized hohlraums
- VFP simulations after 400ps with 7.5T applied field, density 5x10<sup>19</sup>Wcm<sup>-2</sup>. Ray tracing, intensity ~5x10<sup>14</sup>Wcm<sup>-2</sup>



A. Joglekar et al., Phys. Rev. E, 93, 043206 (2016)

• Magnetic fields modify degree of nonlocality



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- Magnetic field evolution non-local



- Magnetic fields modify degree of nonlocality
- Magnetic field evolution non-local
- Need non-local models which selfconsistently include B-field



$$\frac{\partial \boldsymbol{B}}{\partial t} \approx -\frac{1}{e n_e} \nabla n_e \times \nabla T_e - \nabla \cdot (\boldsymbol{v}_N \boldsymbol{B}_z) \qquad \boldsymbol{v}_N \propto \boldsymbol{q}_e$$

- Magnetic fields modify degree of nonlocality
- Magnetic field evolution non-local
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### An idealised view....

**Model building** 

Mismatch motivates improved models

Integrated (ignition scale) experiments

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**Model building** 

Design simplified experiments Mismatch motivates improved models

Integrated (ignition scale) experiments

Simplified experiments







### An idealised view....

Data to impivalidate mod

Simplified experime v predictions

Integrated (ignition scale) experiments

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- Large-scale ICF experiments inherently multi-scale necessitates simple models for kinetic processes
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## Nonlocal transport experiment



M.D. Rosen et al., HEDP, 7, 180 (2011)

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M.D. Rosen et al., HEDP, 7, 180 (2011)

## Nonlocal transport experiment



G. Gregori et al., PRL, 92, 205006 (2004)

'Classic' SNB

$$\left[\frac{1}{\lambda_g} - \nabla \cdot \left(\frac{\lambda_g}{3}\nabla\right)\right] H_g = \nabla \cdot \boldsymbol{U}_g$$

$$\boldsymbol{q} = \boldsymbol{q}_{LTE} - \frac{m_e}{2} \int \frac{\lambda_g}{3} \nabla H_g dv$$

'Classic' SNB

Modified SNB

$$\left[\frac{1}{\lambda_g} - \nabla \cdot \left(\frac{\lambda_g}{3}\nabla\right)\right] H_g = \nabla \cdot \boldsymbol{U}_g$$

$$\left[\frac{1}{\lambda_{ee}^{*}} - \nabla \cdot \left(\frac{\lambda_{ei}^{*}}{3}\nabla\right)\right] H_{g} = \nabla \cdot \boldsymbol{U}_{g}$$

$$\boldsymbol{q} = \boldsymbol{q}_{LTE} - \frac{m_e}{2} \int \frac{\lambda_g}{3} \nabla H_g dv$$

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$$\lambda_{ee}^* = r \lambda_{ee} \qquad r = 2$$

'Classic' SNB

Modified SNB

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$$\lambda_{ee}^* = r \lambda_{ee} \qquad r = 2$$
$$\lambda_{ei}^* = f(Z) \lambda_{ei} \qquad f(Z) = \frac{Z + 0.24}{Z + 4.2}$$



J. Brodrick et al, Phys. Plas., 24, 092309 (2017)

Modified SNB

$$\left[\frac{1}{\lambda_{ee}^*} - \nabla \cdot \left(\frac{\lambda_{ei}^*}{3}\nabla\right)\right] H_g = \nabla \cdot \boldsymbol{U}_g$$

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![](_page_67_Figure_1.jpeg)

![](_page_68_Figure_1.jpeg)

![](_page_69_Figure_1.jpeg)

![](_page_70_Figure_1.jpeg)

K. Ma et al, in press

![](_page_71_Figure_1.jpeg)

K. Ma et al, in press
#### Nonlocal transport model



- 1D HYDRA mid-Z sphere Omega-like simulations
- Modified SNB improvement on classic SNB
- Comparison to K2 error in heat flux at 1ns

K. Ma et al, in press

#### **Magnetic fields & nonlocal transport**



W.A. Farmer et al., Phys. Plasmas, 24, 052703 (2017)

• HYDRA high-foot simulation (CH capsule 0.6mg/cc fill) t=13ns

#### **Magnetic fields & nonlocal transport**





- Magnetic field deflects heat flow
  - 1. Restricts conduction
  - 2. Righi-Leduc

W.A. Farmer et al., Phys. Plasmas, 24, 052703 (2017)





D.H. Froula et al., PRL, 98, 135001 (2007)













C.P. Ridgers et al., PRL, 100, 075003 (2008)

# **Nonlocal transport experiment** $\frac{\partial \mathbf{B}}{\partial t} \approx -\frac{1}{en_e} \nabla n_e \times \nabla T_e - \nabla \cdot (\mathbf{v}_N B_z) \qquad \text{IMPACT VFP code} \\ - \text{f0+f1 code}$

• In simplified geometry with  $B=(0,0,B_{r})$ 

$$\frac{\partial B_z}{\partial t} \approx -\nabla \cdot (\boldsymbol{v}_N B_z) \quad \boldsymbol{v}_N \propto \boldsymbol{q}_e$$

C.P. Ridgers et al., PRL, 100, 075003 (2008)

- 3D (2 spatial, 1 velocity)

#### **Nonlocal transport experiment** $\frac{\partial \mathbf{B}}{\partial t} \approx -\frac{1}{en_{e}} \nabla n_{e} \times \nabla T_{e} - \nabla \cdot (\mathbf{v}_{N} \mathbf{B}_{z})_{1.8}$ B-field at 1.6 Magnetic Field (Normalized to Unity) 440ps 1.4 • In simplified geometry with 1.2 B=(0,0,B)0.8 $\frac{\partial B_z}{\partial t} \approx -\nabla \cdot (\boldsymbol{v}_N B_z) \quad \boldsymbol{v}_N \propto \boldsymbol{q}_e$ B=2T (Frozen) 0.6 -+ B=4T (Full) 0.4 - \* - B=4T (Frozen) 0.2 B=12T Frozen 0 100 200 300 400 500 0

Radial Distance/Microns

C.P. Ridgers et al., PRL, 100, 075003 (2008)

• Gas-jet target  $\rightarrow$  easily change collisionality ( $n_e$ ~5x10<sup>18</sup>-5x10<sup>19</sup>cm<sup>-3</sup>)



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- Impose a magnetic field (~3T)



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- Impose a magnetic field (~3T)
- Diagnose B-field using proton radiography





$$\frac{\partial B_z}{\partial t} \approx -\nabla \cdot (\boldsymbol{v}_N B_z) \quad \boldsymbol{v}_N \propto \boldsymbol{q}_e$$





Proton radiograph and inteferogram at 1.1ns (10<sup>19</sup>cm<sup>-3</sup>) C. Arran et al., 100, arXiV:2105.07414 (2021)





Proton radiograph and inteferogram at 1.1ns (10<sup>19</sup>cm<sup>-3</sup>) C. Arran et al., 10





Proton radiograph and inteferogram at 1.1ns (10<sup>19</sup>cm<sup>-3</sup>) C. Arran et al., 10



#### Nonlocality in magnetised plasmas



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## **Magnetic fields & transport** $\frac{\partial \mathbf{B}}{\partial t} \approx -\frac{1}{en_e} \nabla n_e \times \nabla T_e - \nabla \times (\mathbf{v}_N \times \mathbf{B})$



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- Biermann-producing electric field is  $e n_e E \approx -\nabla P_e$
- But this assumes LTE

R.J. Kingham & A.R. Bell, PRL, 88, 045004 (2002)

- Biermann-producing electric field is  $e n_e \mathbf{E} \approx -\nabla P_e$
- But this assumes LTE
- Relax this assumption  $e n_e \mathbf{E} \approx -\frac{m_e}{6} \frac{\nabla \langle n_e v^5 \rangle}{\langle v^3 \rangle}$

R.J. Kingham & A.R. Bell, PRL, 88, 045004 (2002)

 $\nabla P_{e}$  ..mes LTE ..mes ..• Biermann-producing dB/dt (simulation) electric field is Power law fit • But this assumes LTE • Relax this assumption 0.1

M. Sherlock & J.J Bissell, PRL, 124 055001 (2020)









 Laser intensity 2x10<sup>14</sup>Wcm<sup>-2</sup>, spot size 33µm. Gas density 1.5x10<sup>19</sup>cm<sup>-3</sup>



- Laser intensity 2x10<sup>14</sup>Wcm<sup>-2</sup>, spot size 33µm. Gas density 1.5x10<sup>19</sup>cm<sup>-3</sup>
- 1D IMPACT simulations








C.P. Ridgers et al., Phil. Trans. R. Soc. A, 379, 20200079 (2020)













- 2D MHD simulations with density gradient
- B-field after 60ps



- 2D MHD simulations with density gradient
- B-field after 60ps





intensity map)



# M1 model - work in progress

- M1 model similar to SNB (more complicated)
- Naturally includes B-fields
- LEARN project to assess accuracy

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- M1 model similar to SNB (more complicated)
- Naturally includes B-fields
- LEARN project to assess accuracy
- Relaxation of sine-wave is promising
- LDRD (M. Sherlock) to develop robust, fast VFP





• Ongoing analysis Heater beam of Vulcan expt. 5° Blade mm May see kinetic 8 3.5 mm Biermann 1.04 mm Interferometry probe 10MeV Gas jet beam (2w, 1ps) protons Gas-jet (out of page) . Heater beams (1w, 1ns) Thomson scattering diagnostic (2w, 0.5ns)





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- Full kinetic codes too slow but reduced models exist
- We can benchmark these using full kinetic codes and simplified experiments
- Need to include magnetic fields
- No nonlocal model incuding B-fields has been benchmarked...

• Scrape-Off Layer (SOL) - 'open' field lines outside core plasma onto divertor



- Scrape-Off Layer (SOL) 'open' field lines outside core plasma onto divertor
- Heat load on divertor limitation for high power devices



- Scrape-Off Layer (SOL) 'open' field lines outside core plasma onto divertor
- Heat load on divertor limitation for high power devices
- Transport nonlocal in SOL







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