New Perspectives for the *ab-initio* Simulation and Diagnostics of Warm-dense Matter

T. Dornheim^{1,2}, J. Vorberger², Zh. Moldabekov^{1,2}, M. Böhme^{1,2,3}, K. Ramakrishna^{1,2,3}, M. Bonitz⁴, D. Kraus^{5,2}, T. Döppner⁶, T. Preston⁷, P. Tolias⁸



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Warm Dense Matter (WDM)

• Matter under extreme density / temperature ubiquitous throughout our universe

r_s ~ θ ~ Γ ~ 1

 $\rightarrow r_s = d/a_B$, density parameter, $\theta = k_B T/E_F$, $\Gamma = W/E_{kin}$







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- WDM highly important for technological applications:
- \rightarrow Inertial confinement fusion, etc.



National Ignition Facility (NIF)



Taken from: Lawrence Livermore National Laboratory

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<u>facilities ar</u>	ound the globe!
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But: Rigorous WDM theory indispensable

• <u>Diagnostics</u>: parameters like *T*, *n*, *Z*, etc. cannot be measured and have to be inferred from theory

 \rightarrow X-ray Thomson scattering (XRTS)

Isochorically heated graphite at LCLS (Stanford)



<u>Taken from:</u> D. Kraus *et al.*, *Plasma Phys. Control. Fusion* **61**, 014015 (2019)

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Ab-initio Quantum Monte Carlo (QMC) simulations

Problem:

• **Density functional theory (DFT)** etc. require external input about XC-effects

 \rightarrow finite *T*: XC-<u>free</u> energy f_{xc}

Solution:

• Quantum Monte Carlo methods in principle allow for exact solution of quantum many-body problems <u>without</u> any empirical input

• Finite T: Path Integral Monte Carlo (PIMC)

Taken from: **T. Dornheim**, S. Groth, and M. Bonitz, *Contrib. Plasma Phys.* **59**, e201800157 (2019)

Previous result: XC-free energy of UEG

S. Groth, T. Dornheim, T. Sjostrom, F.D. Malone, W.M.C. Foulkes, and M. Bonitz, PRL 119, 135001 (2017)

Impact on thermal DFT simulation of warm dense hydrogen

Example:

Hydrogen at T=65,000K

r_s = 2

(a) Ground-state LDA by PerdewAnd Zunger, PRB (1980) [PZ](b) our thermal LDA

[GDSMFB]

Taken from: K. Ramakrishna, **T. Dornheim**, and J. Vorberger, *Phys. Rev. B* **101**, 195129 (2020)

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Example:

Part II: Density response of warm dense electrons

Density response functions, local field correction

• Dynamic density response function

$$\chi(q,\omega) = \frac{\chi_0(q,\omega)}{1 - \frac{4\pi}{q^2} [1 - G(q,\omega)] \chi_0(q,\omega)}$$

- $\rightarrow \chi_0(q,\omega)$ ideal density response function
- → $G(q, \omega)$ dynamic local field correction, <u>containing all</u> <u>electronic XC-effects</u>

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Important input for many applications:

- \rightarrow Advanced nonlocal XC-functionals for DFT
- \rightarrow Stopping power, electronic friction, ...
- \rightarrow Effective potentials
- \rightarrow Electrical/thermal conductivity
- \rightarrow Interpretation of XRTS experiments
- → ...

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$$F(\mathbf{q},\tau) = \frac{1}{N} \left\langle \rho(\mathbf{q},\tau) \rho(-\mathbf{q},0) \right\rangle$$

Taken from: **T. Dornheim**, T. Sjostrom, S. Tanaka, and J. Vorberger, Phys. Rev. B **101**, 045129 (2020)

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- Extensive PIMC data for LFC G(q) for ~50 r_{s} - θ combinations

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First results for XC-kernel of hydrogen:

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M. Böhme, Zh. Moldabekov, J. Vorberger, and T. Dornheim, *Phys. Rev. Lett.* 129, 066402 (2022)

Density response of warm dense hydrogen

- Exact exchange-correlation kernel of hydrogen (PIMC)
 - \rightarrow Benchmark Adiabatic LDA (ALDA) etc
 - → Influence of partial localization around ions?
 - → ...
- UEG models break down at low density

Exchange—correlation kernel of warm dense hydrogen

Taken from: M. Böhme, Zh. Moldabekov, J. Vorberger, and **T. Dornheim**, *Phys. Rev. Lett.* **129**, 066402 (2022)

• Future works:

- \rightarrow input for time-dependent DFT, etc
- \rightarrow development of new XC-functionals
- \rightarrow predict NIF experiments
- $\rightarrow \dots$

Density response of real materials

- Compute exchange—correlation kernel from DFT simulations
 - \rightarrow <u>Problem</u>: DFT limited to singe-electron density
 - \rightarrow <u>Solution</u>: Perturb system, compute density response
- → DFT gives <u>access</u> to exchange—correlation kernel

Taken from: Zh. Moldabekov, M. Böhme, J. Vorberger, D. Blaschke, and **T. Dornheim**, arXiv:2209.00928

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- DFT capable to generate electronic XC-effects
 - \rightarrow insights into performance of XC-functionals
 - \rightarrow XC-effects of <u>real materials</u>

 $\rightarrow \dots$

Exchange—correlation kernel of warm dense hydrogen

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Need for dynamic properties of WDM

- WDM Diagnostics: obtain plasma parameters from XRTS experiments
- → <u>Dynamic structure factor</u>

$$F(q,t) = \frac{1}{N} \langle \rho(q,t)\rho(-q,0) \rangle$$
$$\Rightarrow S(q,\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} dt \ F(q,t) \ e^{i\omega t}$$

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- Rigorous description of dynamic properties even more challenging then TD equilibrium
- \rightarrow TD-DFT: adiabatic approximation, no XC-kernel
- \rightarrow Green functions: approximation in coupling
- → PIMC: Imaginary time, analytic continuation possible for UEG!

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First *ab-initio* results for dynamic structure factor of UEG

RPA: G(**q**,ω)=0

SLFC: $G(q,\omega) = G(q,0)$

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DLFC: exact solutions

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- Electron liquid (rs=10):
- \rightarrow non-trivial shape of S(**q**, ω)
- \rightarrow negative dispersion relation

First *ab-initio* results for dynamic behaviour of WDM

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Exampe: UEG at **0**=1

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First *ab-initio* results for dynamic behaviour of WDM

 \rightarrow negative dispersion relation

X-ray Thomson scattering (XRTS)

• <u>Standard way</u>: construct a model for $S(q,\omega)$, convolve with instrument function $R(\omega)$, fit to XRTS signal $I(q,\omega)$

$$I(\mathbf{q},\omega) = S(\mathbf{q},\omega) \circledast R(\omega)$$

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- \rightarrow no direct access to physical information
- \rightarrow approximate theoretical models

Taken from: **T. Dornheim**, Zh. Moldabekov, P. Tolias, M. Böhme, and J. Vorberger, arXiv:2209.02254 (submitted)

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- Imaginary-time domain:
- \rightarrow direct access to physics, e.g. T, ω_{p}
- \rightarrow exact QMC simulations

$$\mathcal{L}\left[S(\mathbf{q},\omega)\right] = \int_{-\infty}^{\infty} \mathrm{d}\omega \ e^{-\tau\omega} \ S(\mathbf{q},\omega)$$

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Model-free temperature from XRTS experiments:

• Detailed balance in the τ-domain:

 $S(\mathbf{q},-\omega) = S(\mathbf{q},\omega)e^{-\beta\omega}$

- \rightarrow works for all wave numbers
- \rightarrow no explicit resolution of plasmon required

Laplace transform:

$$\mathcal{L}\left[S(\mathbf{q},\omega)\right] = \int_{-\infty}^{\infty} \mathrm{d}\omega \ e^{-\tau\omega} \ S(\mathbf{q},\omega)$$

 \rightarrow symmetry around $\tau = (2T)^{-1}$

Taken from: **T. Dornheim** *et al.*, in preparation

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 \rightarrow finite $\omega\text{-range},$ check convergence with x

$$\mathcal{L}_x\left[S(\mathbf{q},\omega) \circledast R(\omega)\right] = \int_{-x}^{x} \mathrm{d}\omega \ e^{-\tau\omega} \left\{S(\mathbf{q},\omega) \circledast R(\omega)\right\}$$

$$\mathcal{L}\left[S(\mathbf{q},\omega)\right] = \frac{\mathcal{L}\left[S(\mathbf{q},\omega) \circledast R(\omega)\right]}{\mathcal{L}\left[R(\omega)\right]}$$

Temperature of warm dense Be [Glenzer (2007)]

Model-free temperature from XRTS experiments:

- Detailed balance in the τ-domain:
- \rightarrow works for all wave numbers

3

 $E_0=2.96$ keV

 \rightarrow no explicit resolution of plasmon required

Experiment

model. T=12eV

 \rightarrow finite $\omega\text{-range},$ check convergence with x

C)

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Taken from: T. Dornheim, M. Böhme, D. Kraus, T. Döppner, T. Preston, Zh. Moldabekov, and J. Vorberger, arXiv:2206.12805

Temperature of warm dense Be [Glenzer (2007)]

100

b)

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Summary and Outlook

Ab initio theory of WDM

- need for finite-T XC functionals based on PIMC results
- static density response: PIMC + neural net
- PIMC results for warm dense hydrogen
- DFT framework for the study of XC-effects
- dynamic density response: PIMC + analytic continuation

Key pre-print:

arXiv:2209.00928

Taken from: **T. Dornheim**, J. Vorberger, S. Groth, N. Hoffmann, Zh. Moldabekov, and M. Bonitz, JCP **151**, 194104 (2019)

Summary and Outlook

Physics in the imaginary time

- Usual ω -representation equivalent to τ -domain
- Model-free T-diagnostics etc.
- Future works: physical insights form the τ -domain

Key pre-prints:

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